

The origin of matter in the nMSSM

or

Why do we need the ILC?

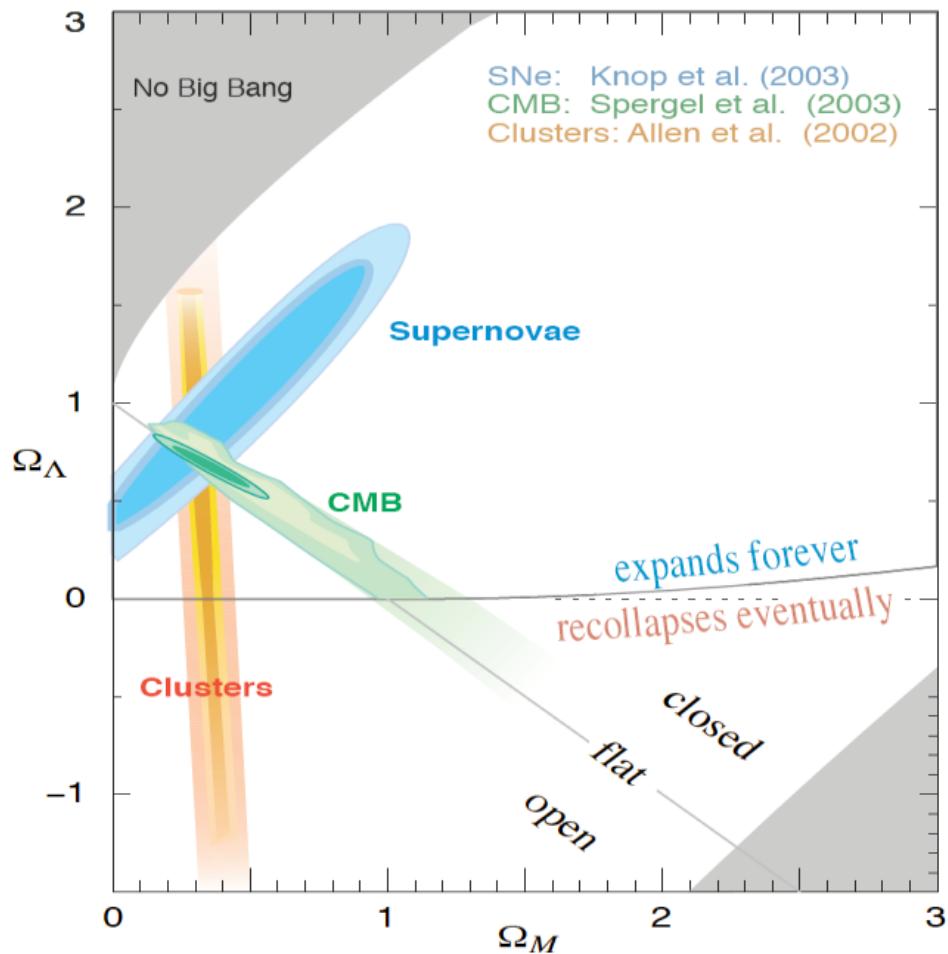
- Origin of matter in the MSSM
- Origin of matter in the nMSSM
- Need for the ILC

C.Balázs, M.Carena, A. Freitas, C.Wagner hep-ph/0606000

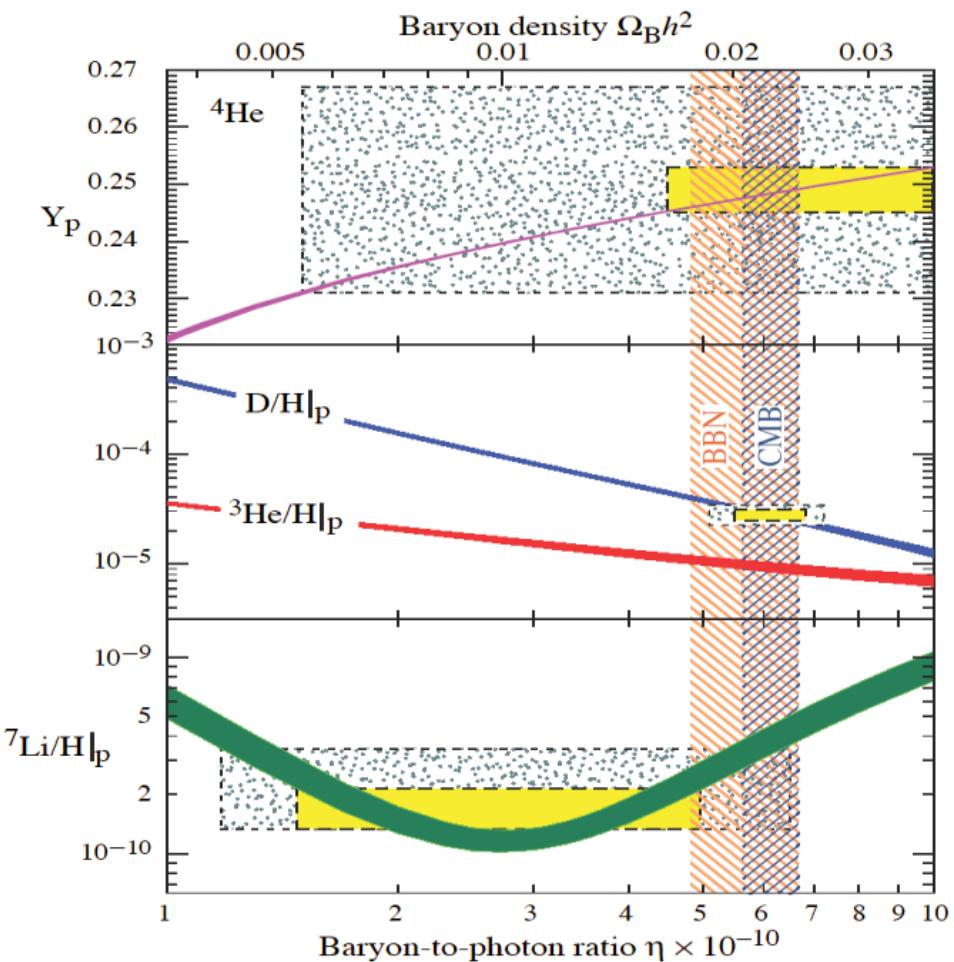
C.Balázs, M.Carena, A. Menon, D.Morrissey, C.Wagner PRD71 075002 ('05)

C.Balázs, M.Carena, C.E.M.Wagner PRD70 015007 ('04)

The matter content of the Universe



$$\Omega_{\text{CDM}} = 0.27 \pm 0.04$$

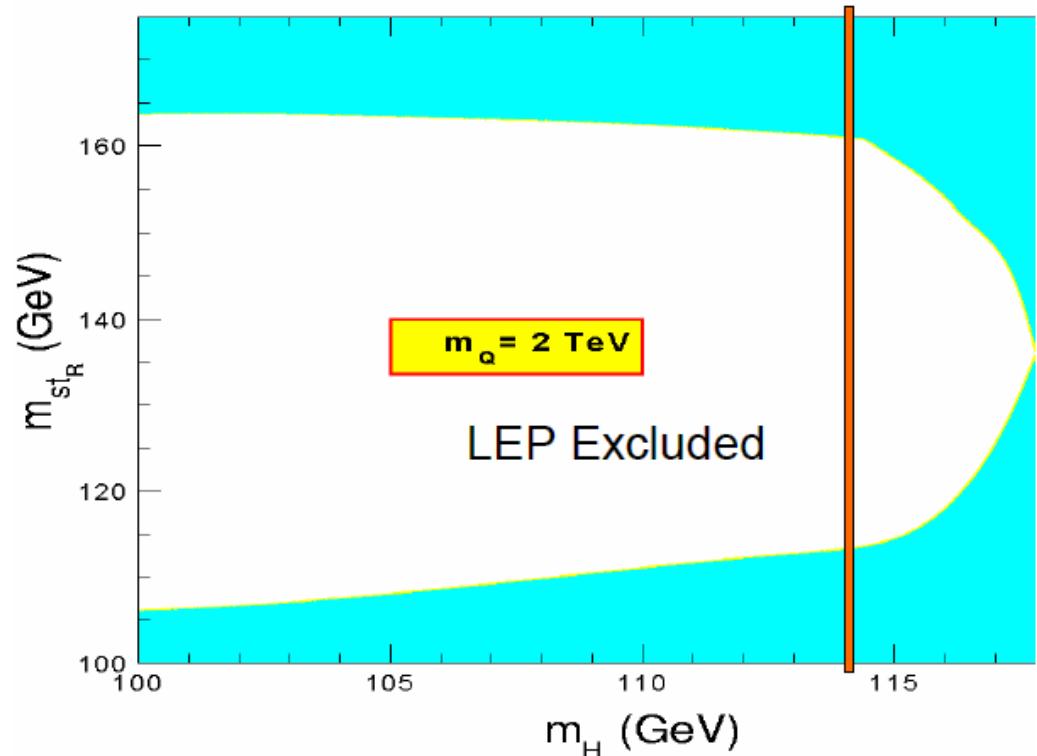


$$\Omega_B = 0.044 \pm 0.001$$

unexplained by the standard models

Baryons in the MSSM

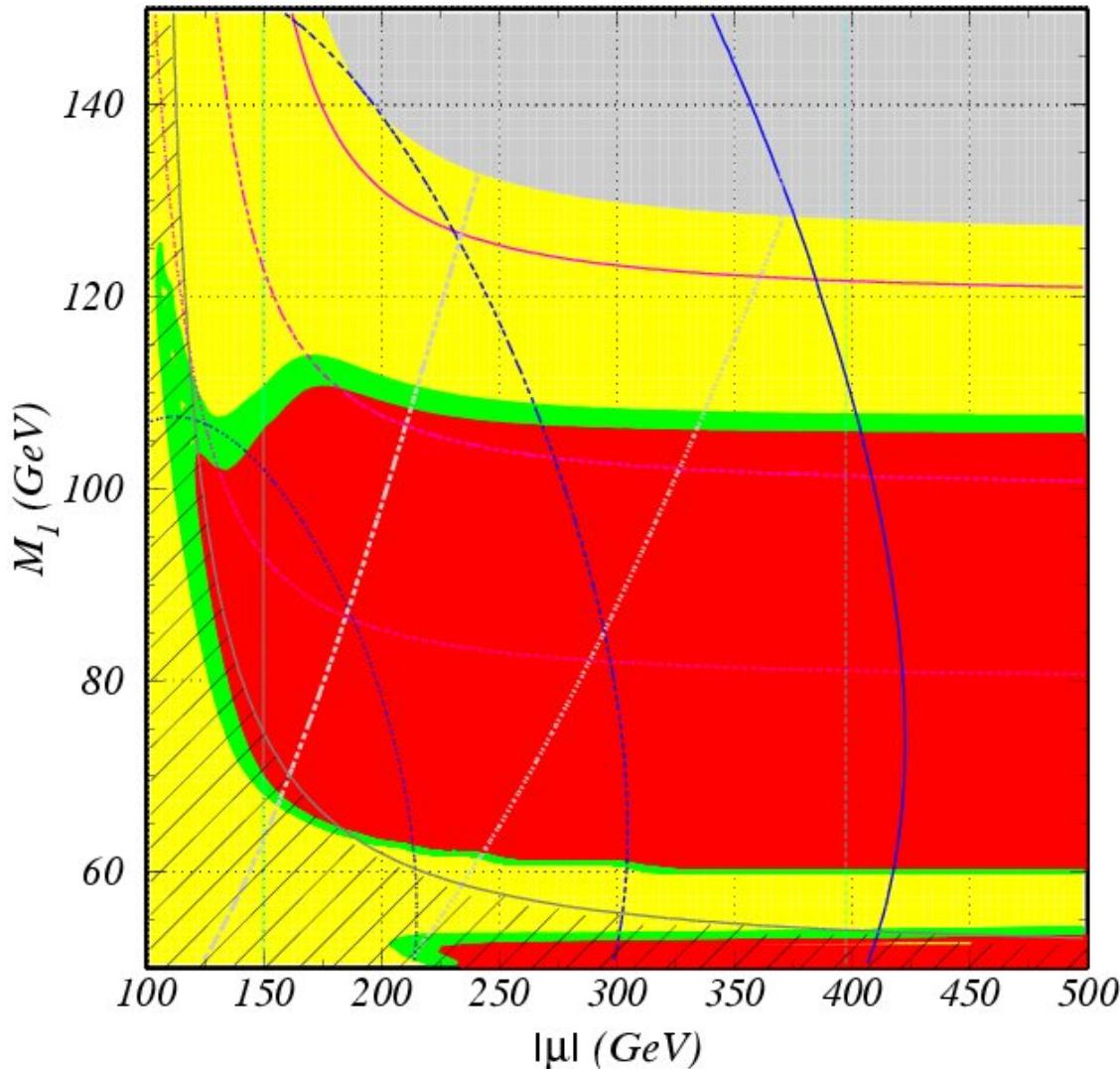
- Electroweak baryogenesis
 - explains Ω_B
 - constrains MSSM parameter space
- EW phase transition
strongly 1st order →
a light stop & higgs
 $m_{\tilde{t}_1} < m_t$, $m_{\tilde{t}_2} \gtrsim 1 \text{ TeV}$,
 $0.3 < |x_t| / m_{\tilde{Q}_3} < 0.5$,
 $m_h \lesssim 120 \text{ GeV}$
- Enough CP →
light charginos & μ
 $M_2, \mu \lesssim 500 \text{ GeV}$,
 $\text{Arg}(M_2 \mu) \gtrsim 0.1$
- EDM limits → heavy 1st & 2nd generation scalars



Carena, Seco, Quiros, Wagner 2002

All matter in the MSSM

- $\Omega_B \notin \Omega_{\text{CDM}}$ simultaneously predicted in MSSM



Input parameters:

$\tan\beta = 7, m_A = 1000 \text{ GeV}, \text{Arg}(\mu) = 1.571$
 $M_2 = M_1 g_2^2/g_1^2, \text{Arg}(M_1) = \text{Arg}(M_2) = 0, M_3 = 1 \text{ TeV}$
 $m_{U3} = 0 \text{ GeV}, m_{Q3} = 1.5 \text{ TeV}, X_t = 0.7 \text{ TeV}$
 $m_{L3}, m_{E3}, m_{D3} = 1 \text{ TeV}$
 $m_{L1,2}, m_{E1,2} = 10 \text{ TeV}$
 $m_{Q1,2}, m_{U1,2}, m_{D1,2} = 10 \text{ TeV}$

Legend:

| | | | |
|-----------------|------------------------------|-----------------------------|--|
| | $m_{t1} > m_{Z1}$ | | $m_{W1} < 103.5 \text{ GeV}$ |
| | $\Omega h^2 > 0.129$ | | $\Omega h^2 < 0.095$ |
| | $0.095 < \Omega h^2 < 0.129$ | | |
| $\sigma_{si} =$ | <u>$3E-08$</u> | <u>$3E-09$</u> | <u>$3E-10 \text{ pb}$</u> |
| $m_{Z1} =$ | <u>120</u> | <u>100</u> | <u>80 GeV</u> |
| $d_e =$ | <u>$1E-27$</u> | <u>$1.2E-27$</u> | <u>$1.4E-27 \text{ e cm}$</u> |

Balázs,Carena,Menon,Morrissey,Wagner 2005

nMSSM

- Discrete symmetries of super- $\&$ Kahler potential

$$Z_5^R, Z_7^R \subset U(1)_{R'} \quad \text{where} \quad R' = 3R + PQ$$

to prevent domain walls and large tadpoles

- Superpotential

$$W = W_{\text{MSSM}} + \lambda \hat{S} H_1 \cdot H_2 + \frac{m_{12}^2}{\lambda} \hat{S}$$

- Scalar potential

$$V = V_{\text{MSSM}} + m_S^2 |S|^2 + t_S(S + h.c.) + \\ a_\lambda(S H_1 \cdot H_2 + h.c.)$$

- New parameters

$$v_S, \lambda, a_\lambda, m_S, m_{12}, t_S$$

nMSSM

- Solves μ problem naturally

$$W = W_{\text{MSSM}} + \lambda \hat{S} H_1 \cdot H_2 + \frac{m_{12}^2}{\lambda} \hat{S}$$

- $\mu = \lambda \langle S \rangle = \lambda v_S$ set by EW scale

- Alleviates fine tuning in Higgs/stop sector

$$m_h^2 \leq m_Z^2 \left(\cos^2 2\beta + \frac{2\lambda^2}{g^2} \sin^2 2\beta \right)$$

- tree level lightest Higgs mass limit relaxed

- Tree level cubic term of scalar potential

$$V = V_{\text{MSSM}} + m_S^2 |S|^2 + t_S (S + h.c.) + \\ a_\lambda (S H_1 \cdot H_2 + h.c.)$$

assists a strongly 1st order EW phase transition

Baryons in the nMSSM

- Measured $\eta_B \leftrightarrow$ strongly 1st order EWPT \leftrightarrow large OP

$$\phi_c/T_c \gtrsim 1$$

- strength of EWPT \leftrightarrow minimum of finite T eff. potential

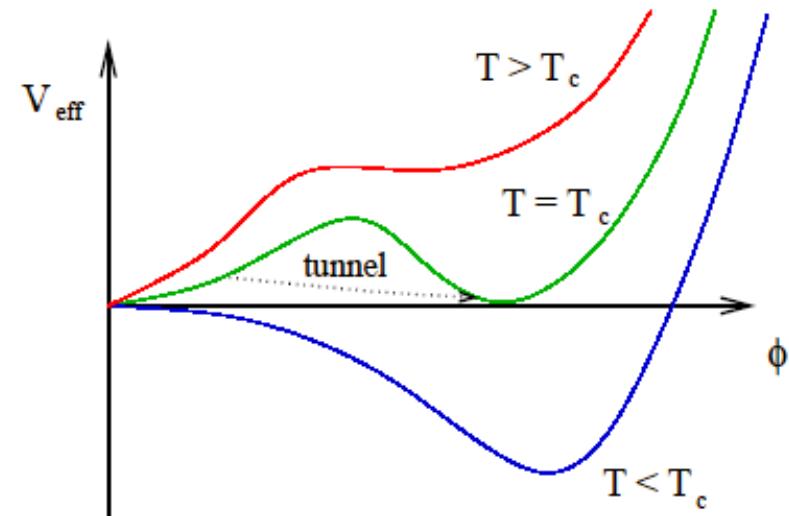
$$V_{\text{eff}}(\phi, T) = (-\mu^2 + \alpha T^2)\phi^2 - \gamma T \phi^3 + \frac{\lambda}{4} \phi^4 + \dots$$

- V_{eff} minimal for $0 < \phi$ if
 $\frac{\phi_c}{T_c} \sim \frac{\gamma}{\lambda} \rightarrow \gamma$ determines OP
- γ generated by

SM: bosonic loops $\rightarrow \gamma \sim g^3$

MSSM: sc. loops $\rightarrow \gamma \sim y^3$

nMSSM: tree level $\rightarrow \gamma \sim a_\chi$



- MSSM: light stop induces strongly 1st order EWPT

nMSSM: no need for light stop

Dark matter in nMSSM

— Neutralinos

$$M_{\tilde{Z}} = \begin{pmatrix} M_1 & & & & & \\ & \ddots & & & & \\ 0 & & M_2 & & & \\ -c_\beta s_w m_Z & c_\beta s_w m_Z & 0 & & & \\ s_\beta c_w m_Z & -s_\beta c_w m_Z & \lambda v_s & 0 & & \\ 0 & 0 & \lambda v_2 & \lambda v_1 & 0 & \end{pmatrix}$$

- unification assumption: $M_2 = \alpha_2/\alpha_1 M_1$
- EWBG: low $\tan\beta$
- typical lightest neutralino: mostly singlino

$$m_{\tilde{Z}_1} \sim 2\lambda v_1 v_2 v_s / (v_1^2 + v_2^2 + v_s^2) \lesssim 60 \text{ GeV}$$

- complex phase: $\lambda = |\lambda| e^{i\phi_\lambda}$

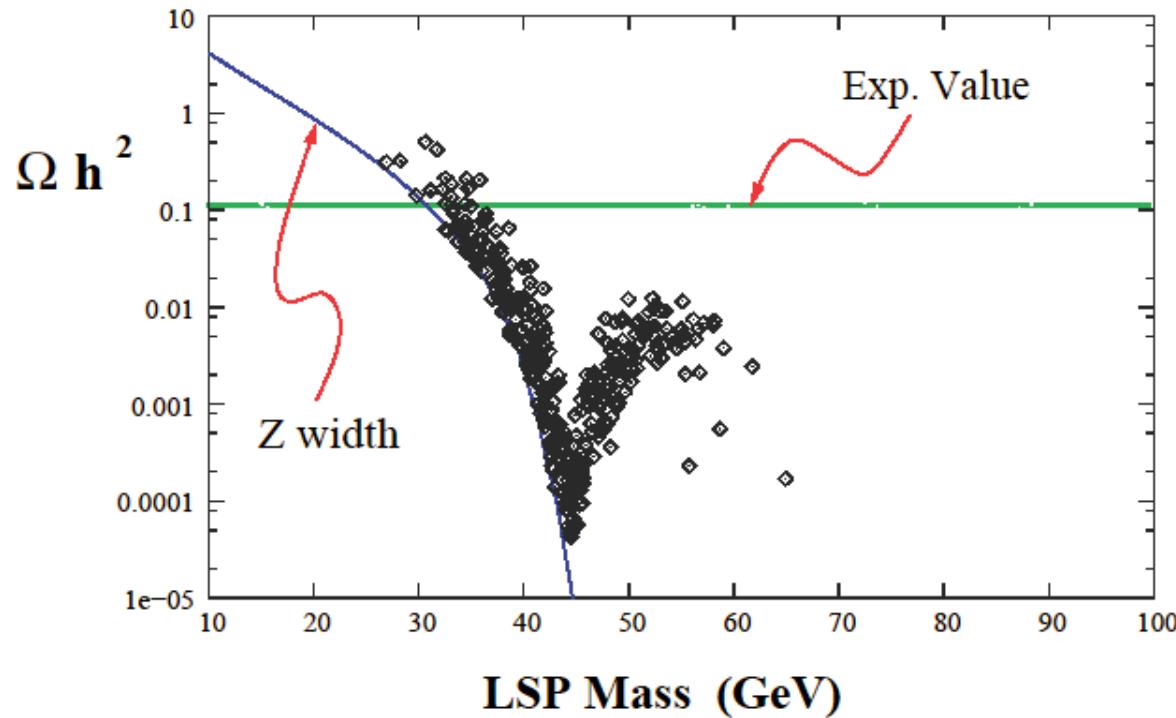
All matter in nMSSM

— Neutralino relic density

• \tilde{Z}_1 light →

no coannihilations

dominant annihilation channel: $\tilde{Z}_1 \tilde{Z}_1 \rightarrow Z \rightarrow 2X$



Menon, Morrissey, Wagner 2004

Aside: Higgs sector

$$H_1 = \begin{pmatrix} v_1 + (\phi_1 + i a_1)/\sqrt{2} \\ \phi_1^- \end{pmatrix} \quad H_2 = \begin{pmatrix} \phi_2^+ \\ v_2 + (\phi_2 + i a_2)/\sqrt{2} \end{pmatrix}$$

$$S = v_S + (\phi_S + i a_S)/\sqrt{2}$$

- $6 - 1 = 5$ neutral physical Higgses

$$\begin{pmatrix} S_1 \\ S_2 \\ S_3 \end{pmatrix} = O^S \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_S \end{pmatrix} \quad \begin{pmatrix} P_1 \\ P_2 \end{pmatrix} = O^P \begin{pmatrix} A^0 \\ a_S \end{pmatrix}$$

- Lightest Higgs (~ 115 GeV) decay modes

$$\text{Br}(S_1 \rightarrow b\bar{b}) = 8\% \quad \text{Br}(S_1 \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0) = 91\%$$

- LHC: detection via WBF & Zh production

mass determination via ratio of WBF & Zh prod

Benchmark scans

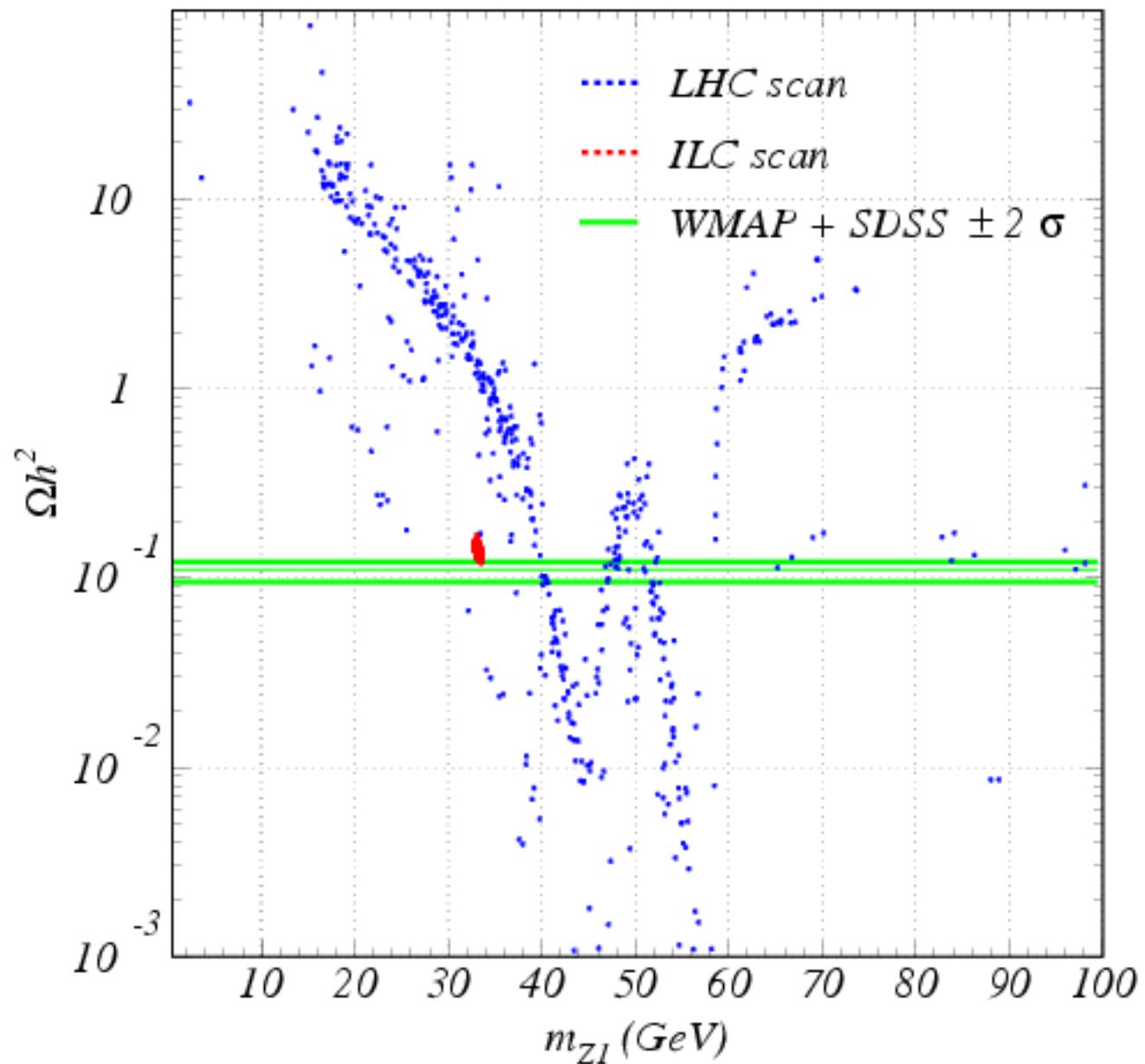
— Benchmarks

| | $\tan\beta$ | λ | v_s | a_λ | m_a | M_2 | ϕ_λ |
|---|-------------|-----------|-------|-------------|-------|-------|----------------|
| A | 1.70 | 0.619 | -384 | 373 | 923 | 245 | 0.14 |
| B | 1.99 | 0.676 | -220 | 305 | 914 | 418 | 2.57 |
| C | 1.10 | 0.920 | -276 | 386 | 514 | 462 | 2.38 |
| | | | GeV | GeV | GeV | GeV | |

— "Scan"

- generate LHC & ILC events (tree & parton level w/ BGs, jet broadening, ...)
- construct appropriate invariant mass distributions
- reconstruct masses (couplings) from distributions
- determine central values and precision

Need for the ILC



Need for the ILC

- Typical production/decay chain:

squarks/gluinos →

charginos/neutralinos →

leptons/jets

- typical invariant mass spectra (lumi 300 fb^{-1})

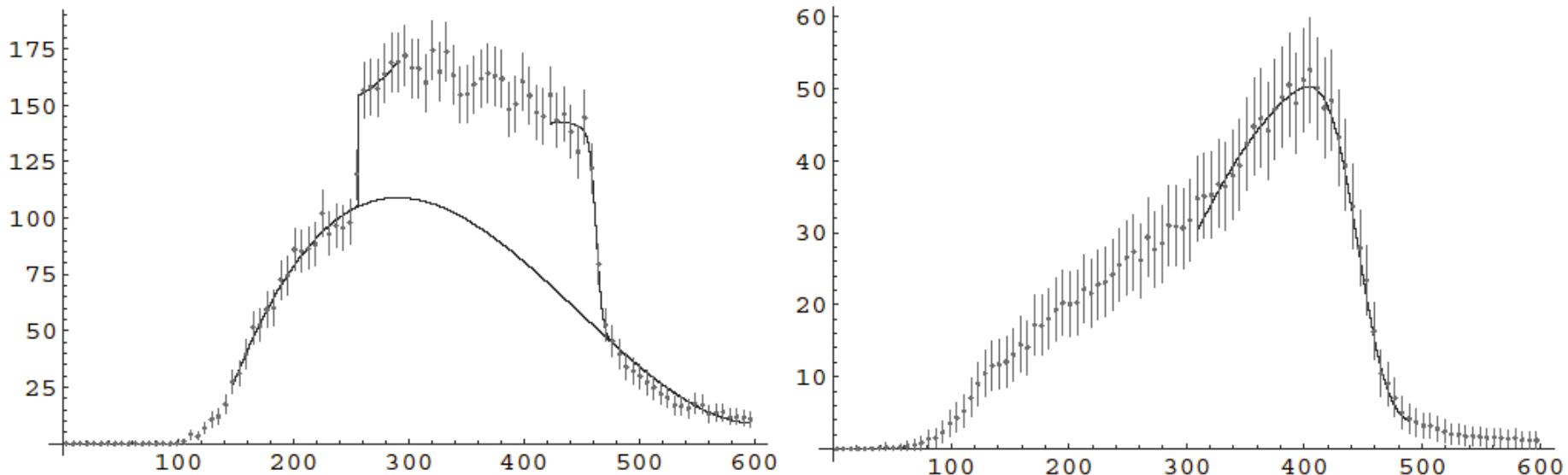


Figure 2: Fits to m_{jll} distribution for $\tilde{\chi}_3^0$ and $\tilde{\chi}_2^0$ production at LHC.

Need for the ILC

- four kinematic (mass) edges \rightarrow four mass parameters

$$m_{\text{JL,max}} = m_{\tilde{Z}_2} - m_{\tilde{Z}_1} = 73.5 \pm 0.6 \text{ GeV}$$

$$m_{\text{JL,max,2}}^2 = (m_{\tilde{Z}_2}^2 - m_{\tilde{Z}_1}^2)(m_{\tilde{b}_1}^2 - m_{\tilde{Z}_2}^2)/m_{\tilde{Z}_2}^2 = 447.0 \pm 20.0 \text{ GeV}$$

$$m_{\text{JL,min,3}}^2 = f(m_{\tilde{Z}_1}, m_{\tilde{Z}_3}, m_{\tilde{Z}_2}, m_{\tilde{b}_1}) = 256.2 \pm 7.0 \text{ GeV}$$

$$m_{\text{JL,max,3}}^2 = f(m_{\tilde{Z}_1}, m_{\tilde{Z}_3}, m_{\tilde{Z}_2}, m_{\tilde{b}_1}) = 463.5 \pm 9.0 \text{ GeV}$$

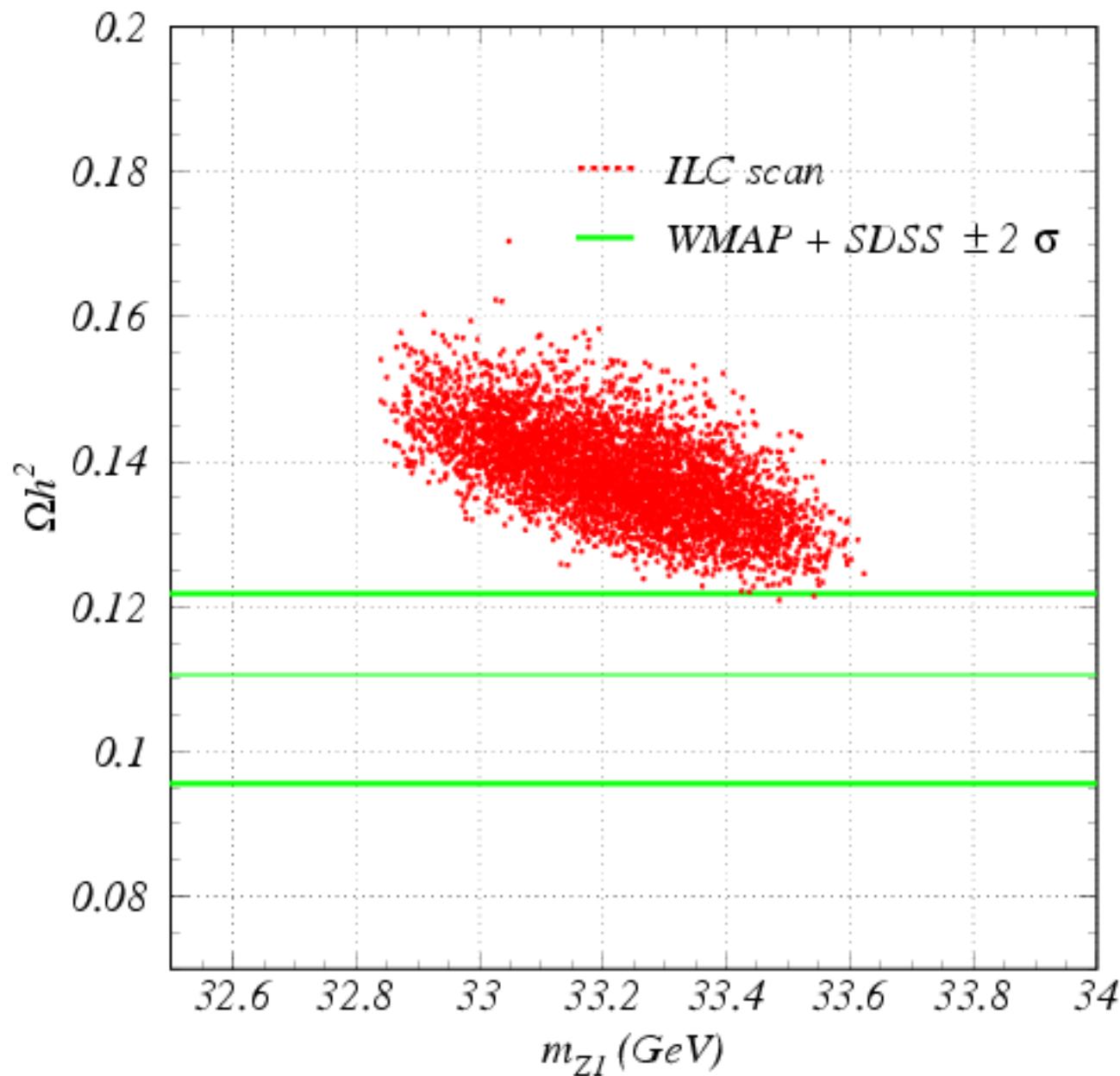
- individual masses

$$m_{\tilde{Z}_1} = 33^{+32}_{-18} \text{ GeV} \quad m_{\tilde{Z}_2} = 107^{+33}_{-18} \text{ GeV}$$

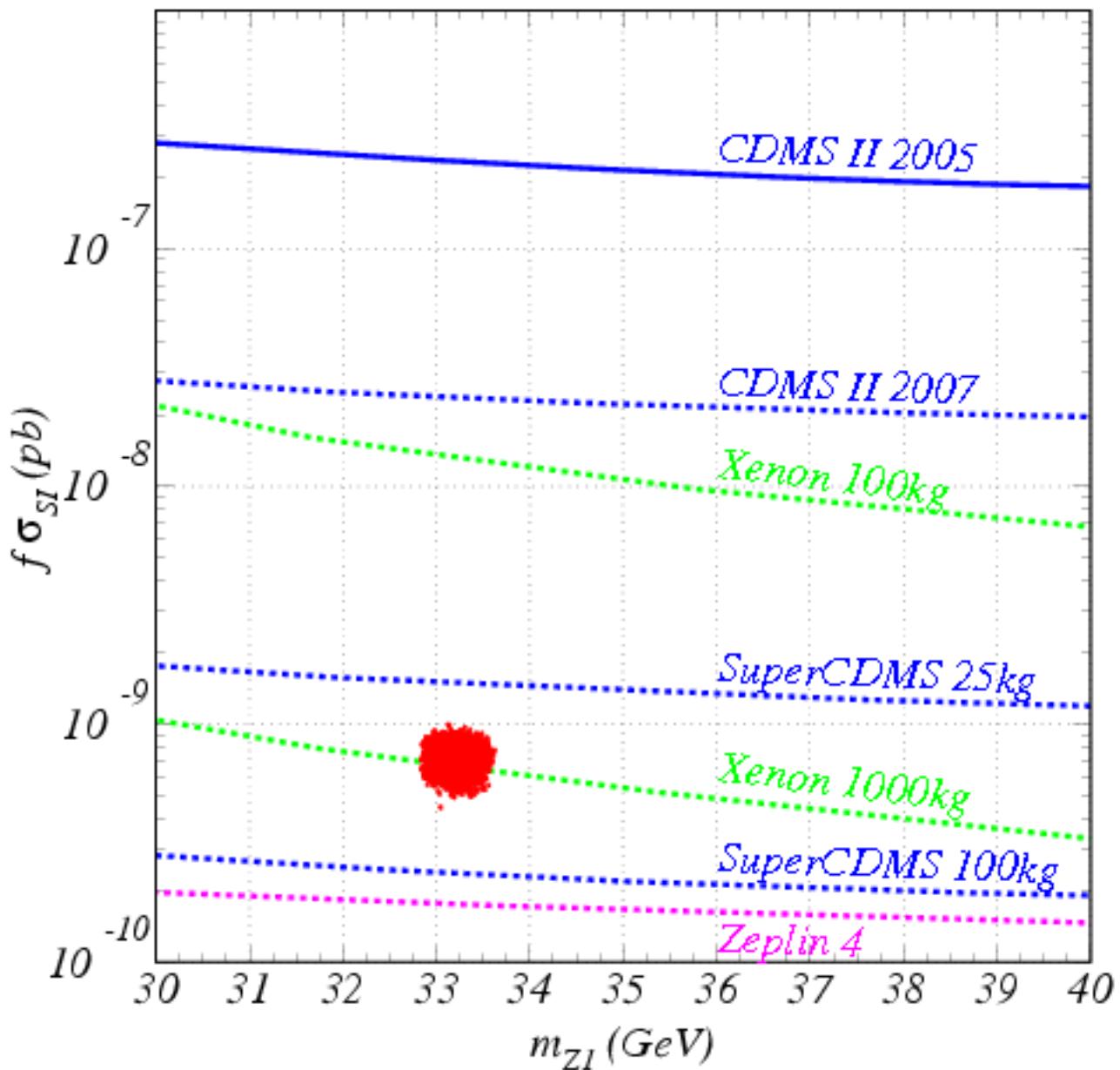
$$m_{\tilde{Z}_3} = 181^{+20}_{-10} \text{ GeV} \quad m_{\tilde{b}_1} = 499^{+30}_{-17} \text{ GeV}$$

- absolute precision is reasonable but \tilde{Z}_1 is very light!

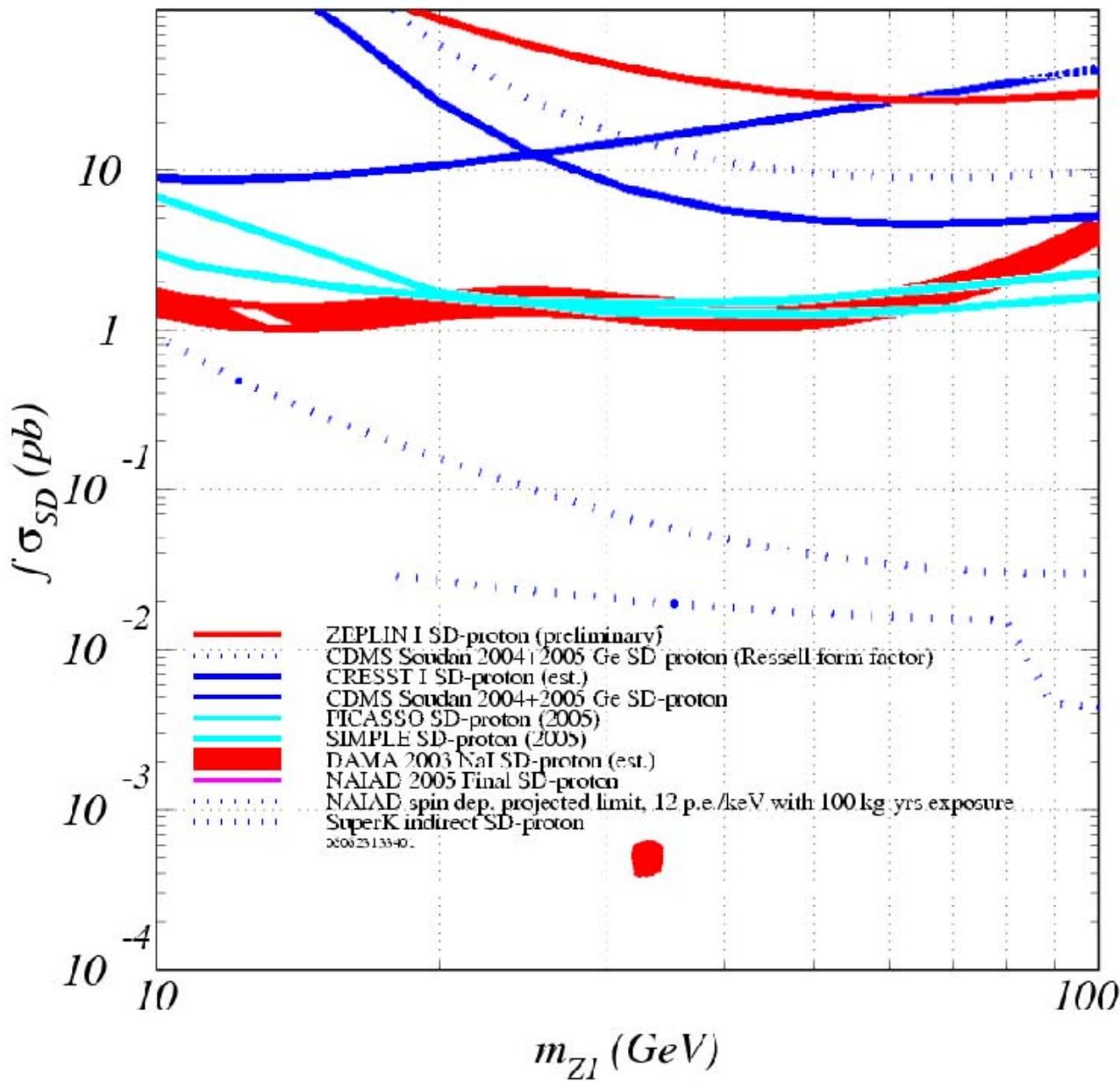
Need for the ILC



Direct detection



(in)Direct detection



Conclusions

nMSSM improves on MSSM: no μ problem, no (less) fine-tuning, less constrained spectrum

Electroweak baryogenesis less constrained in nMSSM due to tree level cubic contribution to scalar potential

All matter in the Universe can be simultaneously generated in the nMSSM

ILC precision is critical for determining astrophysical parameters: relic density, WIMP-nucleon scattering, ...