

A FORTRAN code for $\gamma\gamma\rightarrow ZZ$ in SM and MSSM

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Based on the publications:

- Th. Diakonidis , GJG, J. Layssac, hep-ph/0610085
- GJG, P. Porfyriadis, F.M. Renard, Eur. Phys. J. C19:57 (2001), hep-ph/0010006.
- GJG, P. Porfyriadis, J. Layssac, F.M. Renard, Eur. Phys. J. C13:79(2000), hep-ph/9909243.
- **The whole code is contained in the file gamgamZZ.tar.gz downloaded from**

<http://users.auth.gr/~gounaris/FORTRANcodes/>

It contains 4 sub-codes called, sm1, mssm1, sm2 and mssm2.

A Readme.dat file explains everything.

Real SM and MSSM parameters are assumed.

The $\gamma(\lambda_1) \gamma(\lambda_2) \rightarrow Z(\lambda_3) Z(\lambda_4)$ helicity amplitudes

$\lambda_j = \text{helicities}$

Helicity amplitudes $\Rightarrow F_{\lambda_1 \lambda_2 \lambda_3 \lambda_4}(\beta_Z, t, u)$, $s = s_{\gamma\gamma} = \frac{4m_Z^2}{1 - \beta_Z^2}$

Bose $\Rightarrow F_{\lambda_1 \lambda_2 \lambda_3 \lambda_4}(\beta_Z, t, u) = F_{\lambda_2 \lambda_1 \lambda_4 \lambda_3}(\beta_Z, t, u) (-1)^{\lambda_3 - \lambda_4}$

$$F_{\lambda_1 \lambda_2 \lambda_3 \lambda_4}(\beta_Z, t, u) = F_{\lambda_2 \lambda_1 \lambda_3 \lambda_4}(\beta_Z, u, t) (-1)^{\lambda_3 - \lambda_4}$$

$$F_{\lambda_1 \lambda_2 \lambda_3 \lambda_4}(\beta_Z, t, u) = F_{\lambda_1 \lambda_2 \lambda_4 \lambda_3}(\beta_Z, u, t)$$

CP $\Rightarrow F_{\lambda_1 \lambda_2 \lambda_3 \lambda_4}(\beta_Z, t, u) = F_{-\lambda_1, -\lambda_2, -\lambda_3, -\lambda_4}(\beta_Z, t, u) (-1)^{\lambda_3 - \lambda_4}$

Thus, the 10 independent helicity amplitudes are:

$$F_{++++}, F_{+--+}, F_{+-00}$$

$$F_{++++-}, F_{+---+}, F_{++00}$$

$$F_{+++0}, F_{++--}, F_{++-0}$$

$$F_{+-+0}$$

The three **magenta** amplitudes respect the asymptotic **Helicity Conservation (HC)** rule

$$\lambda_1 + \lambda_2 = \lambda_3 + \lambda_4$$

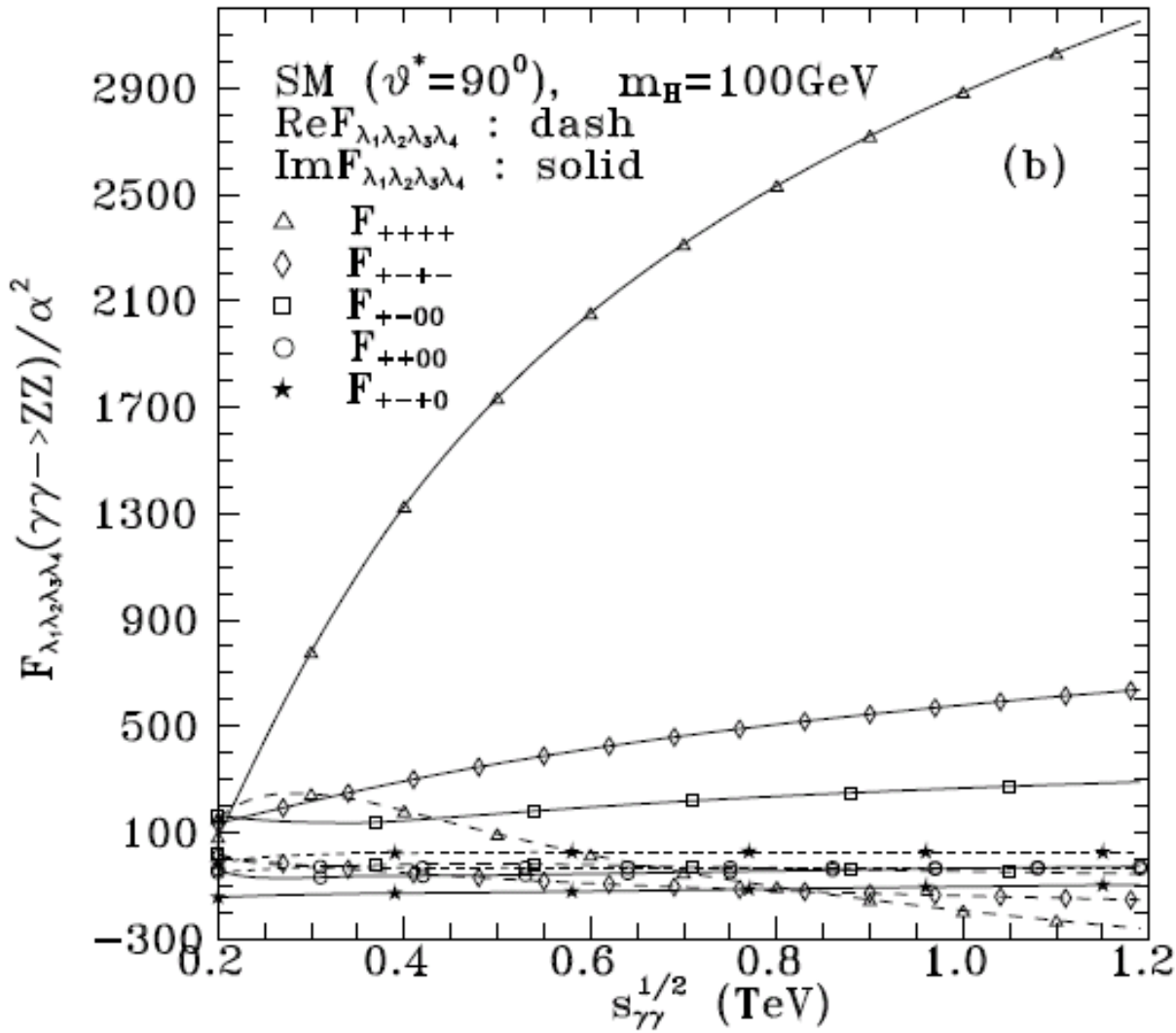
They are the largest ones, and **mostly imaginary**.

The rest of these amplitudes violate **HC** and they are very small above 200GeV.

HC: Renard + G, hep-ph/0501046, hep-ph/0604041

Codes **sm1** and **mssm1** calculate these amplitudes for **SM** and **MSSM** respectively.

Example...



Imaginary and real parts of the helicity amplitudes in SM at 90° .

HC respecting amplitudes mostly imaginary.

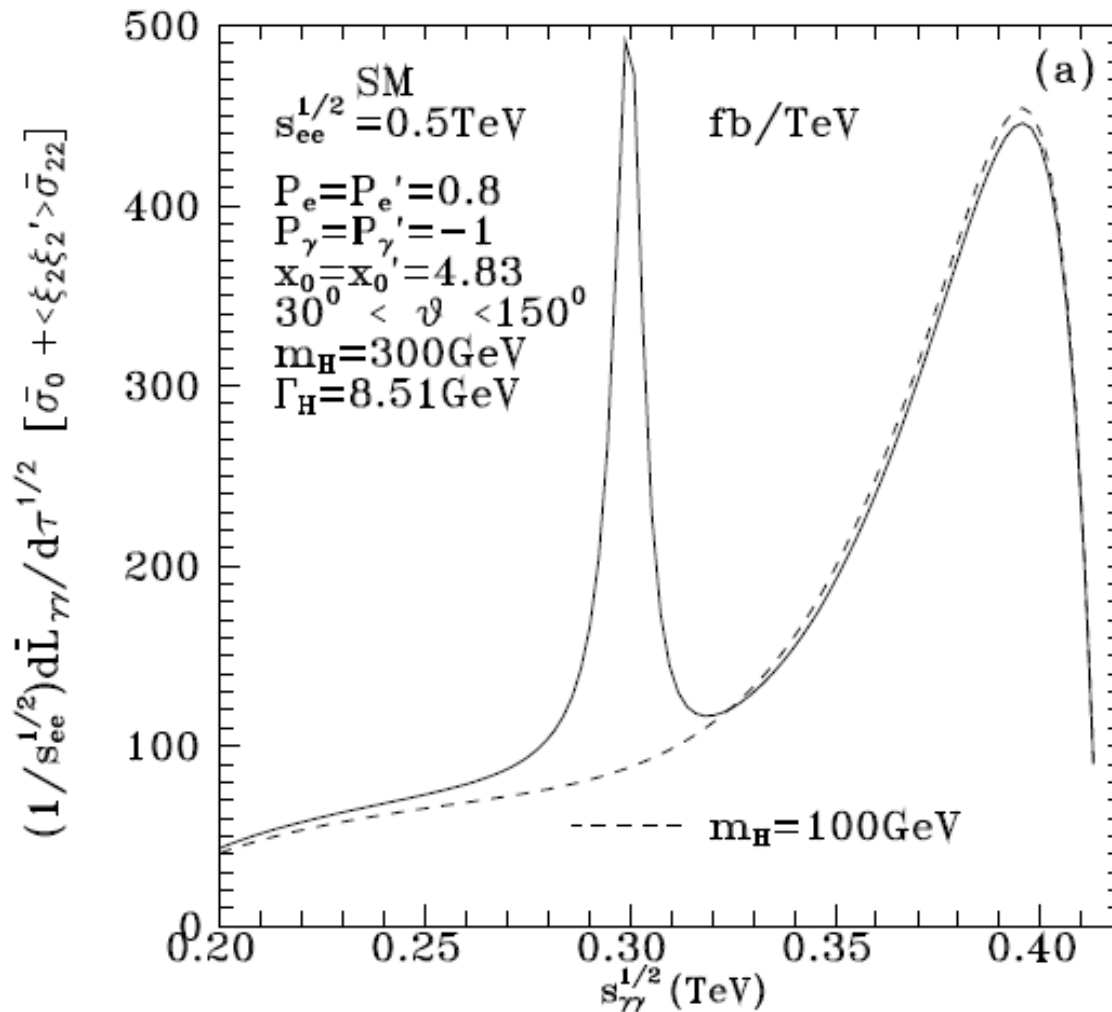
HC violating amplitudes very small at high (s, t)

Cross sections for unpolarized Z's, given by the codes **sm2** and **mssm2** for **SM** and **MSSM** respectively.

These are:

$$\begin{aligned} \frac{d\sigma_0(\gamma\gamma \rightarrow ZZ)}{d\cos\theta} &= \left(\frac{\beta_Z}{128\pi\hat{s}} \right) \sum_{\lambda_3\lambda_4} [|F_{++\lambda_3\lambda_4}|^2 + |F_{+-\lambda_3\lambda_4}|^2], \\ \frac{d\sigma_{22}(\gamma\gamma \rightarrow ZZ)}{d\cos\theta} &= \left(\frac{\beta_Z}{128\pi\hat{s}} \right) \sum_{\lambda_3\lambda_4} [|F_{++\lambda_3\lambda_4}|^2 - |F_{+-\lambda_3\lambda_4}|^2], \\ \frac{d\sigma_3(\gamma\gamma \rightarrow ZZ)}{d\cos\theta} &= \left(\frac{-\beta_Z}{64\pi\hat{s}} \right) \sum_{\lambda_3\lambda_4} \text{Re}[F_{++\lambda_3\lambda_4}F_{-+\lambda_3\lambda_4}^*], \\ \frac{d\sigma_{33}(\gamma\gamma \rightarrow ZZ)}{d\cos\theta} &= \left(\frac{\beta_Z}{128\pi\hat{s}} \right) \sum_{\lambda_3\lambda_4} \text{Re}[F_{+-\lambda_3\lambda_4}F_{-+\lambda_3\lambda_4}^*], \\ \frac{d\sigma'_{33}(\gamma\gamma \rightarrow ZZ)}{d\cos\theta} &= \left(\frac{\beta_Z}{128\pi\hat{s}} \right) \sum_{\lambda_3\lambda_4} \text{Re}[F_{++\lambda_3\lambda_4}F_{--\lambda_3\lambda_4}^*], \\ \frac{d\sigma_{23}(\gamma\gamma \rightarrow ZZ)}{d\cos\theta} &= \left(\frac{\beta_Z}{64\pi\hat{s}} \right) \sum_{\lambda_3\lambda_4} \text{Im}[F_{++\lambda_3\lambda_4}F_{+-\lambda_3\lambda_4}^*] \end{aligned}$$

Cross sections for polarized e^\pm beams, integrated over the azimuthal angles in SM



$$\bar{\sigma}_j \equiv \sigma_j$$

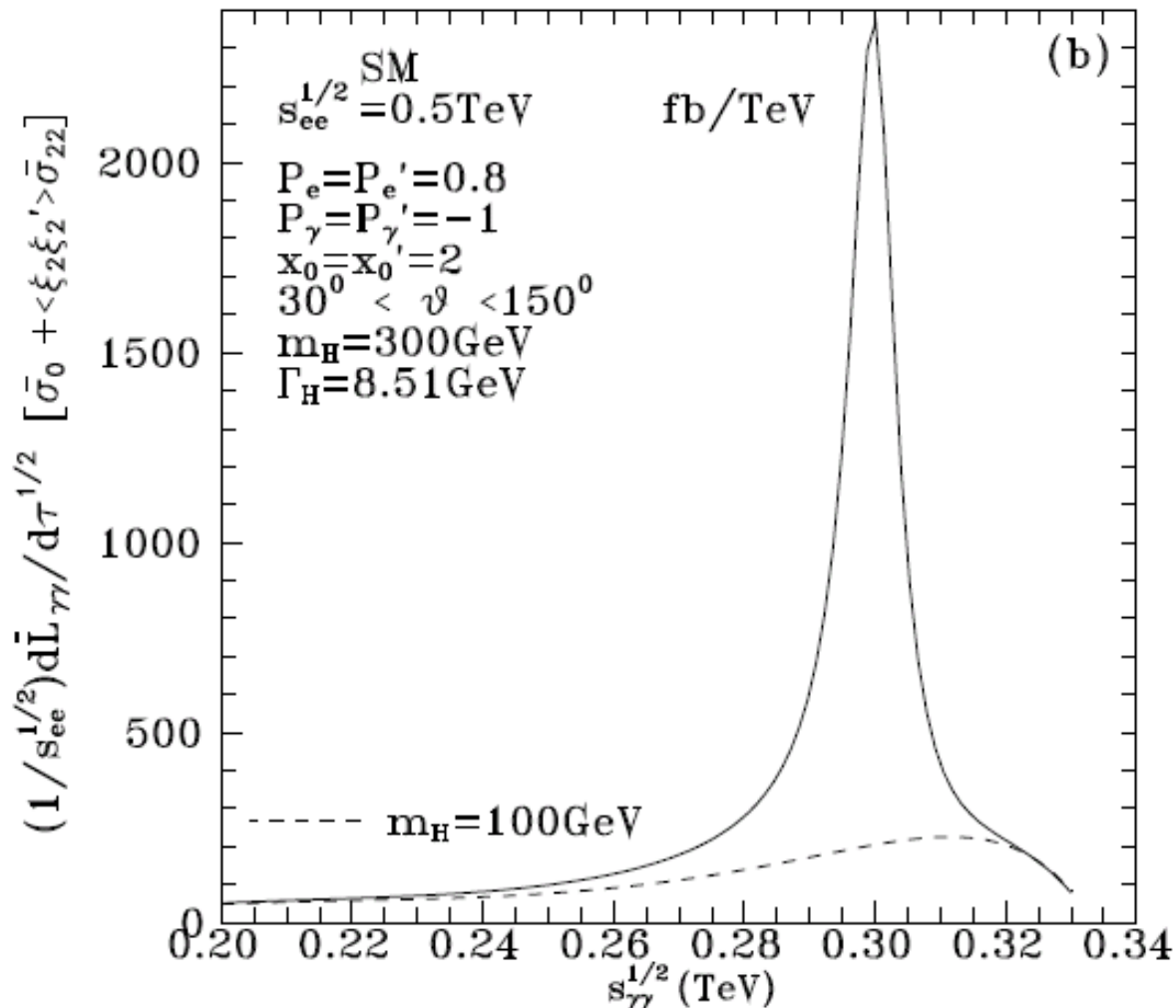
$$\tau = \frac{s_{\gamma\gamma}}{4E_e^2}$$

$$x_0 = \frac{4E_e \omega_{\text{laser}}}{m_e^2}$$

E_e = energy of
 the e^\pm beams.

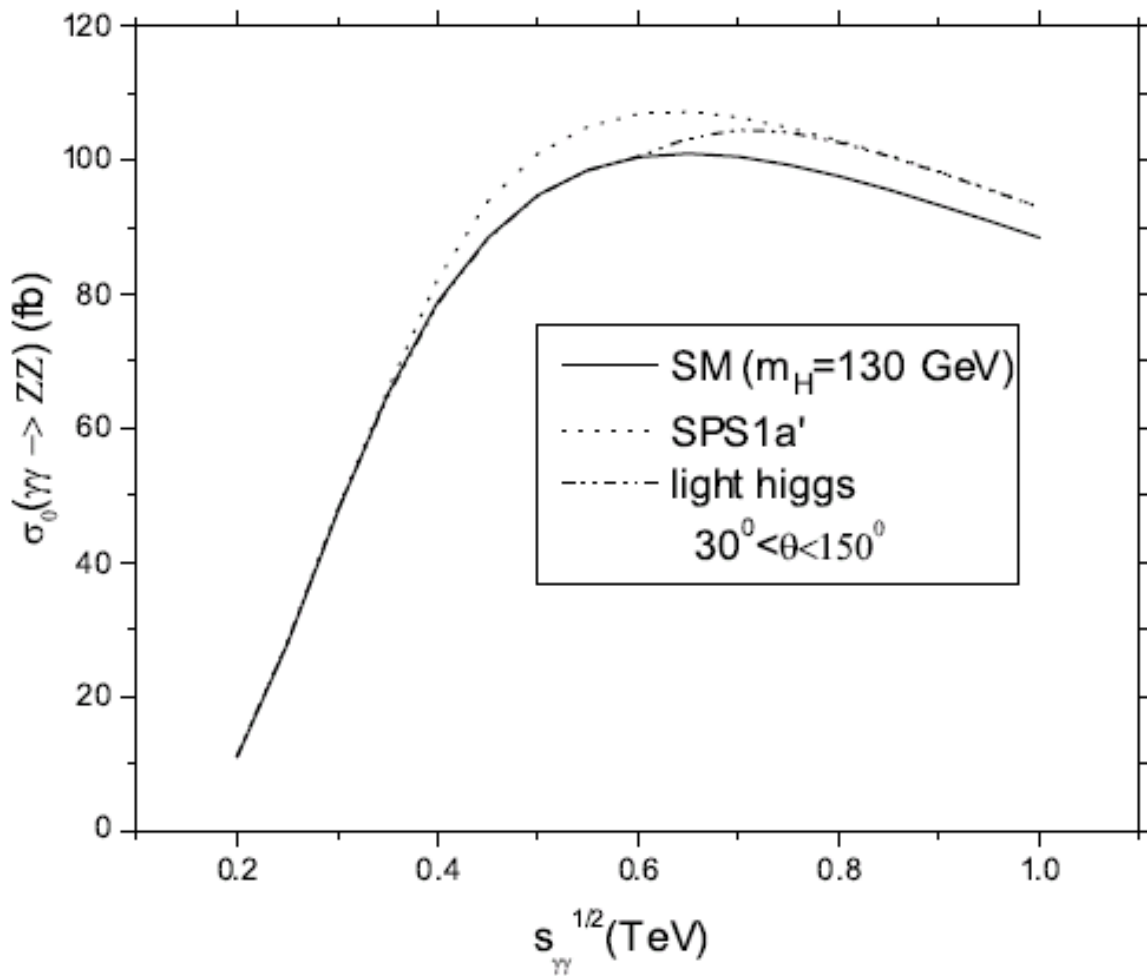
**A Higgs peak may
 be visible in SM,
 if m_H large.**

Cross sections for polarized e^\pm beams, integrated over the azimuthal angles in SM



Effect stronger for smaller x_0 .

σ_0 cross sections in MSSM and SM



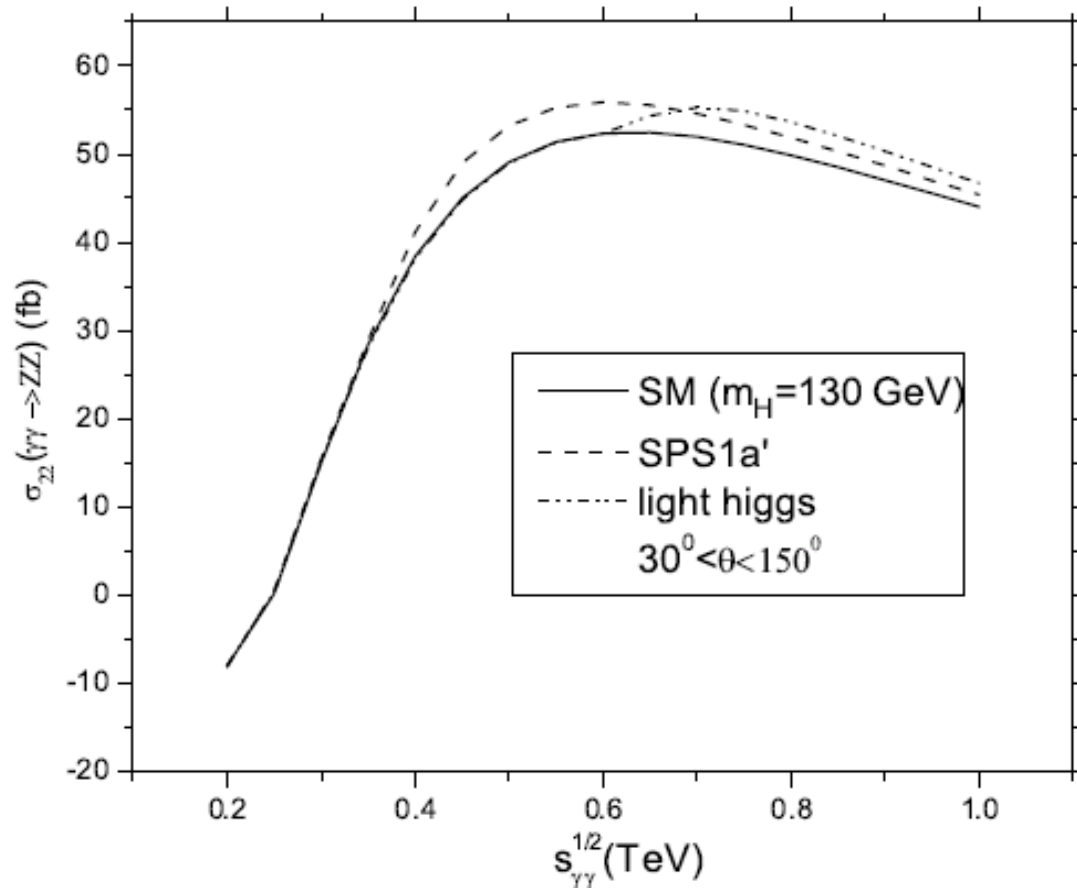
SPS1a'

$m_0=70\text{GeV}$,
 $m_{1/2}=250\text{GeV}$,
 $A_0=-300\text{GeV}$,
 $\tan\beta=10$, $\mu>0$

Here H^0 lies at 424GeV , but it couples so weakly, and its width is so large, that no peak is visible in σ_0 , (integrated over the indicated angular range).

“light higgs” is an MSSM model from hep-ph/0609079, where H^0 is below the ZZ threshold.

σ_{22} cross sections in MSSM and SM



Again, no H^0 peak is visible in σ_{22} , in SPS1a'.

Nevertheless, there exist MSSM examples, where H^0 peaks may be visible in both, σ_0 and σ_{22} .

Summary

The purpose of this talk is to inform the community that a code called **gamgamZZ.tar.gz** exists in

<http://users.auth.gr/~gounaris/FORTRANcodes/>

which calculates:

- All helicity amplitudes for $\gamma(\lambda_1) \gamma(\lambda_2) \rightarrow Z(\lambda_3) Z(\lambda_4)$.

This should be useful in cases where Z-polarization effects must be studied.

- All cross sections where Z-polarizations are summed over. These give all physical observables, in case Z-polarization is not looked at.
- MSSM parameters are assumed real.
- Please send me suggestions for improvements...