

Precise Prediction of M_W in the MSSM

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based on collaboration with
W. Hollik, D. Stöckinger, A.M. Weber and G. Weiglein

1. Motivation
2. Calculation of M_W in the MSSM
3. Numerical results
4. Remaining theoretical uncertainties
5. Conclusions

1. Motivation

M_W is a precision observable

Experimental situation:

$$\begin{aligned} \text{today: LEP2, Tevatron} &\Rightarrow \delta M_W^{\text{exp}} = 32 \text{ MeV} \\ \text{Tevatron (8 if}^{-1}\text{)} &\Rightarrow \delta M_W^{\text{exp}} = 20 \text{ MeV} \\ \text{LHC} &\Rightarrow \delta M_W^{\text{exp}} = 15 \text{ MeV} \\ \text{ILC} &\Rightarrow \delta M_W^{\text{exp}} = 10 \text{ MeV} \\ \text{GigaZ} &\Rightarrow \delta M_W^{\text{exp}} = 7 \text{ MeV} \end{aligned}$$

Prediction in the SM:

$$\begin{aligned} \delta M_W^{\text{theory,SM,today}} &\approx \pm 4 \text{ MeV} \\ \delta M_W^{\text{theory,SM,future}} &\approx \pm 2 \text{ MeV} \\ \delta m_t = 2.3 \text{ GeV} : \quad \delta M_W^{m_t, \text{para, today}} &\approx \pm 14 \text{ MeV} \\ \delta m_t = 0.1 \text{ GeV} : \quad \delta M_W^{m_t, \text{para, future}} &\approx \pm 1 \text{ MeV} \\ \delta(\Delta\alpha_{\text{had}}) : \quad \delta M_W^{\Delta\alpha, \text{para, today}} &\approx \pm 6.5 \text{ MeV} \\ \delta(\Delta\alpha_{\text{had}}) : \quad \delta M_W^{\Delta\alpha, \text{para, future}} &\approx \pm 1 \text{ MeV} \end{aligned}$$

1. Motivation

M_W is a precision observable

Experimental situation:

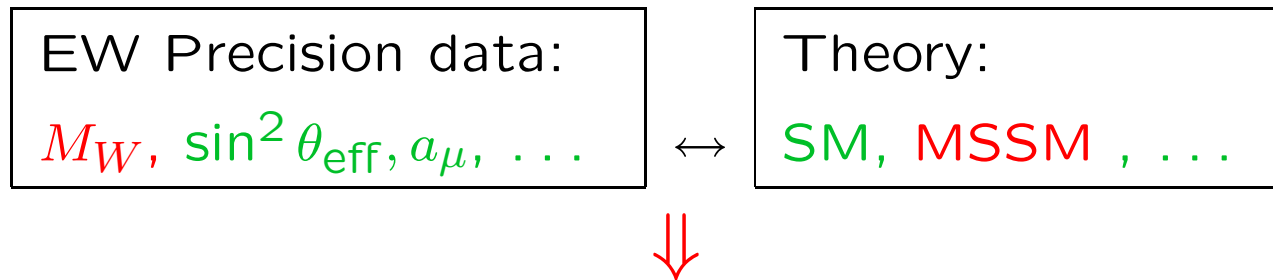
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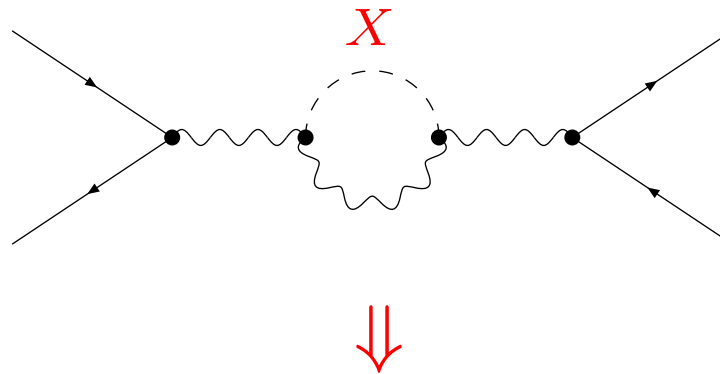
$$\begin{aligned} \delta M_W^{\text{theory,SM,today}} &\approx \pm 4 \text{ MeV} && \text{SUSY?} \\ \delta M_W^{\text{theory,SM,future}} &\approx \pm 2 \text{ MeV} \\ \delta m_t = 2.3 \text{ GeV} : \quad \delta M_W^{m_t, \text{para, today}} &\approx \pm 14 \text{ MeV} \\ \delta m_t = 0.1 \text{ GeV} : \quad \delta M_W^{m_t, \text{para, future}} &\approx \pm 1 \text{ MeV} \\ \delta(\Delta\alpha_{\text{had}}) : \quad \delta M_W^{\Delta\alpha, \text{para, today}} &\approx \pm 6.5 \text{ MeV} \\ \delta(\Delta\alpha_{\text{had}}) : \quad \delta M_W^{\Delta\alpha, \text{para, future}} &\approx \pm 1 \text{ MeV} \end{aligned}$$

Precision Observables (POs):

Comparison of electro-weak precision observables with theory:



Test of theory at quantum level: Sensitivity to loop corrections

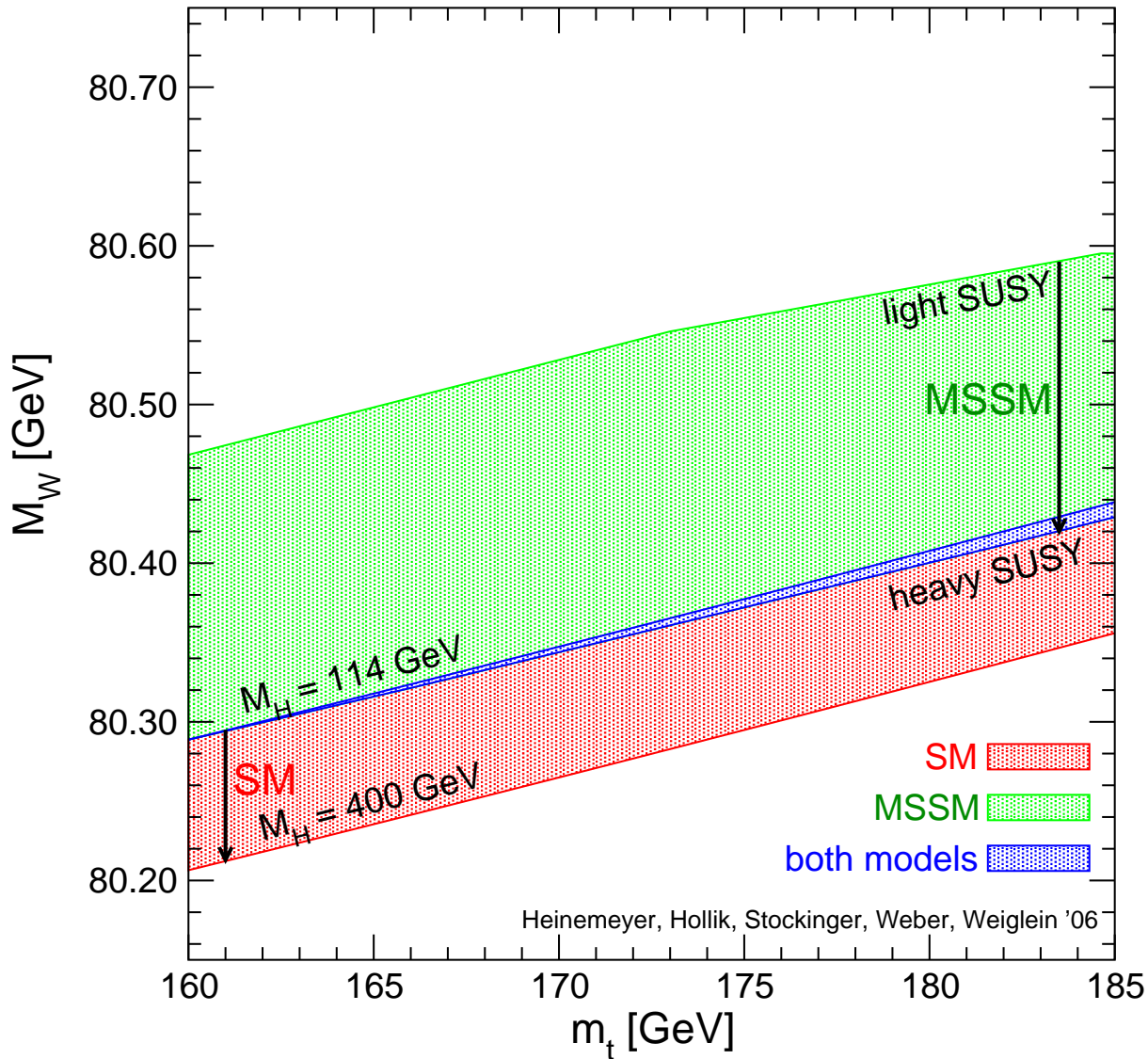


Very high accuracy of measurements and theoretical predictions needed

- Which model fits better?
- Does the prediction of a model contradict the experimental data?

Prediction for M_W in the SM and the MSSM :

[S.H., W. Hollik, D. Stockinger, A.M. Weber, G. Weiglein '06]



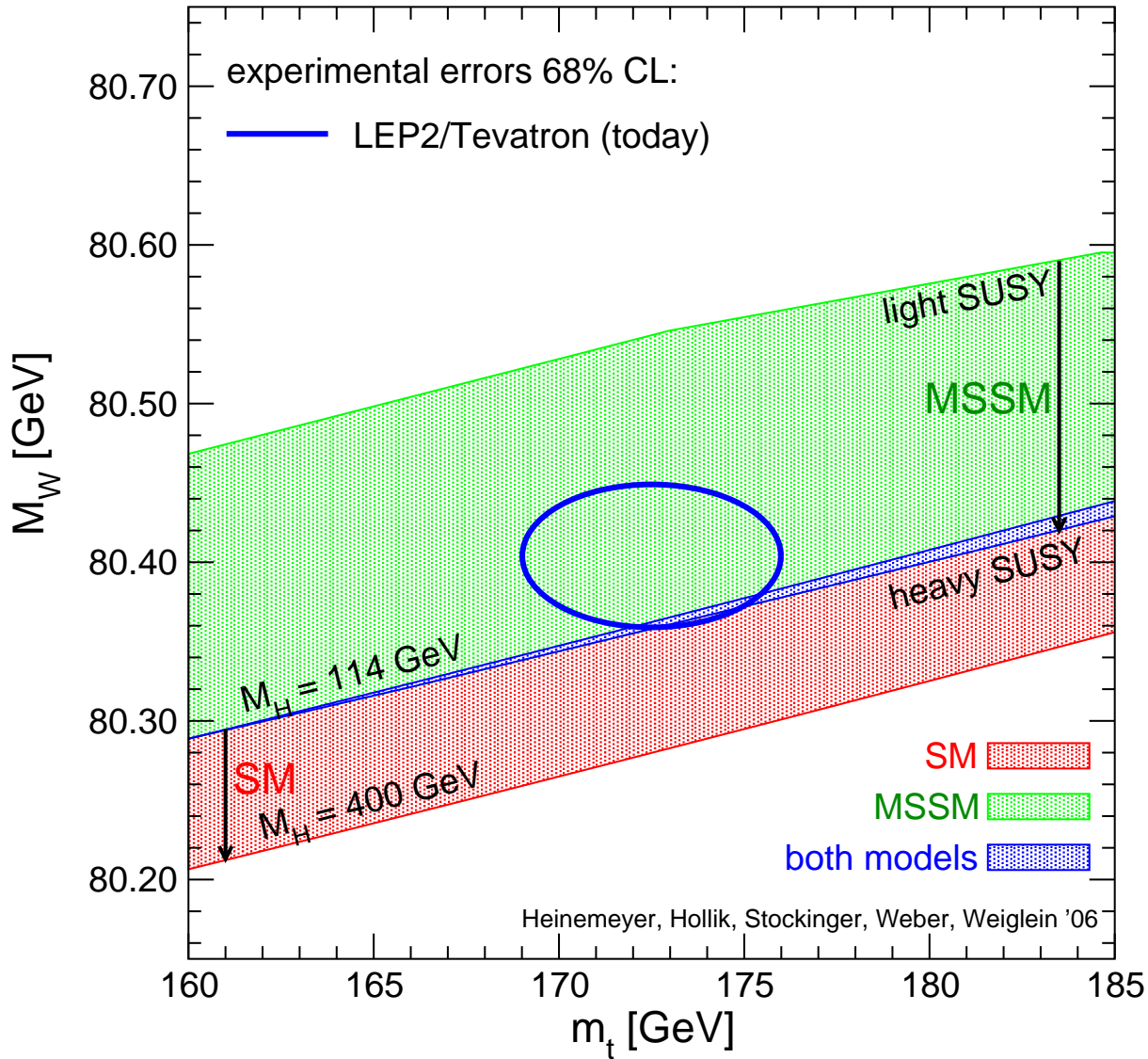
MSSM band:
scan over
SUSY masses

overlap:
SM is MSSM-like
MSSM is SM-like

SM band:
variation of M_H^{SM}

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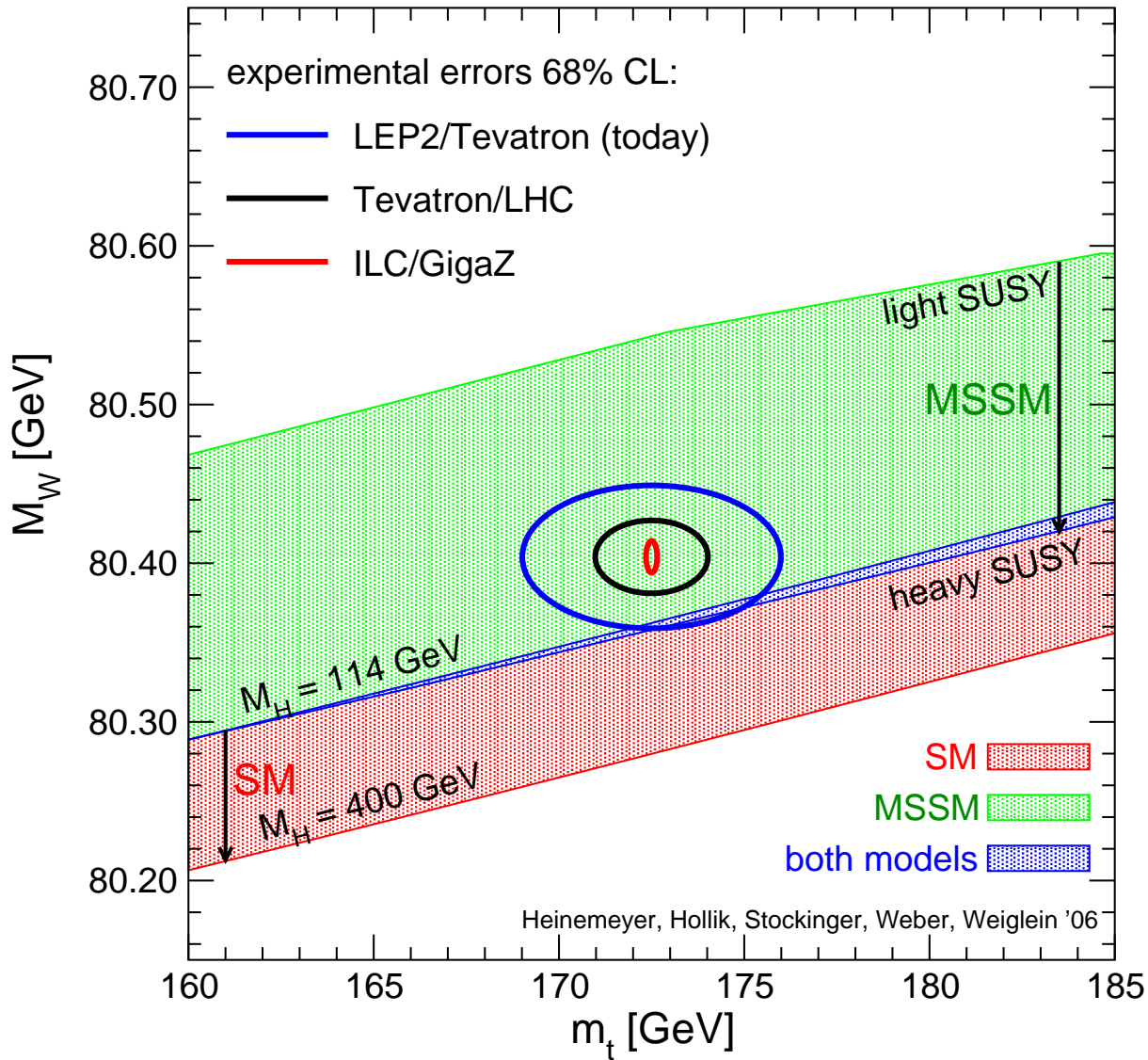
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MSSM band:

scan over
SUSY masses

overlap:

SM is MSSM-like
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SM band:

variation of M_H^{SM}

2. Calculation of M_W in the MSSM

Theoretical prediction for M_W in terms of $M_Z, \alpha, G_\mu, \Delta r$:

$$M_W^2 \left(1 - \frac{M_W^2}{M_Z^2} \right) = \frac{\pi \alpha}{\sqrt{2} G_\mu} (1 + \Delta r)$$



loop corrections

Evaluate Δr from μ decay $\Rightarrow M_W$

One-loop result for M_W in the SM:

[A. Sirlin '80] , [W. Marciano, A. Sirlin '80]

$$\begin{aligned} \Delta r_{1\text{-loop}} = & \Delta\alpha & - & \frac{c_W^2}{s_W^2} \Delta\rho & + & \Delta r_{\text{rem}}(M_H) \\ & \sim \log \frac{M_Z}{m_f} & & \sim m_t^2 & & \\ & \sim 6\% & & \sim 3.3\% & & \sim 1\% \end{aligned}$$

Combination of SM and MSSM result

- use **best available SM** result
- add **all available MSSM** corrections
- subtract double counting

⇒ **decoupling limit ok** ($M_{\text{SUSY}} \rightarrow \infty$)

... but some corrections included only via their **SM part**

⇒ **best solution**

Status of SM calculation

- Most recently full **electroweak two-loop** result available
[Awramik, Czakon, Freitas, Hollik, Onishenko, Veretin, Walter, Weiglein '03, '04]
- Most recent **electroweak three-loop** corrections via $\Delta\rho$
[Faisst, Kühn, Seidensticker, Veretin '04]
- **Compact parametric formula** used
[Awramik, Czakon, Freitas, Weiglein '04]

Status of MSSM calculations of M_W

- MSSM, Δr : full one-loop corrections
[P. Chankowski, A. Dabelstein, W. Hollik, W. Möhle, S. Pokorski, J. Rosiek '94]
[D. Garcia, J. Solà '94]
- MSSM, $\Delta\rho$: leading $\mathcal{O}(\alpha\alpha_s)$ corrections
[A. Djouadi, P. Gambino, S.H., W. Hollik, C. Jünger, G. Weiglein '97]
- MSSM, Δr : leading gluonic $\mathcal{O}(\alpha\alpha_s)$ corr.
[S.H. '98] [S.H., W. Hollik, G. Weiglein '04]
- MSSM, $\Delta\rho$: leading $\mathcal{O}(\alpha_t^2, \alpha_t\alpha_b, \alpha_b^2)$ corrections
[S.H., G. Weiglein '01, '03]
[J. Haestier, S.H., D. Stöckinger, G. Weiglein '05]

Most recent:

- MSSM, Δr : one-loop with complex parameters
[S.H., W. Hollik, D. Stöckinger, A.M. Weber, G. Weiglein '06]

Most dominant:

- $\Delta\rho$: one- or two-loop, contribution from \tilde{t}/\tilde{b} sector

More details about our calculation:

- Complex phases in the squark sector enter **only** via **shift in squark masses** (explicit dependence drops out)

$$\begin{aligned}|X_t|^2 &= |A_t|^2 + |\mu \cot \beta|^2 - 2|A_t| \cdot |\mu| \cot \beta \cos(\phi_{A_t} + \phi_\mu) \\ |X_b|^2 &= |A_b|^2 + |\mu \tan \beta|^2 - 2|A_b| \cdot |\mu| \tan \beta \cos(\phi_{A_b} + \phi_\mu)\end{aligned}$$

Only **some phase combinations are physical**, other phases can be rotated away.

Examples for physical combinations:

$$\begin{aligned}\phi_{A_t} + \phi_\mu \\ \phi_{A_b} + \phi_\mu\end{aligned}$$

- **Higgs mass** dependence of the **two-loop** contributions is known to be very strong
⇒ we use **FeynHiggs** (www.feynhiggs.de)
- All **one-loop** calculations have been performed with **FeynArts** and **FormCalc**
[T. Hahn et al '00 - '05]

Treatment of the phase dependence beyond one-loop order

Phase dependence at the two-loop level approximated by a simple interpolation based on:

full phase dependence at the one-loop level, $M_W^{1L}(\phi)$,

two-loop results for real parameters, $M_W^{\text{full}}(0)$, $M_W^{\text{full}}(\pi)$

⇒ Two-loop result for complex phase ϕ :

$$M_W^{\text{full}}(\phi) = M_W^{1L}(\phi) + \left[M_W^{\text{full}}(0) - M_W^{1L}(0) \right] \times \frac{1 + \cos \phi}{2} \\ + \left[M_W^{\text{full}}(\pi) - M_W^{1L}(\pi) \right] \times \frac{1 - \cos \phi}{2}$$

3. Numerical results

3A) phase dependence from squark sector

Complex phases in the squark sector enter **only** via **shift in squark masses** (explicit dependence drops out)
⇒ phase dependence must be **reflected in the squark masses**

Only **some phase combinations are physical**, other phases can be rotated away.

Examples for physical combinations:

$$\phi_{A_t} + \phi_\mu$$

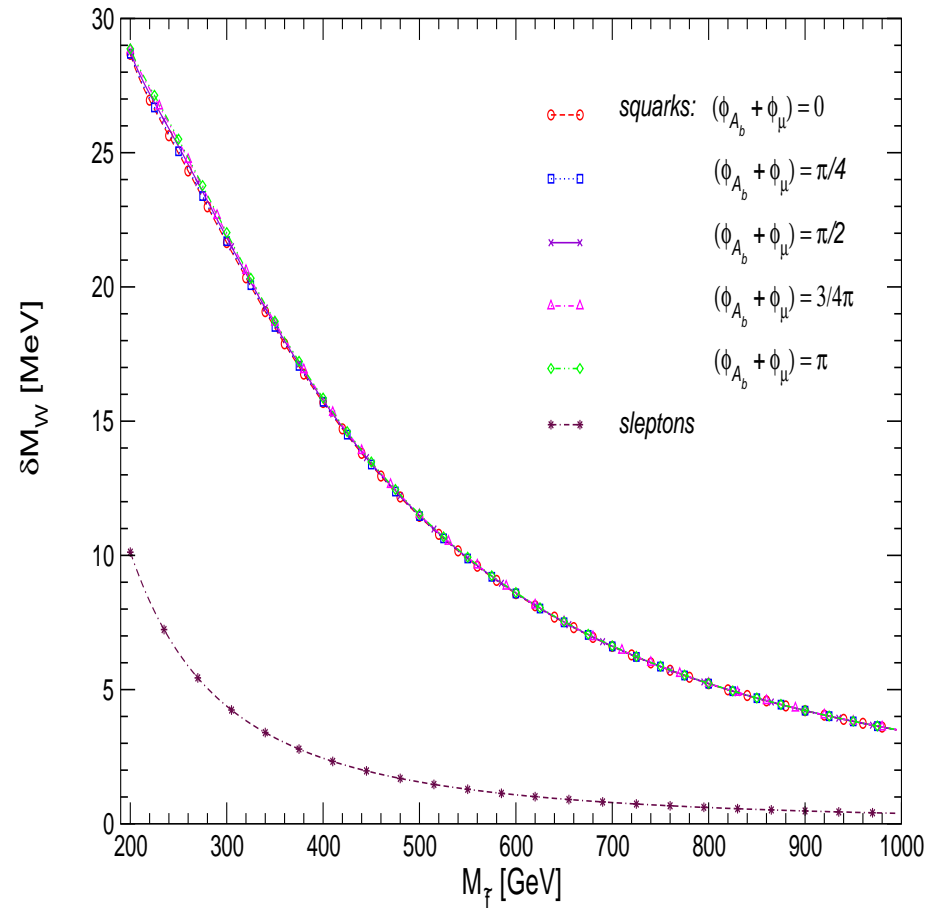
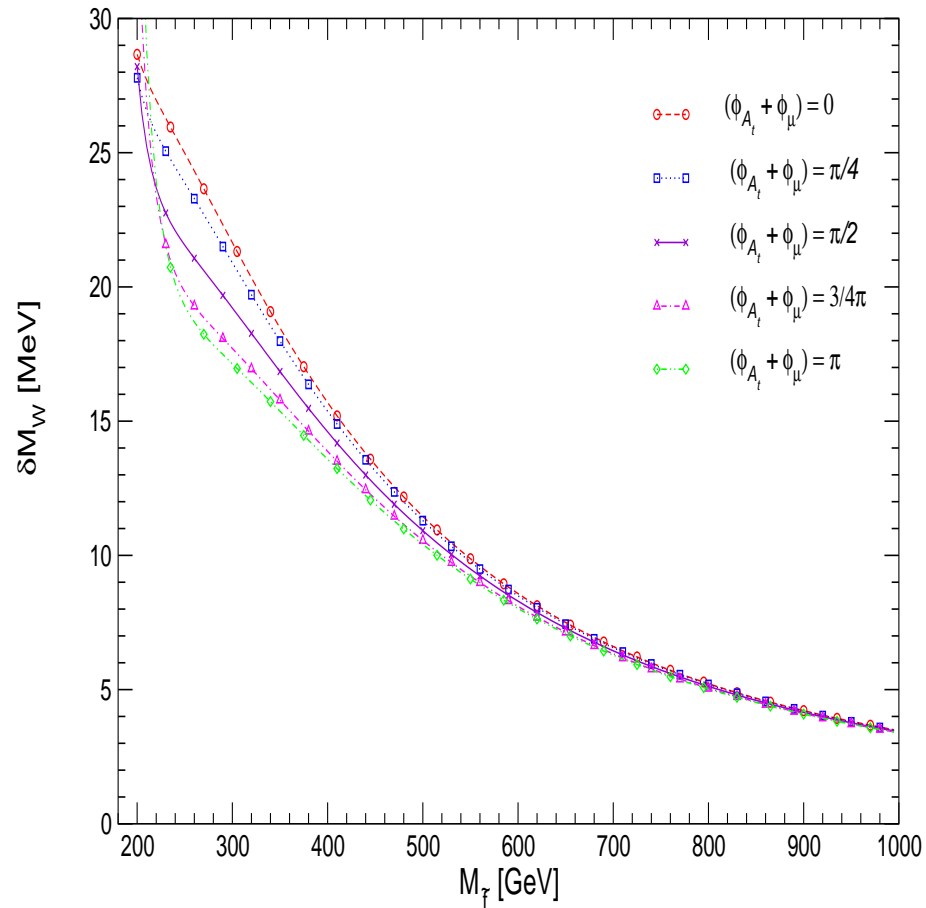
$$\phi_{A_b} + \phi_\mu$$

δM_W evaluated from

$$\delta M_W = -\frac{M_W}{2} \frac{s_W^2}{c_W^2 - s_W^2} \Delta r$$

δM_W dependence on ϕ_{A_t} and ϕ_{A_b} (I):

($|A_{t,b}| = 350$ GeV, $\mu = 300$ GeV, $\tan \beta = 10$)

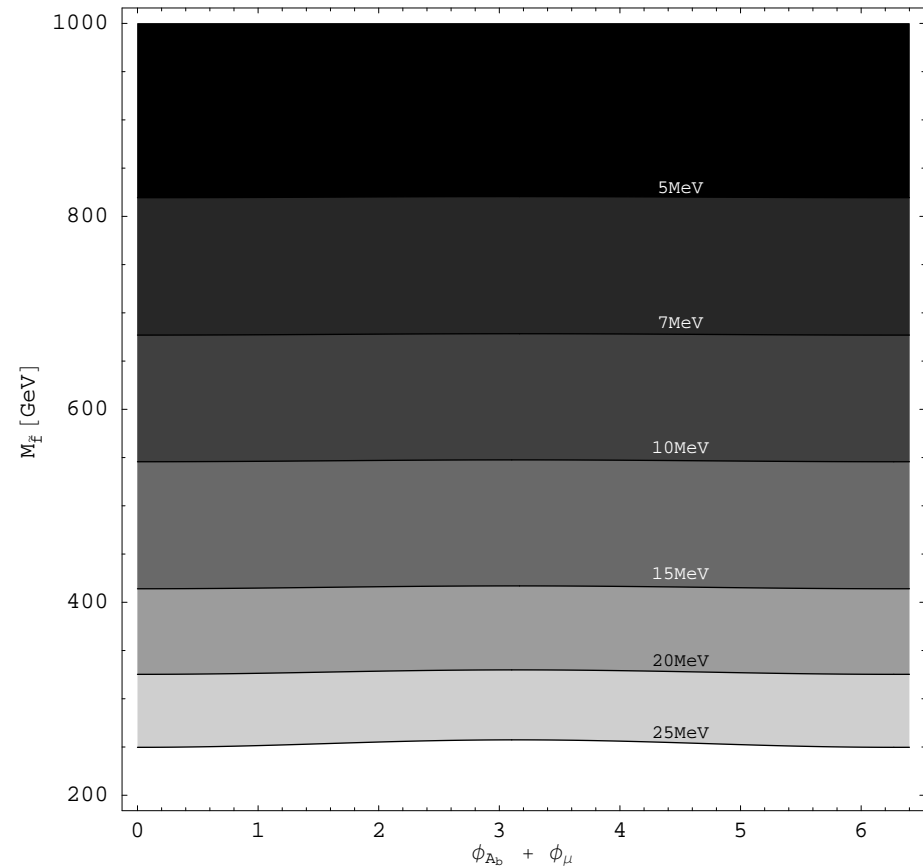
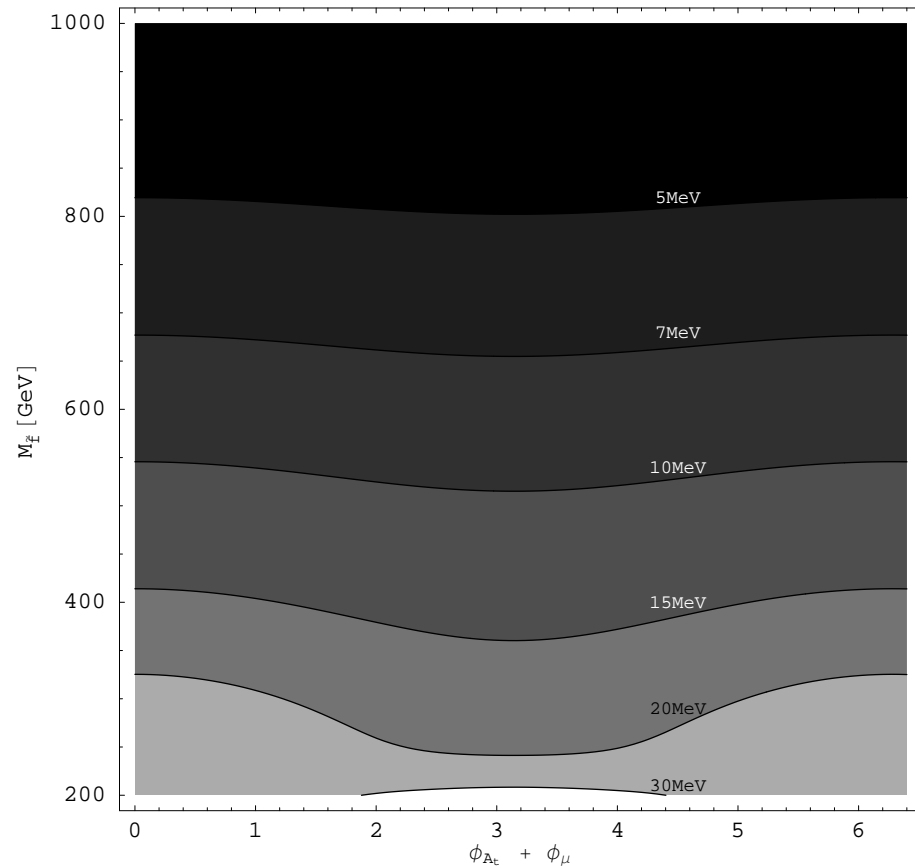


\Rightarrow large squark contribution, decoupling with M_{SUSY}

$\Rightarrow \phi_{A_t}$ dependence, but no ϕ_{A_b} effects

δM_W dependence on ϕ_{A_t} and ϕ_{A_b} (II):

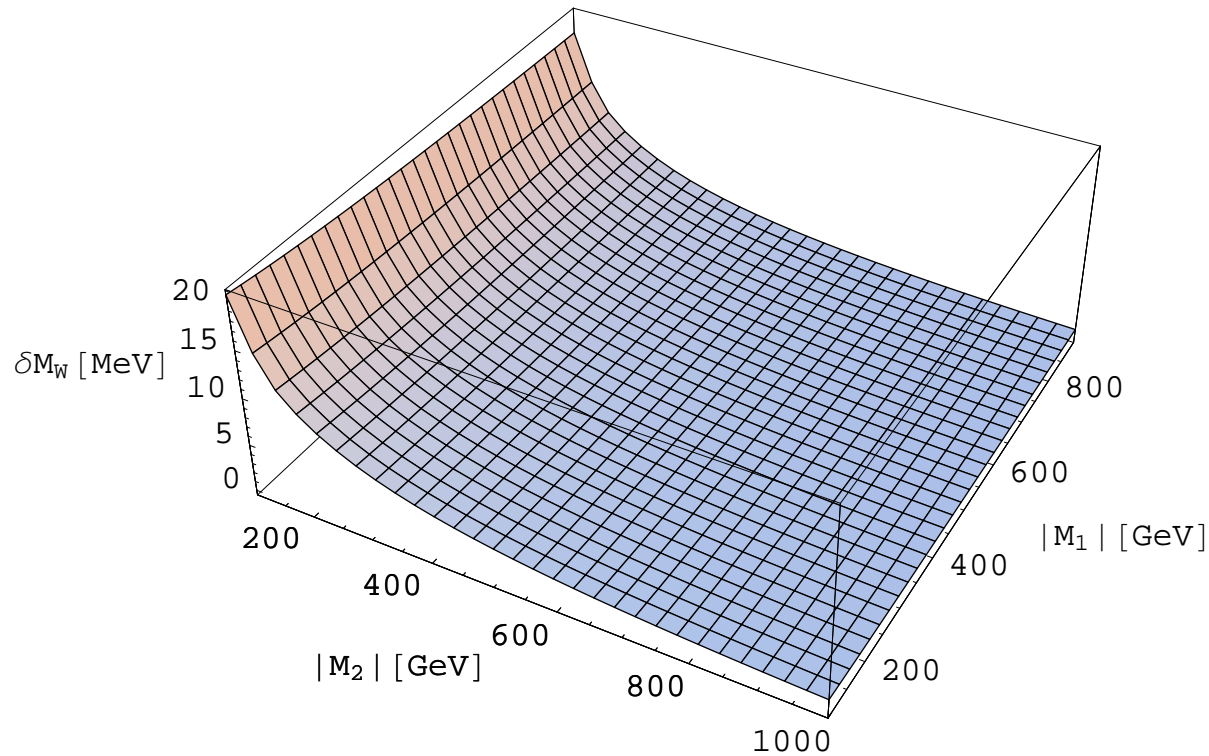
($|A_{t,b}| = 350$ GeV, $\mu = 300$ GeV, $\tan \beta = 10$)



\Rightarrow large squark contribution, decoupling with M_{SUSY}

$\Rightarrow \phi_{A_t}$ dependence, but no ϕ_{A_b} effects

3B) dependence on chargino/neutralino sector

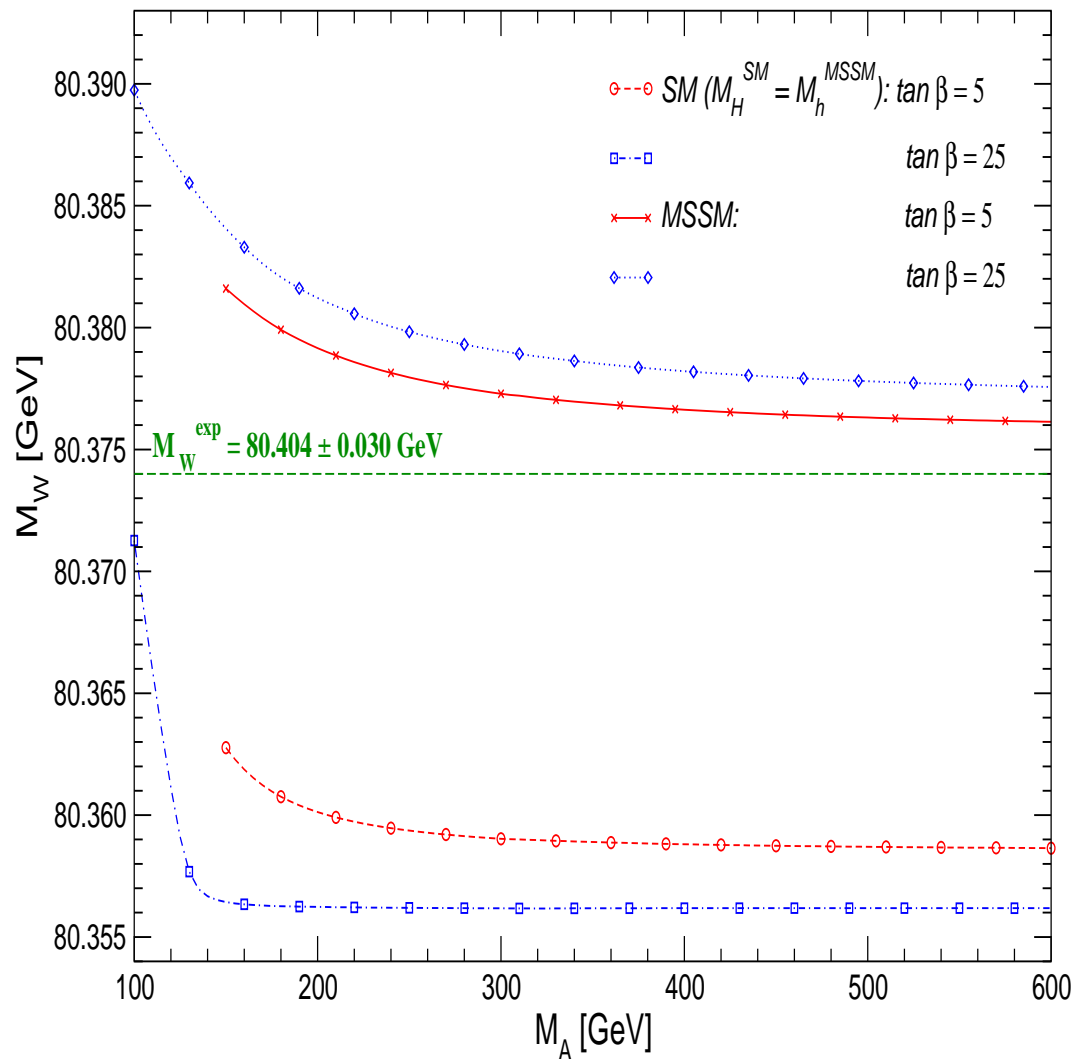


- hardly any M_1 dependence
- up to 20 MeV contribution via M_2

other parameters:
 $M_{\text{SUSY}} = 250$ GeV
 $\mu = 300$ GeV,
 $\phi_{M_{1,2}} = 0$
 $\tan \beta = 10$

3C) Prediction of M_W

Comparison with SM:



other parameters:

$$M_{\text{SUSY}} = 600 \text{ GeV}$$

$$A_{t,b} = 1200 \text{ GeV}$$

$$\mu = M_2 = m_{\tilde{g}} = 300 \text{ GeV},$$

$$\phi_x = 0$$

$$M_W^{\text{SUSY}} - M_W^{\text{SM}} \gtrsim 20 \text{ MeV}$$

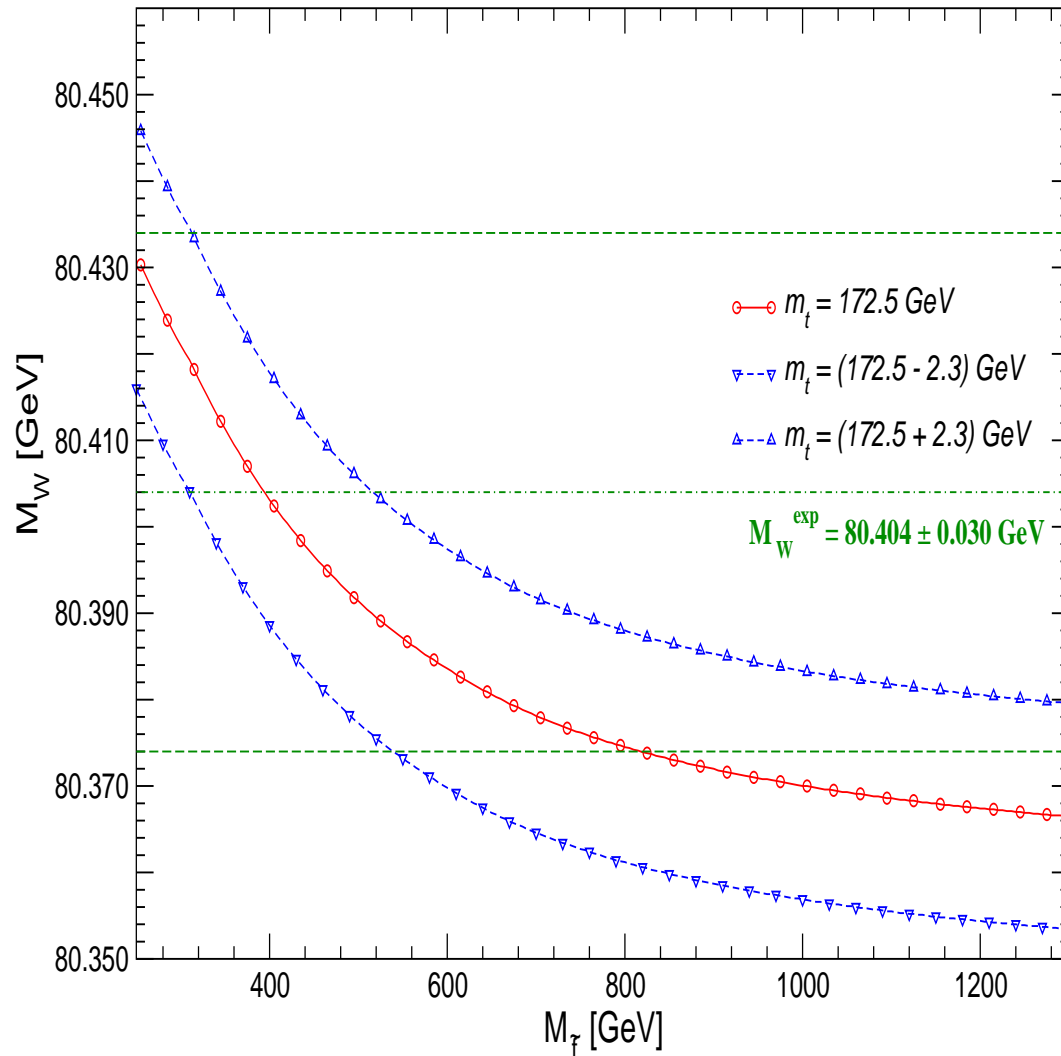
(typical difference)

Comparison with M_W^{exp} :

MSSM agrees at 1σ

SM agrees at 2σ

Prediction of M_W : variation of m_t



other parameters:

$$A_{t,b} = 2 M_{\text{SUSY}}$$

$$\mu = M_2 = m_{\tilde{g}} = 300 \text{ GeV},$$

$$\phi_x = 0$$

$$\tan \beta = 10$$

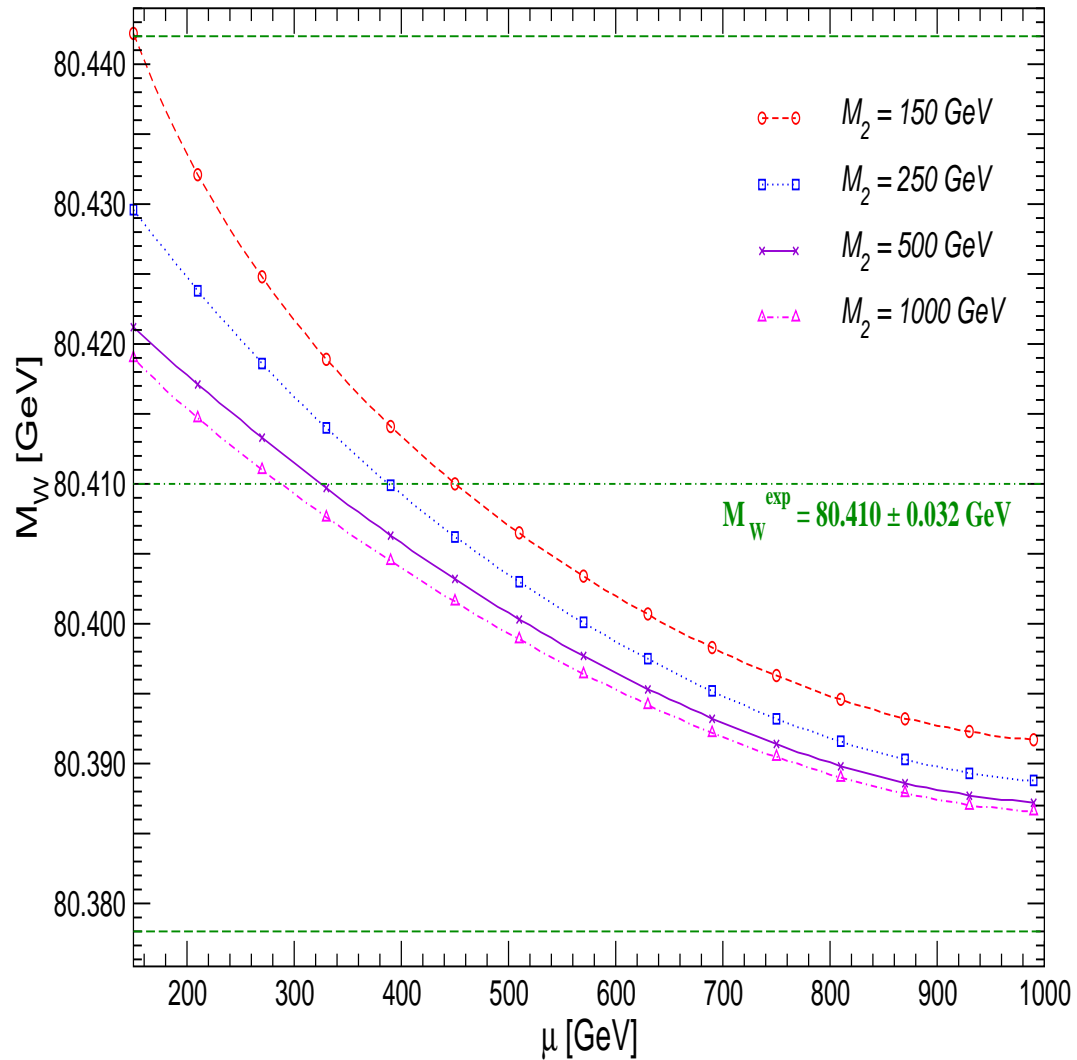
$$\delta m_t = \pm 2.3 \text{ GeV}$$

$$\Rightarrow \delta M_W \approx \pm 14 \text{ MeV}$$

lower m_t disfavored

\Rightarrow largest parametric uncertainty

Prediction of M_W : variation of μ and M_2



other parameters:

$$M_{\text{SUSY}} = 300 \text{ GeV}$$

$$A_{t,b} = 2 M_{\text{SUSY}}$$

$$m_{\tilde{g}} = 300 \text{ GeV},$$

$$\phi_x = 0$$

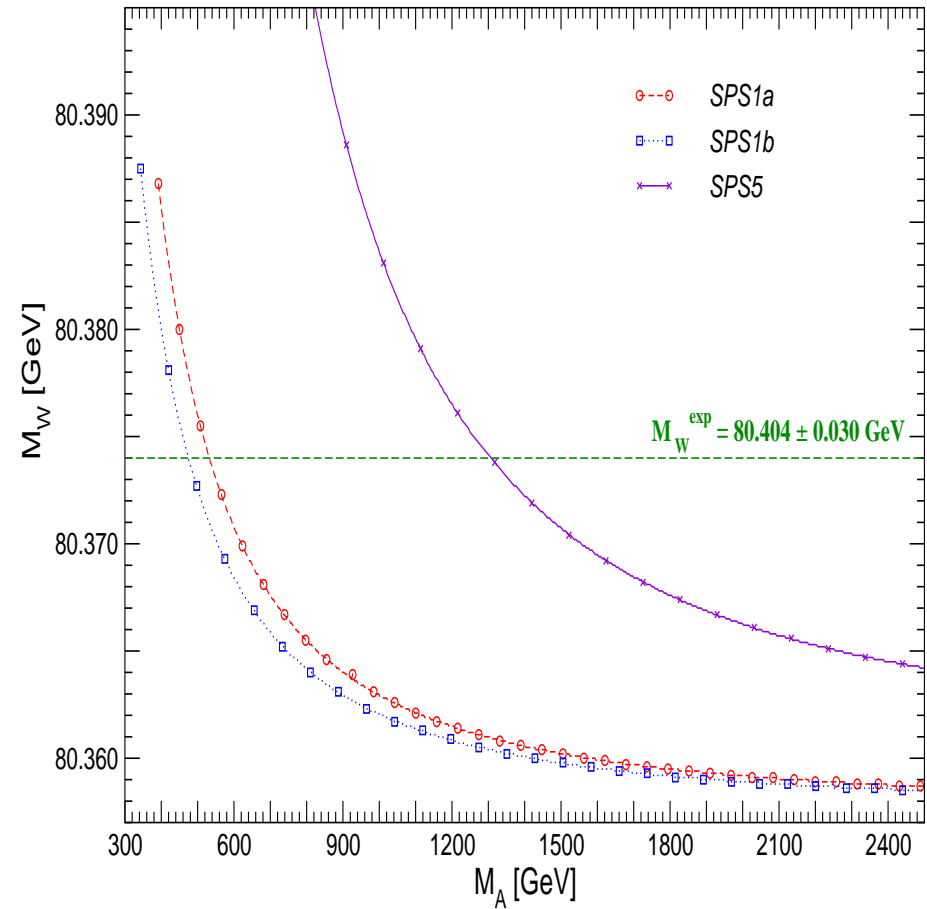
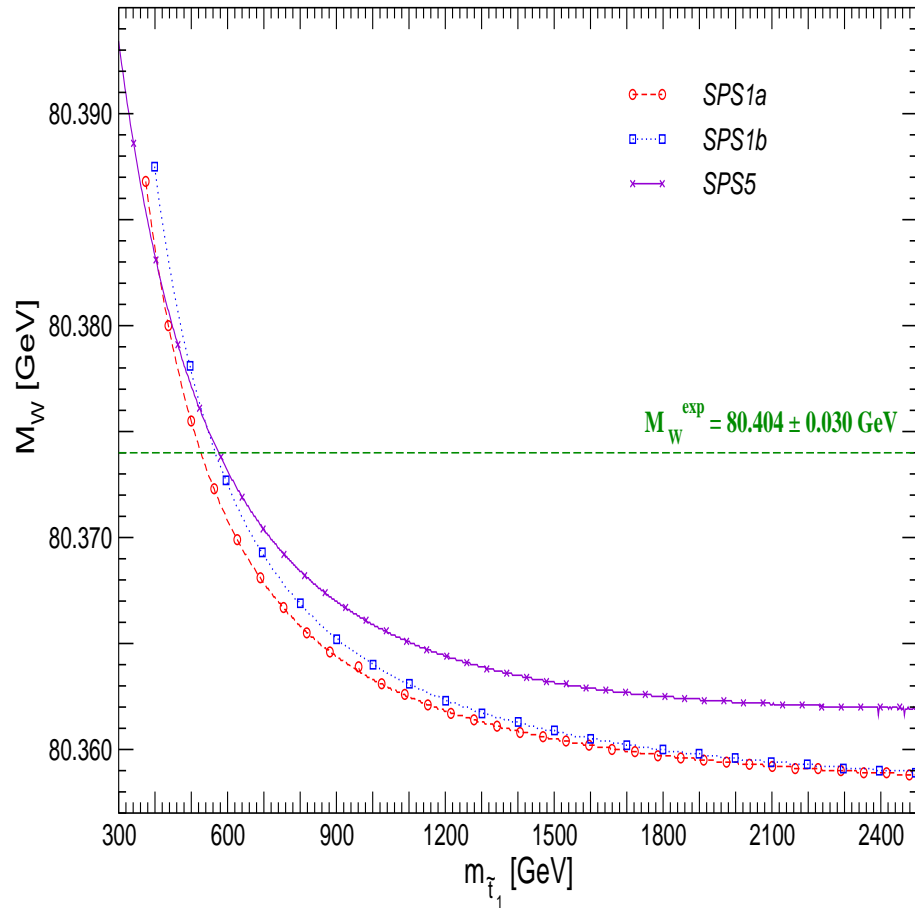
$$\tan \beta = 10$$

$$M_A = 1000 \text{ GeV}$$

\Rightarrow large variation with μ

\Rightarrow smaller variation with M_2

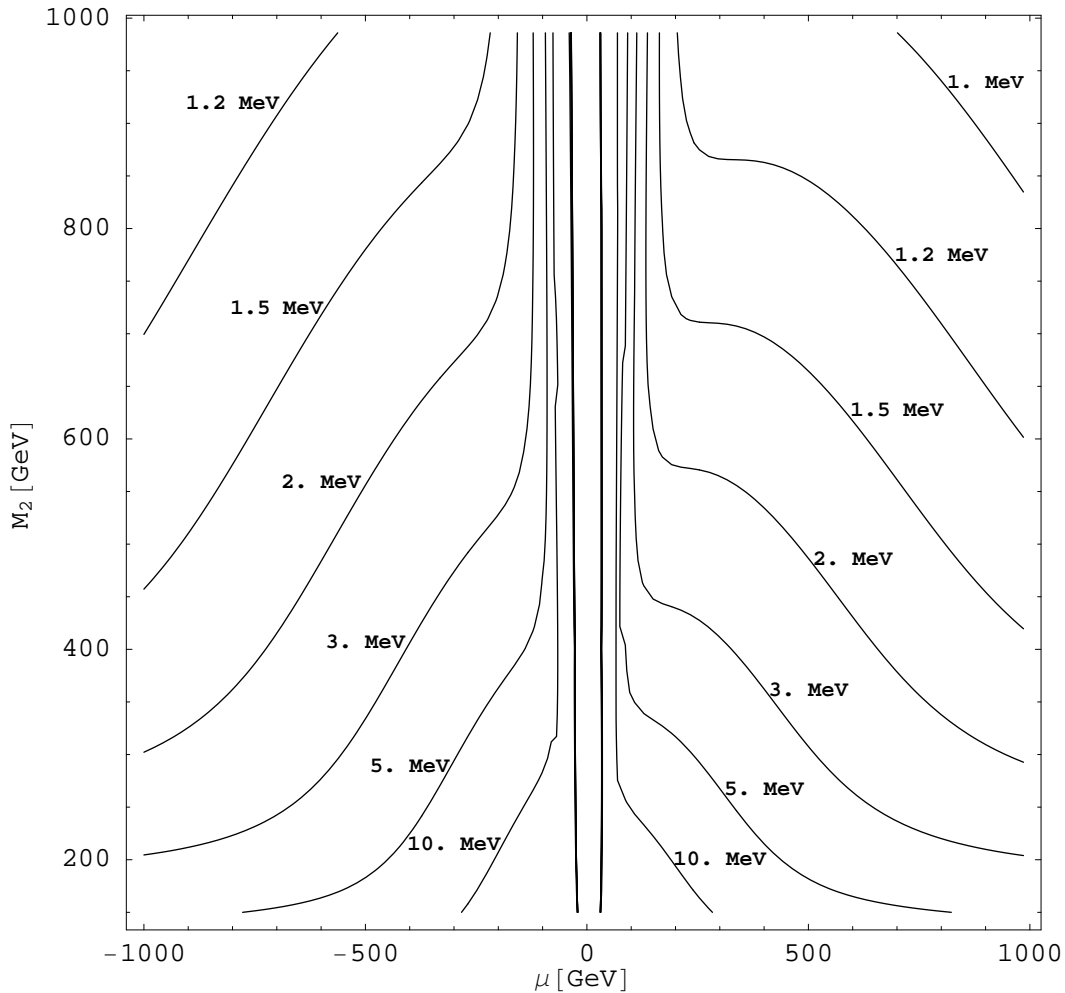
Prediction of M_W : typical scenarios: SPS 1a, 1b, 5



⇒ only “lower” masses show agreement at 1σ

Prediction of M_W : Split SUSY

[N. Arkani-Hamed, S. Dimopoulos '04] [G. Giudice, A. Romanino '04]



Difference to M_W^{SM}
with $M_H^{\text{SM}} = M_h$

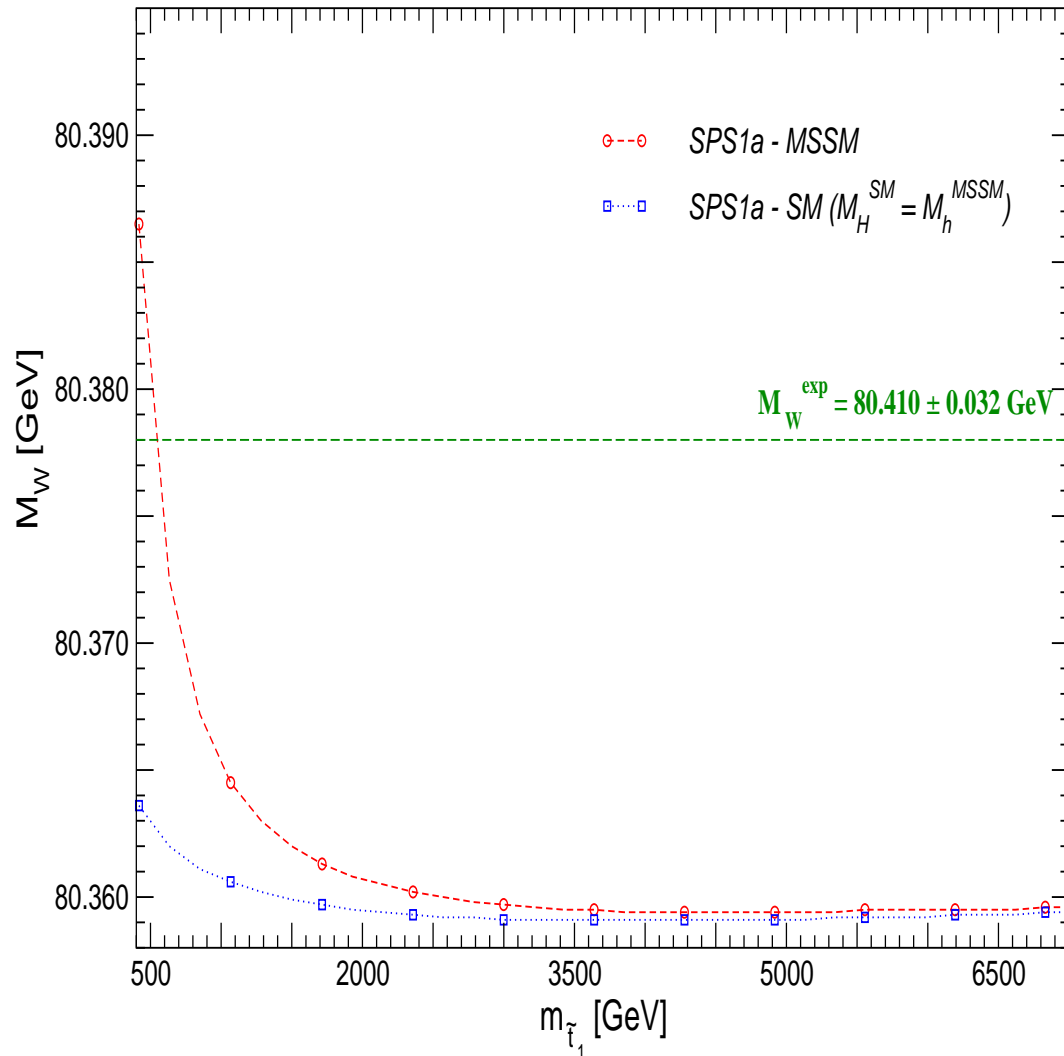
⇒ no deviation for
Tevatron, LHC, ILC
precision

⇒ GigaZ can confirm only
very light masses

(as expected ...)

Split SUSY disagrees with
experiment at the 2σ level

Prediction of M_W : Decoupling limit



Compared to SM with

$$M_H^{\text{SM}} = M_h$$

$$M_{\text{SUSY}} = 500 \text{ GeV}$$

$$\Rightarrow \delta M_W^{\text{MSSM-SM}} > 20 \text{ MeV}$$

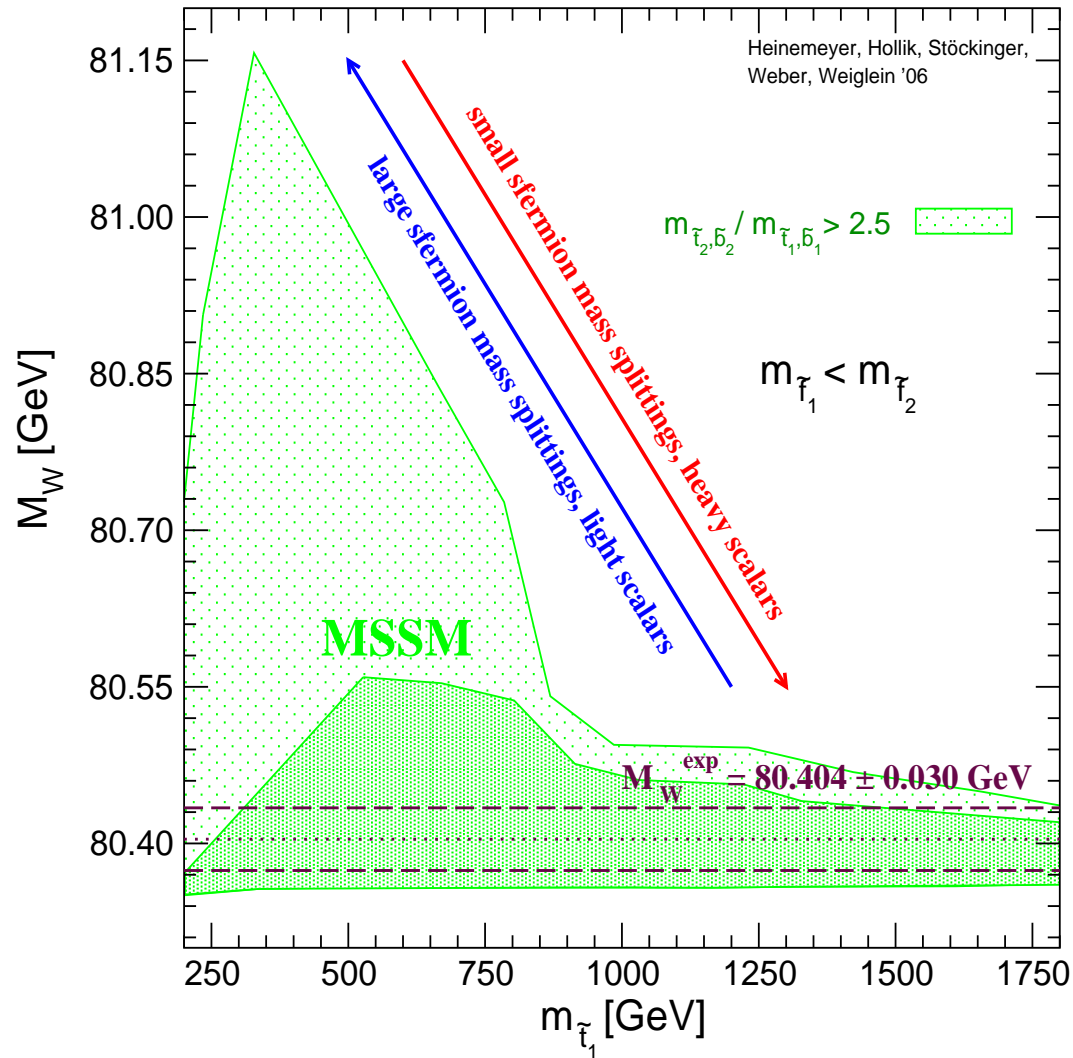
$$M_{\text{SUSY}} = 1 \text{ TeV}$$

$$\Rightarrow \delta M_W^{\text{MSSM-SM}} < 10 \text{ MeV}$$

$$M_{\text{SUSY}} > 3 \text{ TeV}$$

$$\Rightarrow \delta M_W^{\text{MSSM-SM}} < 1 \text{ MeV}$$

Prediction of M_W : Parameter scan



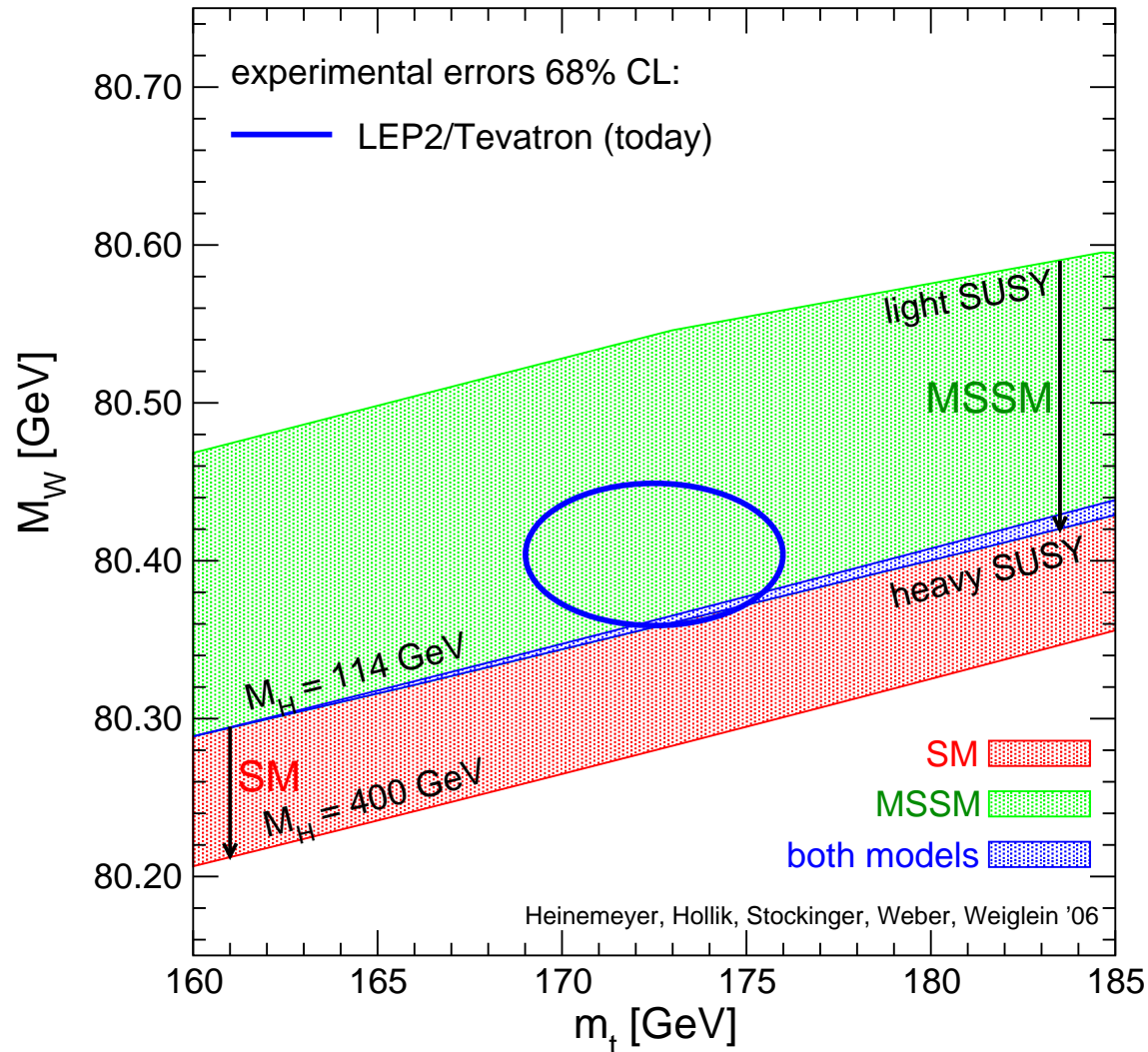
Scatter **all** relevant parameters **independently** over **wide** ranges:
masses $\lesssim 2 \text{ TeV}$
 $\tan \beta = 1.1 \dots 60$

Experimental constraints from LEP, Tevatron included

huge splitting
 \Rightarrow huge correction to M_W
(experimentally excluded)

Impact of M_W on preferred SUSY masses?

Prediction for M_W in the **SM** and the **MSSM** :



MSSM band:

scan over
SUSY masses

overlap:

SM is MSSM-like
MSSM is SM-like

SM band:

variation of M_H^{SM}

⇒ M_W is not enough

⇒ use more precision obs.

→ talk by S.H. later today :-)

4. Remaining theoretical (intrinsic) uncertainties

[J. Haestier, S.H., D. Stöckinger, G. Weiglein '05]

[S.H., W. Hollik, D. Stöckinger, A.M. Weber, G. Weiglein '06]

Estimate missing SUSY corrections order by order:

- $\mathcal{O}(\alpha_t^2, \alpha_t \alpha_b, \alpha_b^2)$: beyond existing leading contributions
- $\mathcal{O}(\alpha \alpha_s)$: beyond $\Delta\rho$ approx.
- $\mathcal{O}(\alpha \alpha_s^2)$
- $\mathcal{O}(\alpha^2 \alpha_s)$
- $\mathcal{O}(\alpha^3)$
- missing phase dependence at two-loop

⇒ evaluate for $M_{\text{SUSY}} = 300, 500, 1000$ GeV

Combine with SM uncertainty: $\delta M_W^{\text{SM, intr.}} = 4$ MeV

$$\delta M_W^{\text{SUSY, intr.}} = 5 - 11 \text{ MeV}$$

(depending on M_{SUSY})

5. Conclusinos

- Precision observables
 - can give valuable information about the “true” Lagrangian
 - can provide bounds on SUSY parameter space
 - ⇒ M_W prominent example
- Combination of
 - full SM result
 - all available MSSM corr. ⇒ best prediction of M_W in the MSSM
 - ⇒ remaining uncertainty: $\delta M_W^{\text{intr.}} = 5 - 11 \text{ MeV}$ (depending on M_{SUSY})
- Effects of certain sectors:
 - \tilde{t}/\tilde{b} sector: $\delta M_W \gtrsim 20 \text{ MeV}$, slepton sector: much smaller
 - complex phases in \tilde{t}/\tilde{b} sector: $\delta M_W \gtrsim 5 \text{ MeV}$
 - chargino/neutralino sector: $\delta M_W \gtrsim 20 \text{ MeV}$
 - $\delta m_t = \pm 2.3 \text{ GeV} \Rightarrow \delta M_W \approx \pm 14 \text{ MeV} \Rightarrow$ largest parametric unc.
- Prediction of M_W in special scenarios:
 - SPS: only “lower” masses show agreement at 1σ
 - Split SUSY: basically no visible effect
 - Decoupling limit: $M_{\text{SUSY}} > 3 \text{ TeV} \Rightarrow \delta M_W^{\text{MSSM-SM}} < 1 \text{ MeV}$
 - Parameter scan: very large M_W only for very large splitting in \tilde{t}/\tilde{b}