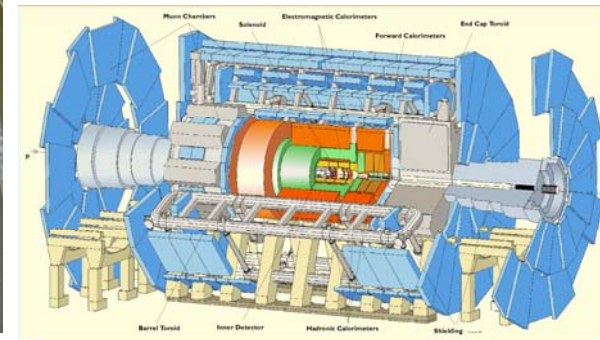
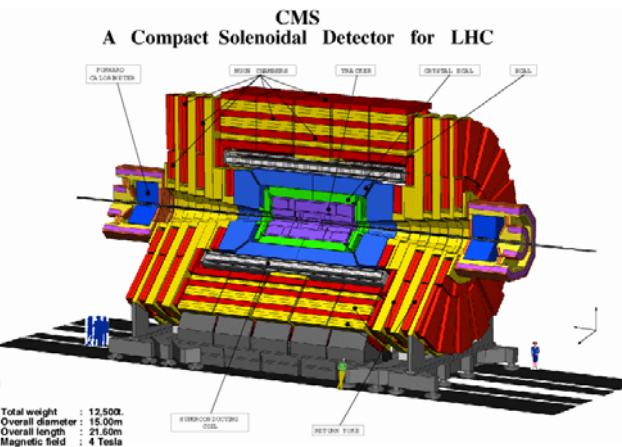


# LHC Physics Landscape

What more will we know a few years from now?  
Possible scenarios



G. Azuelos  
TRIUMF/ U. Montreal



*apologies for ATLAS bias...*

## → EWSB

- SM Higgs
- SUSY
  - Higgs
  - particle spectrum → *talk by Gudrid Moortgat-Pick*
- alternatives: little Higgs, Higgsless models,  
(strong EWSB → *talk by Tim Barklow*)

## → gauge theories with extended symmetry

- $Z', W'$  → *talk by Steve Godfrey*
- LR model, E6

## → ew precision measurements

- *top* → *talk by Aurelio Juste on top couplings*

## → compositeness

## → extra dimensions

## → other topics...

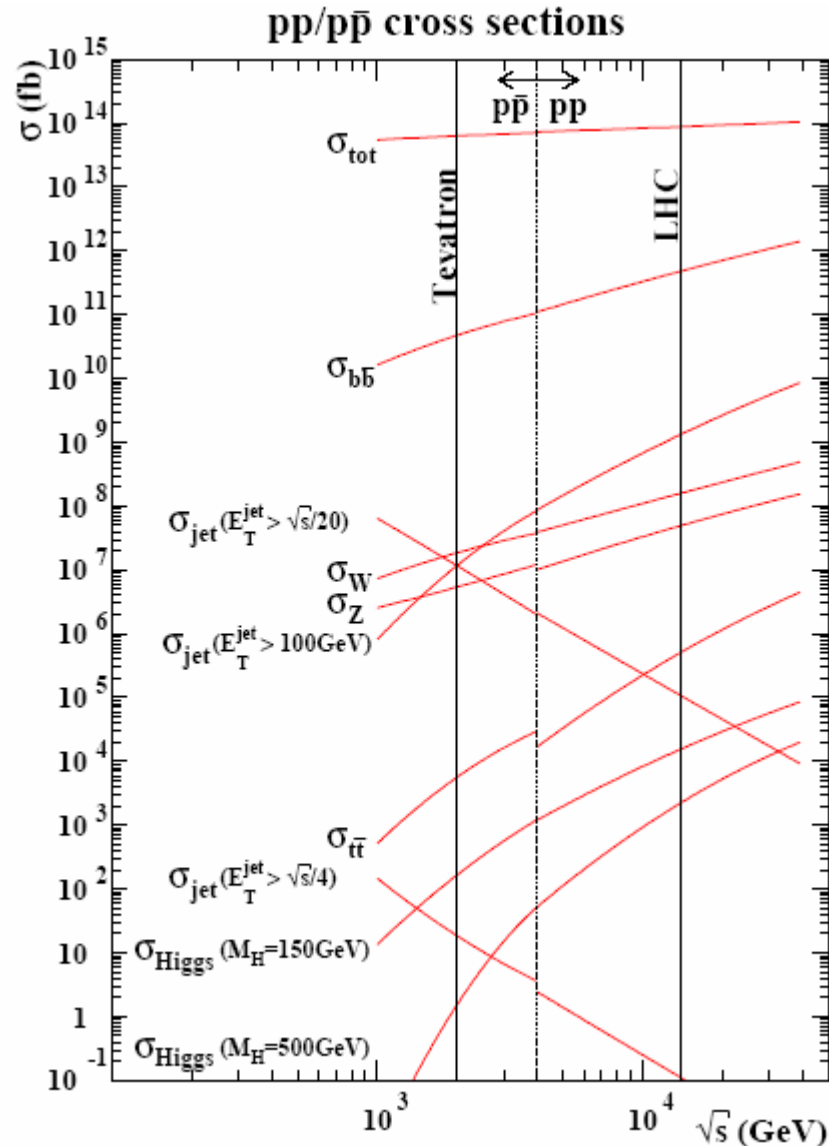
## → QCD

## → *B (LHCb), Heavy Ions (ALICE)*

*For an excellent review:*

*Physics Interplay of the LHC and the ILC, G. Weiglein, ed., (hep-ph/0410364)*

# LHC: basic features



➔ **p p collider,  $\sqrt{s} = 14$  TeV**

➤ (also heavy ions)

➔ **startup: 2007-2008**

➔ **nominal luminosity:  $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$**

➔ **bunch crossing: every 25 ns**

☺ **high energy:**

➤ parton collisions spanning up to  $\sqrt{\hat{s}} \sim 7$  TeV (gg collisions)

☹ **very high QCD cross section**

- NLO corrections important
- pileup:  $\sim 20$  min. bias events/bunch crossing
  - resolution effects and backgrounds
- overwhelming background for ew processes
  - good particle ID crucial

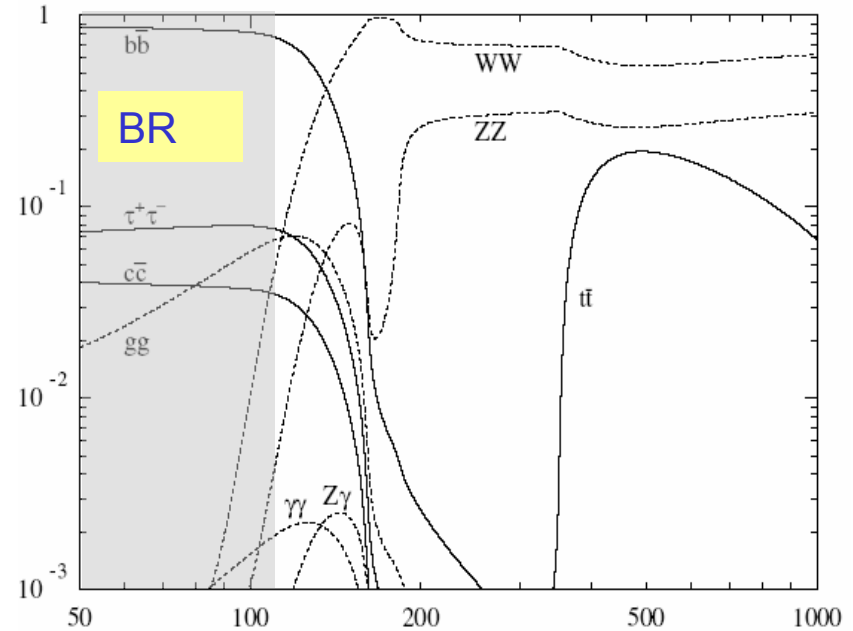
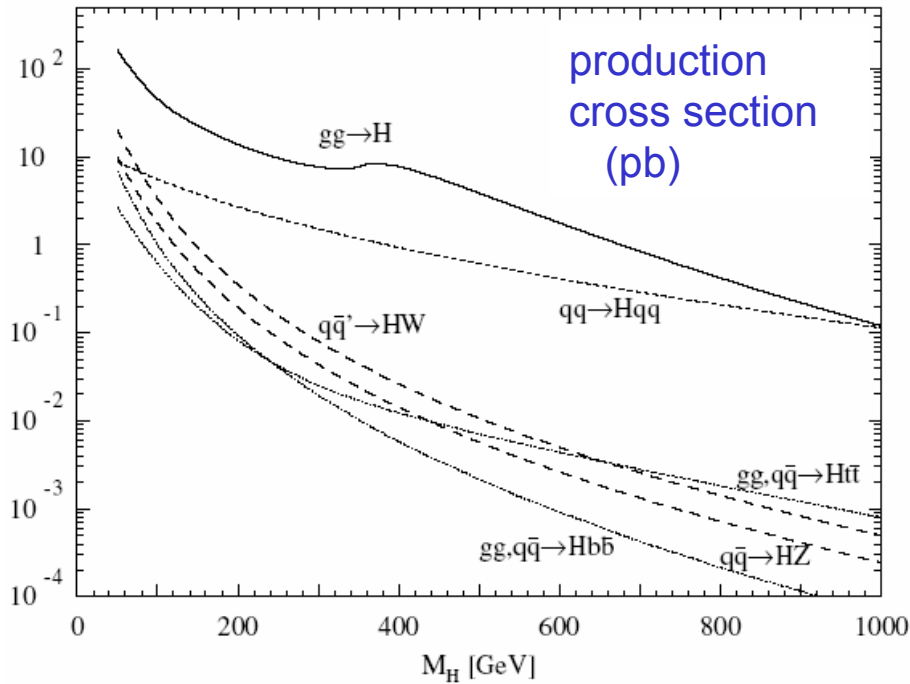
➔ **ILC:**

☹ lower energy, fixed cm energy  
but

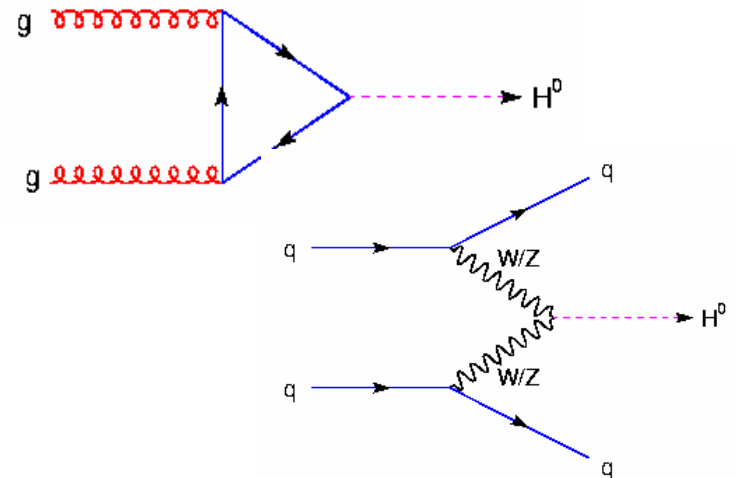
☺ precision measurements, with clean events,  
known cm energy, luminosity, and polarization

from: G. Weiglein et al., *Physics Interplay of the LHC and the ILC*, hep-ph/0410364

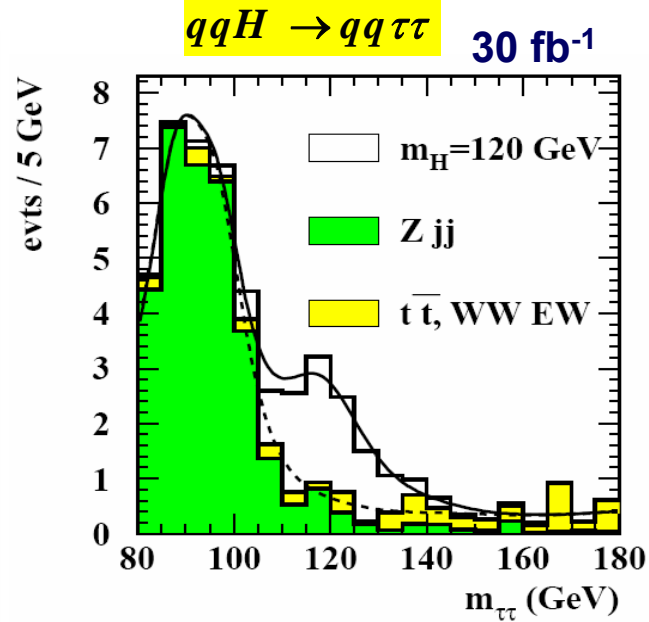
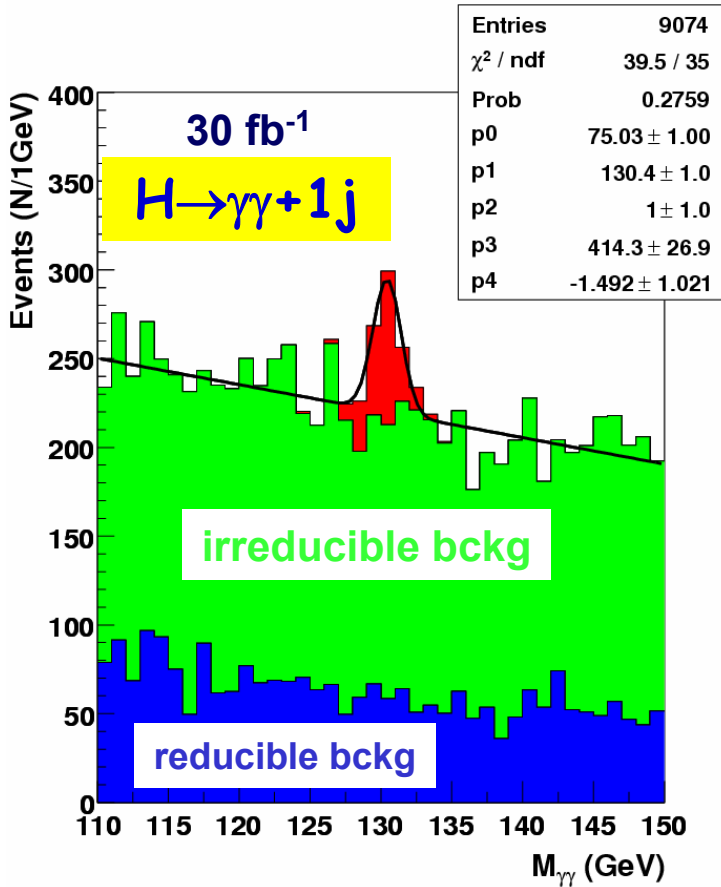
# EWSB: Standard Model Scenario



- production principally by  $gg$  fusion, but also by  $VB$  fusion
- QCD background limits study of  $H$  decay channels
- need to measure properties: *mass, gauge and Yukawa couplings, spin, CP quantum number, self-couplings*



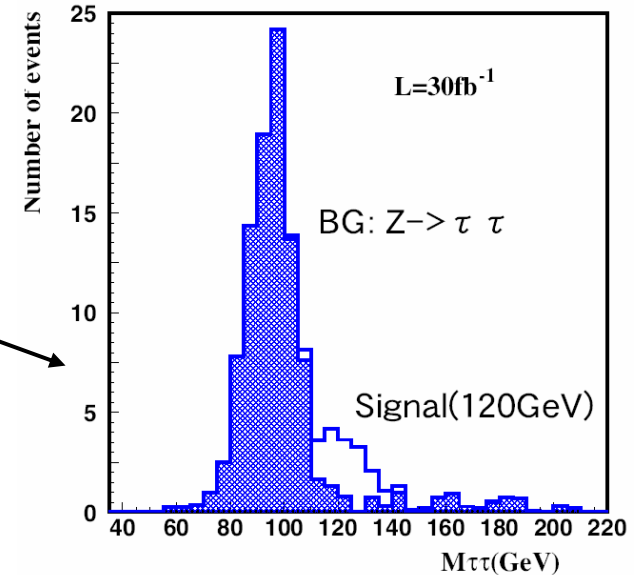
# EWSB: Standard Model Scenario: light Higgs



Asai et al.,  
 Eur.Phys.J. C32S2 (2004) 19  
 (hep-ph/0402254)

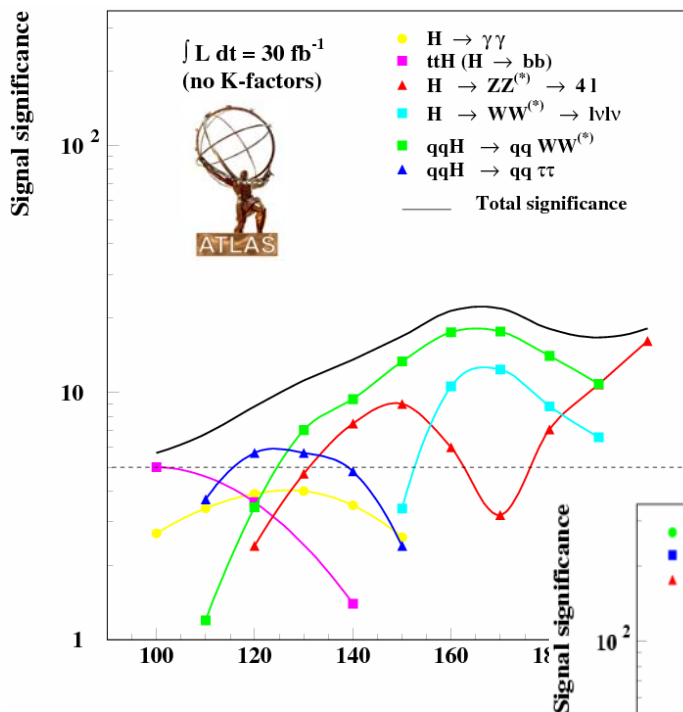
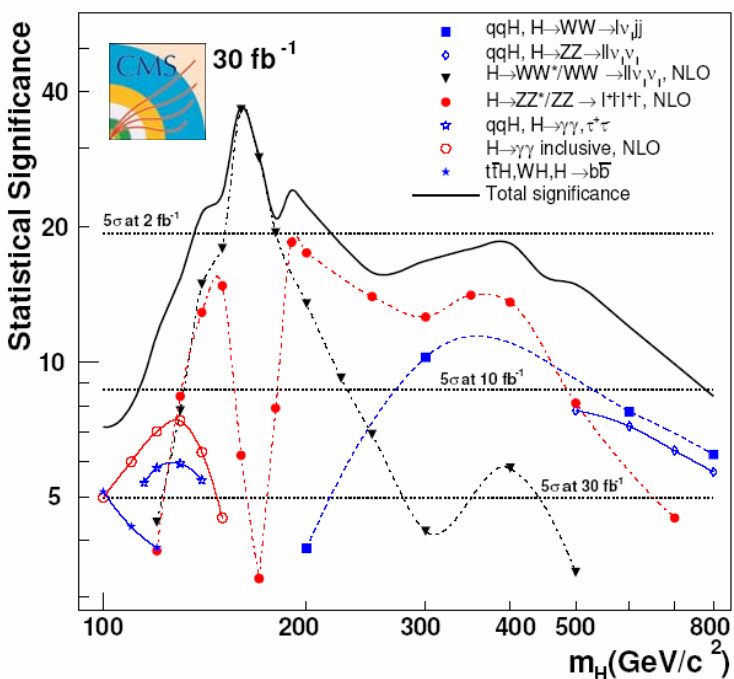
$\tau\tau \rightarrow ll + \cancel{E}_T$

$\tau\tau \rightarrow lh + \cancel{E}_T$



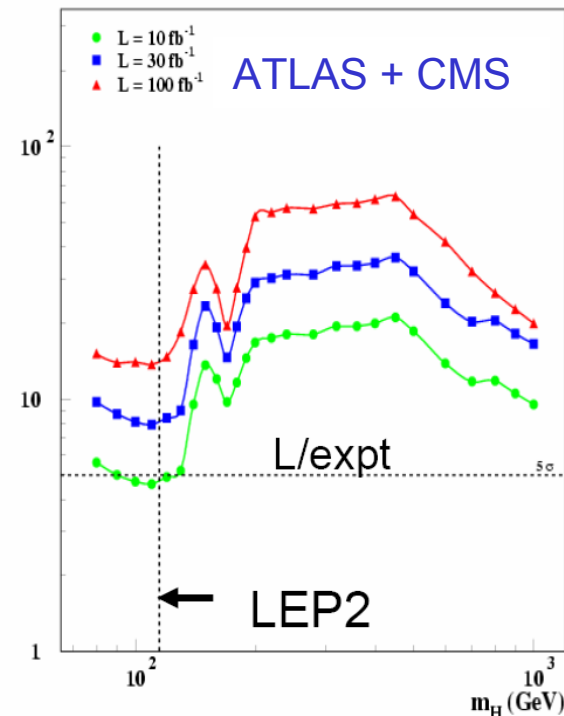
B. Mellado, ATLAS Phys workshop, 7/06/05

# EWSB: Standard Model Scenario



- Discovery assured in full mass range with  $\sim 10 \text{ fb}^{-1}$
- more than one channel of discovery

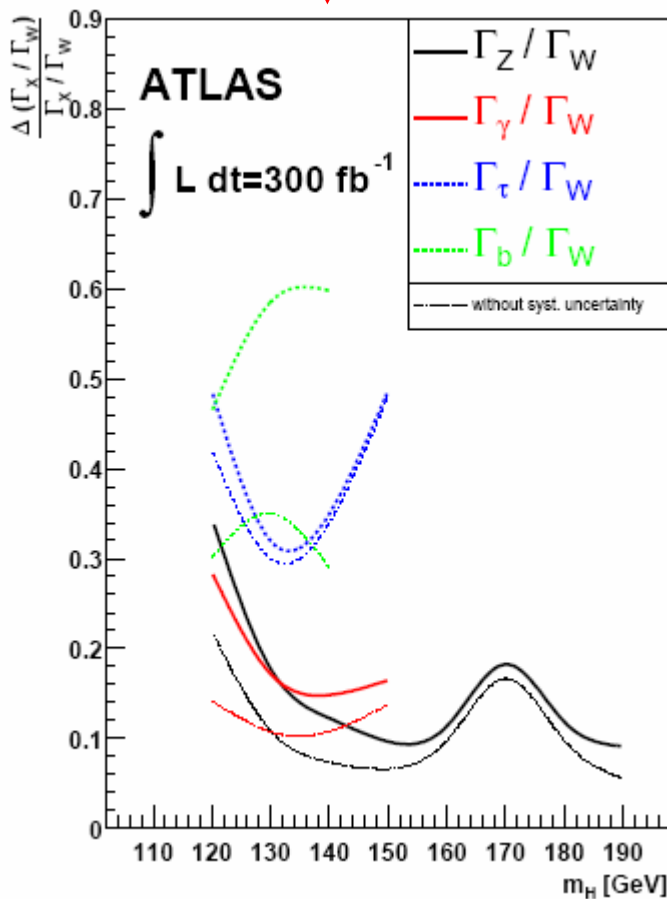
*at ILC,  
a final-state-independent mass  
determination is possible from mass  
recoil against Z in HZ production.*



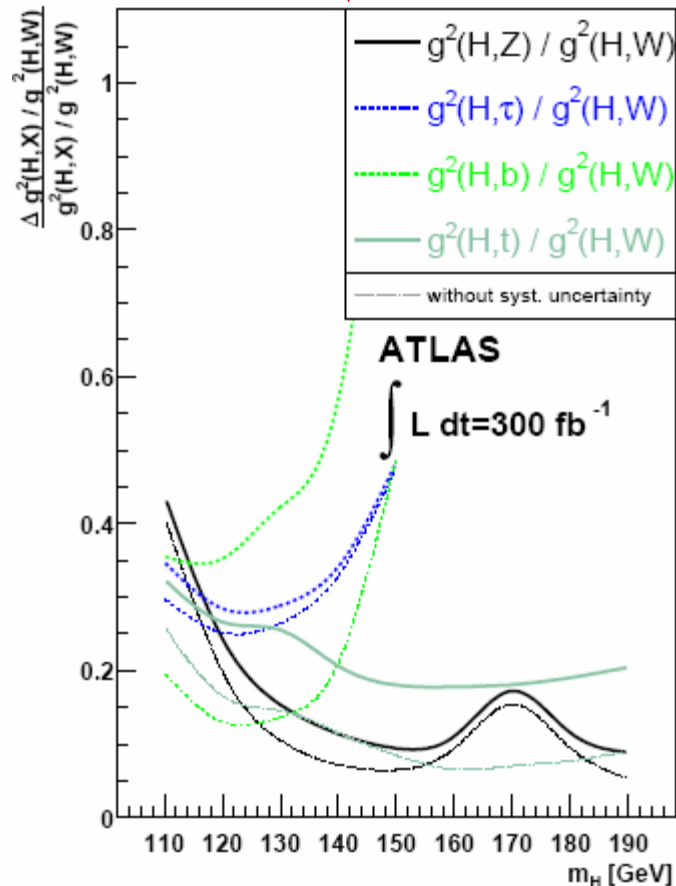
# EWSB: – Is this the SM Higgs ??

## Fractional errors

ratios of BR's



ratios of couplings



assume:

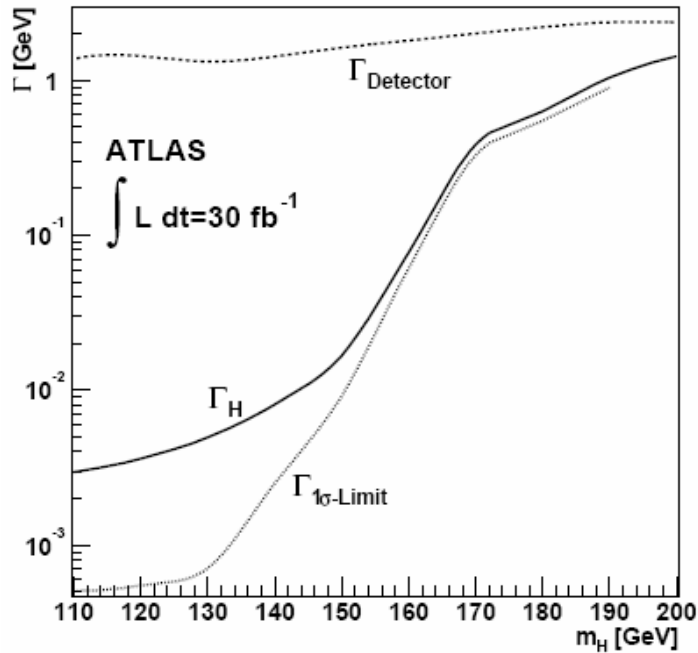
- CP even, spin-0 scalar
- only one Higgs boson
- no extra fermions

some theoretical model-dependence in evaluation of production and decay diagrams

from a compilation of various ATLAS studies, M. Dührssen, ATL-PHYS-2003-030



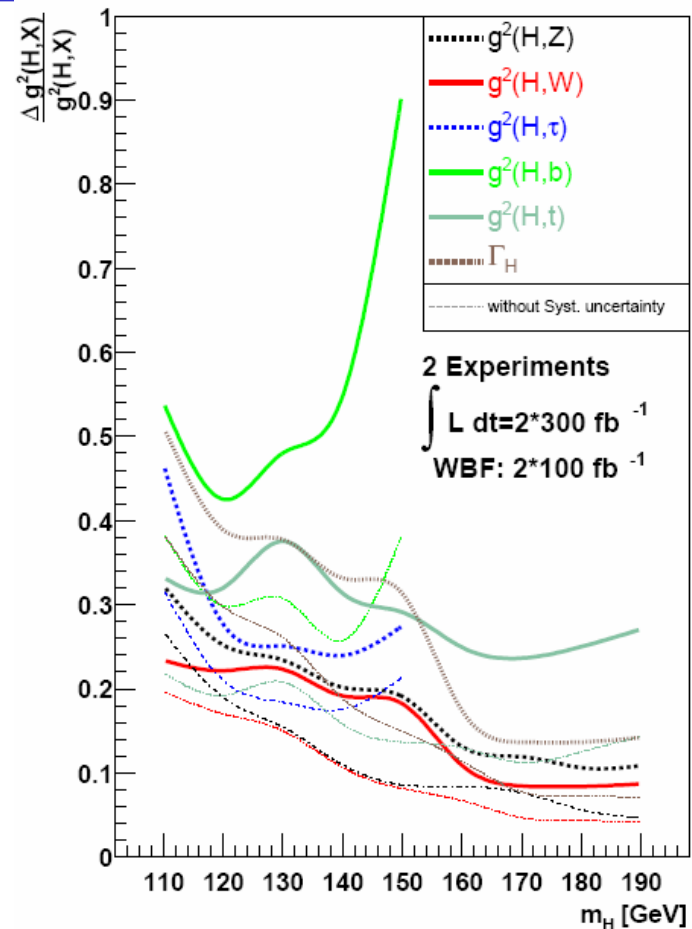
# Absolute couplings



lower limit on width of H boson  
 (from sum of all detectable decays)

M. Dührssen ATL-PHYS-2003-030

*at ILC, much higher precision attainable  
 (top-Yukawa requires high energy)*



M. Dührssen et al., XXXIXth Rencontres de Moriond, La Thuile, March 2004 (hep-ph/0407190)

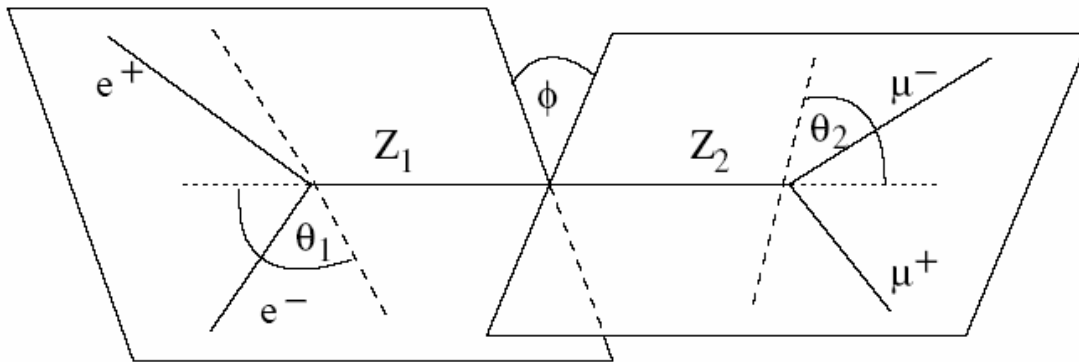
assumes gauge boson couplings not larger than in SM  
 (valid for generic 2HDM)



# Spin and CP Eigenstate

- production vertex not measurable  
gg fusion rules out spin 1 (Yang's theorem)
- decay vertex: decay to resonance pair necessary to measure CP

For  $m_H > 2m_Z$  :

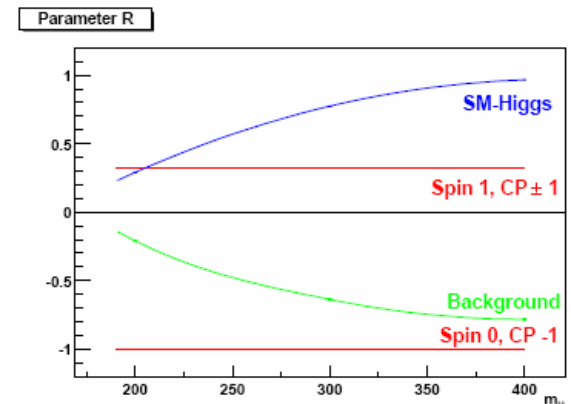
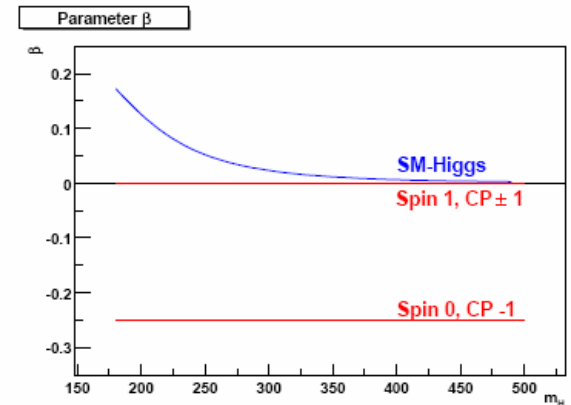
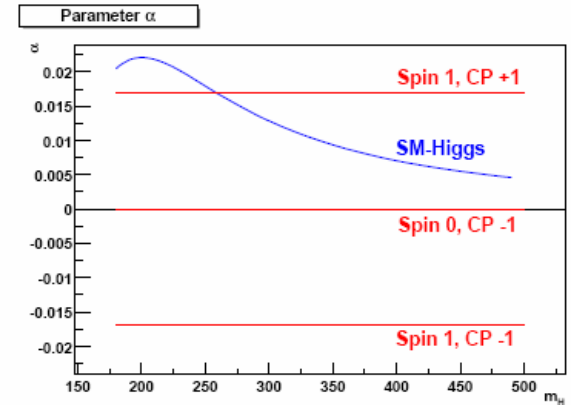


$$F(\phi) = 1 + \alpha \cos(\phi) + \beta \cos(2\phi)$$

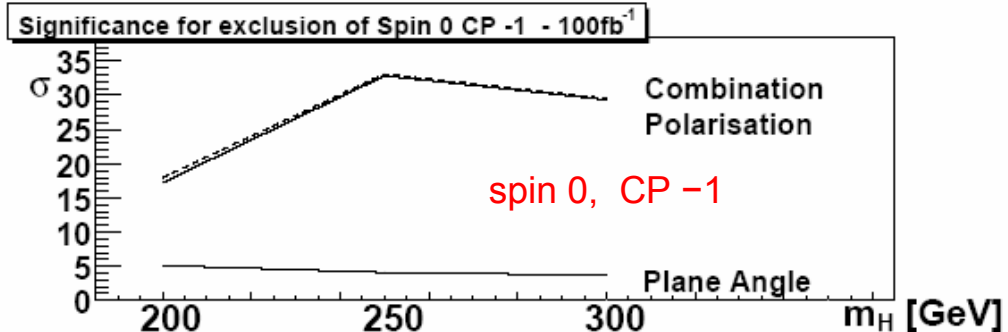
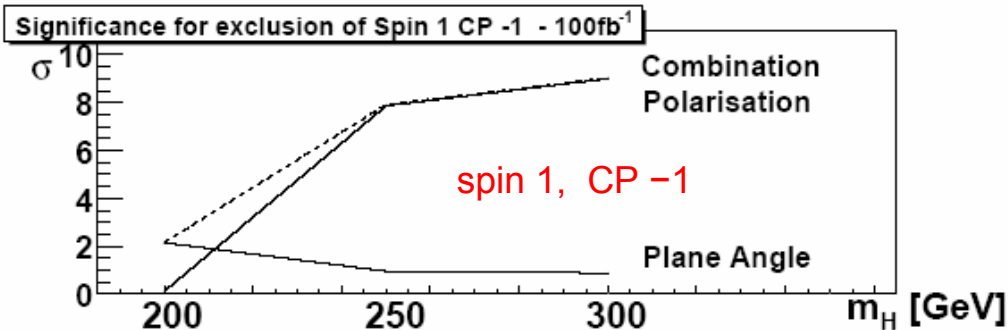
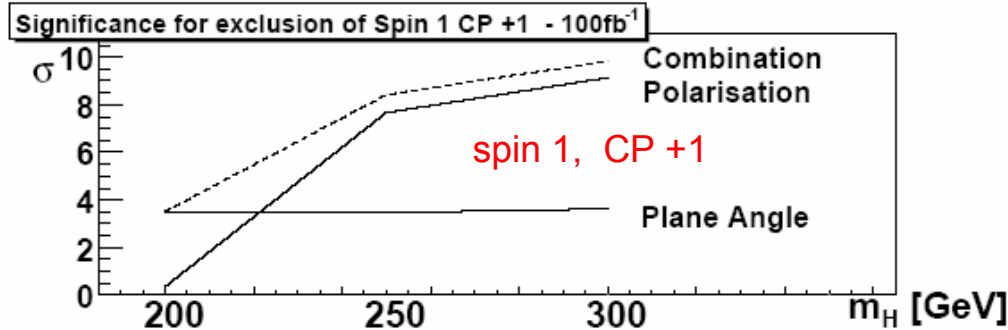
$$G(\theta) = T(1 + \cos^2 \theta) + L \sin^2 \theta$$

$$R = \frac{L - T}{L + T}$$

C.P. Buszello et al., *Eur. Phys. J C32 (2004) 209* (hep-ph/0212396)



# Spin and CP Eigenstate



- rule out spin 1
  - for  $m_H > 230$  GeV with  $100 \text{ fb}^{-1}$
  - for  $m_H \sim 200$  GeV, higher Luminosity
- rule out  $J=0, CP=-1$ 
  - for  $m_H > 200$  GeV with  $100 \text{ fb}^{-1}$

*In a general 2HDM, or MSSM with CP violation, the 3 neutral H bosons mix. Measuring the amount of CP violation will need combination of LHC/ILC.*

C.P. Buszello et al., *Eur. Phys. J C32 (2004) 209*  
(hep-ph/0212396)

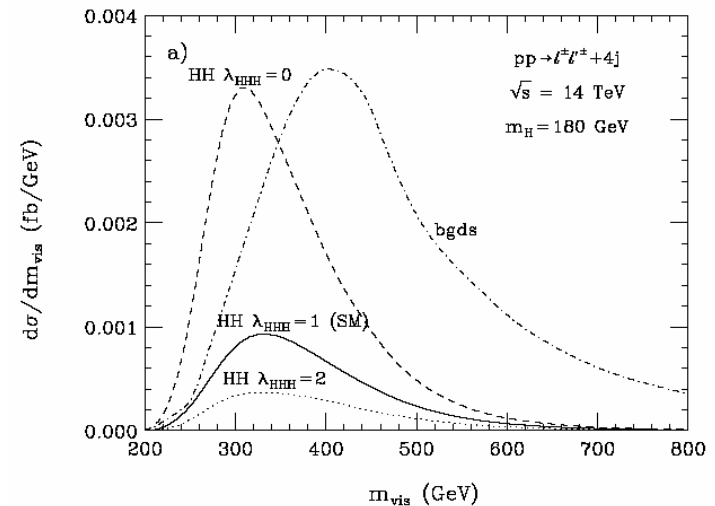
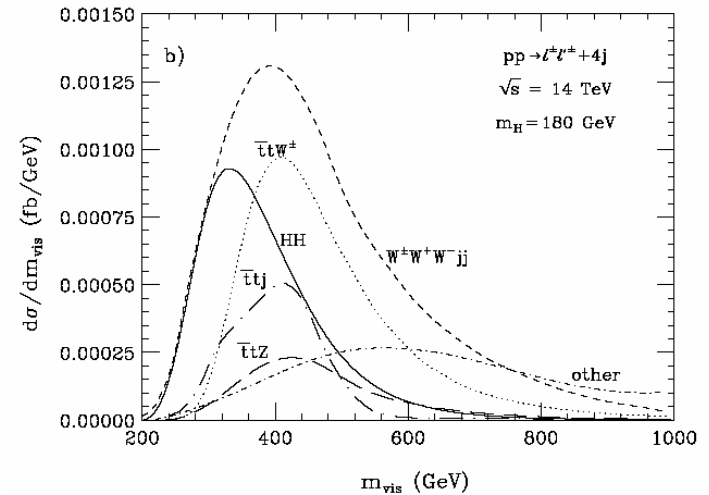
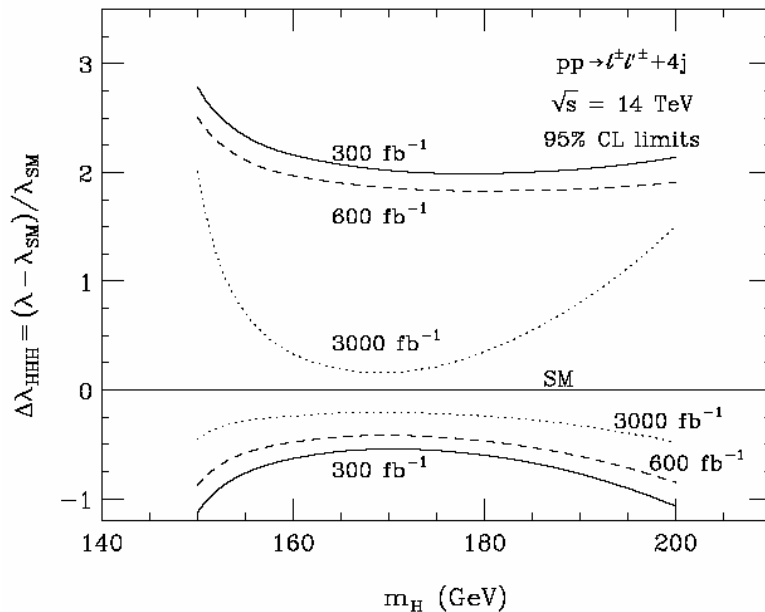
# Higgs self-coupling

Higgs self-coupling is a crucial ingredient of the SM

$$\text{Higgs potential: } V(\phi) = -\mu^2(\phi^\dagger\phi) - \lambda(\phi^\dagger\phi)^2$$

$$\hookrightarrow \mu^2 = -\lambda v^2 = -m_H^2/2$$

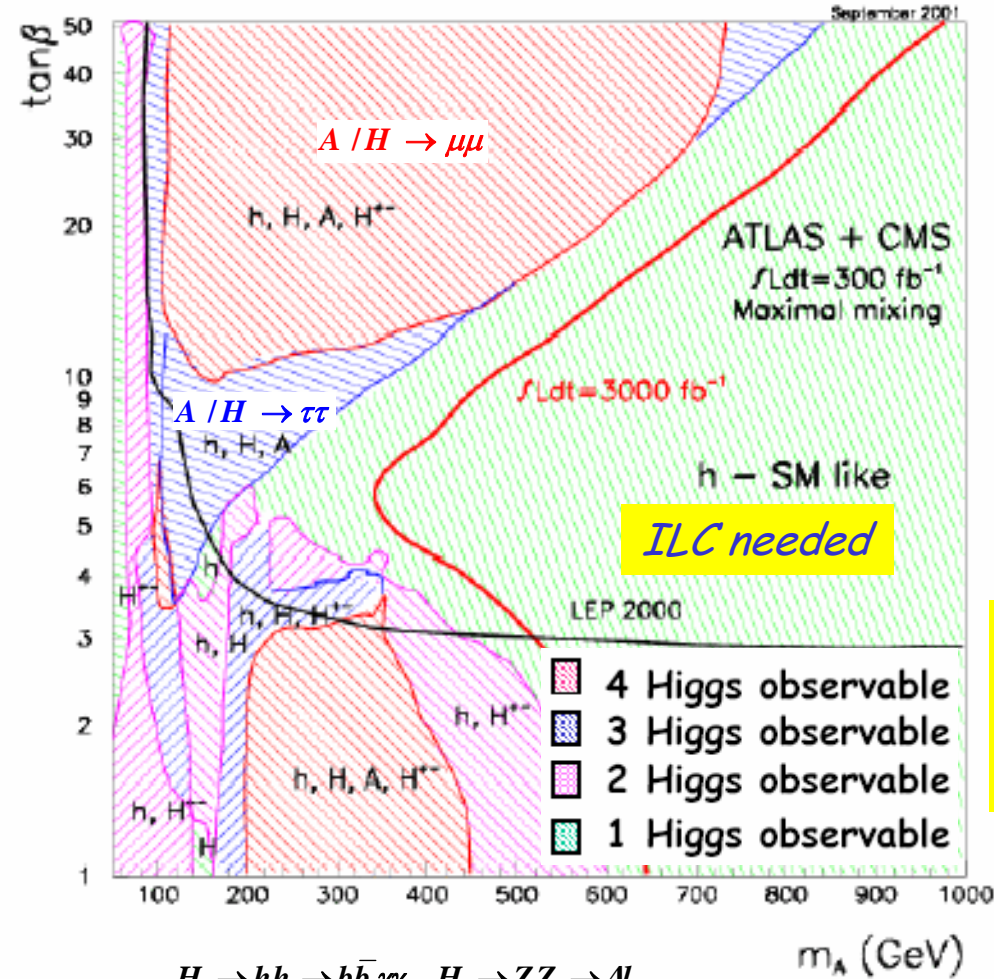
cross section for Higgs pair production at LHC very low, and backgrounds important, but some limits can be set



at ILC (1 TeV), higher precision for  $m_H < 140$  GeV ( $bb \ bb$  channel at ILC,  $\gamma \ bb$  at LHC). Precision Higgs properties important for LHC measurement at higher mass (top Yukawa,  $HWW$  coupling, total width...)  $\rightarrow$  see talk by T. Barklow

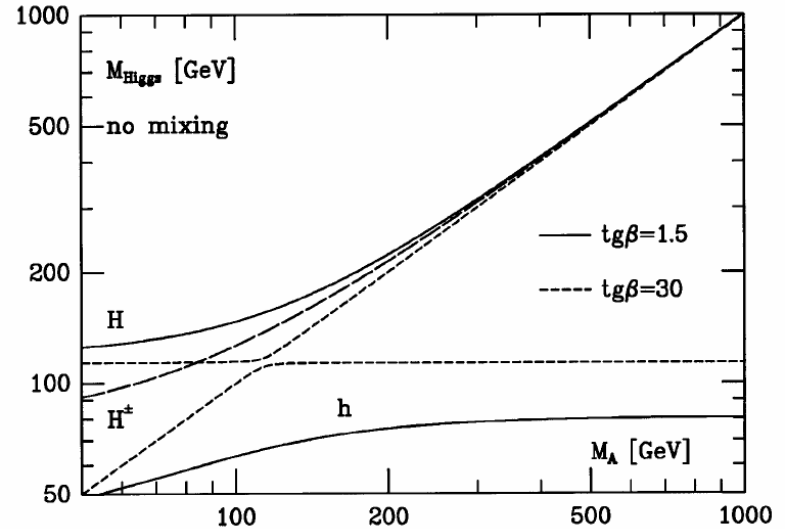
# EWSB: MSSM scenario

## 5 Higgs bosons, two parameters:



$$H \rightarrow hh \rightarrow b\bar{b}\gamma\gamma \quad H \rightarrow ZZ \rightarrow 4l$$

$$A \rightarrow Zh \rightarrow l\bar{l}b\bar{b} \quad A/H \rightarrow t\bar{t}$$



*in the decoupling limit, the ILC can help:*

- precision measurements
- $HA$  and  $H^+H^-$  production (need high energy)
- $\gamma\gamma \rightarrow H/A$  possible

*In general, in the presence of SUSY particles, the discovery reach changes*

## ➤ NMSSM

- extend MSSM by adding one singlet superfield → new  $h_3$  and  $A_2$ 
  - LHC:  $WW \rightarrow h \rightarrow aa \rightarrow bb \tau\tau$
  - ILC : also  $Zh \rightarrow Zaa; a \rightarrow bb \text{ or } \tau\tau$ , → talk by J. Gunion

## ➤ invisible channel

- difficult at LHC (VBF), but easier at ILC ( $Zh$ )

## ➤ radion (or other scalars)

- scalar state in Randall-Sundrum extra-dimension model
  - mixes with SM Higgs, can decay to  $hh$
  - some channels of discovery at LHC,

$$\phi \rightarrow hh \rightarrow \gamma\gamma b\bar{b}, \tau\tau b\bar{b}; \quad \phi \rightarrow ZZ^{(*)} \rightarrow 4l$$

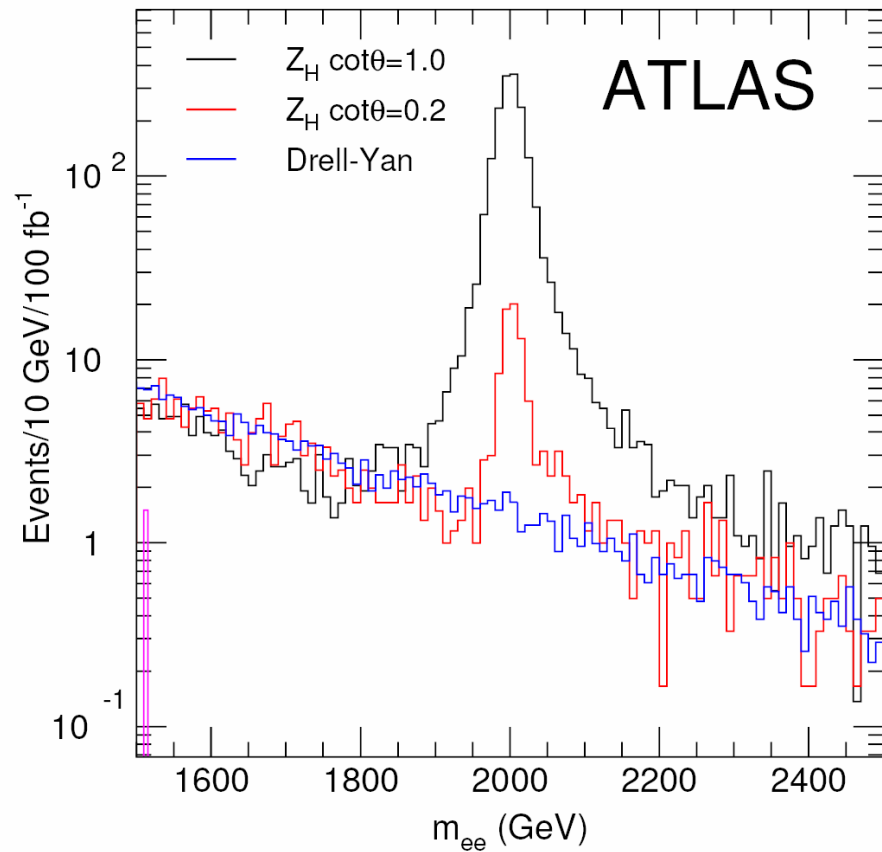
but distinguishing from MSSM Higgs difficult in some mass regions

## ➤ fermiophobic Higgs

- if mass originates from other mechanism...
  - use  $\gamma\gamma$  production at ILC

# little Higgs model

- Higgs is naturally light (pseudo Goldstone boson from breaking of higher symmetry)
- new states from higher symmetry (and isosinglet top) cancel quadratic divergences from radiative corrections
- LHC: discovery possible of new states in large part of parameter space
- with T-parity: missing energy signal could fake SUSY
- ILC: sensitive to presence of new particles in loops (TGC,  $h\gamma\gamma$ ,...)

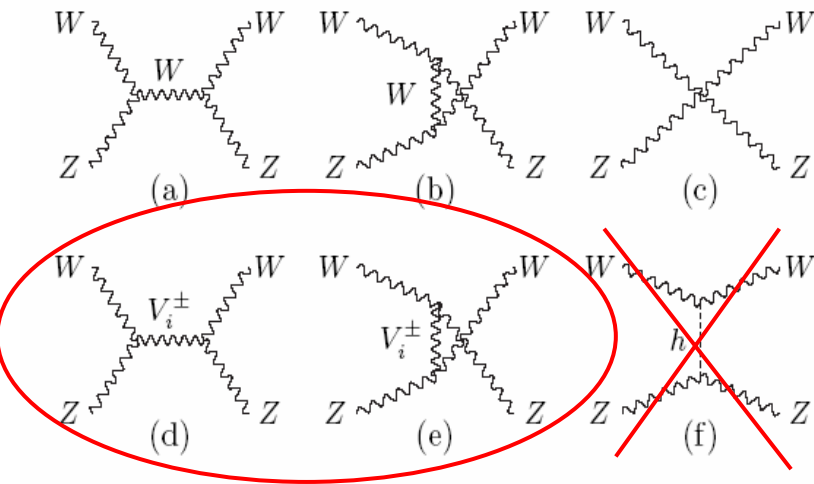


**SN-ATLAS-2004-038**

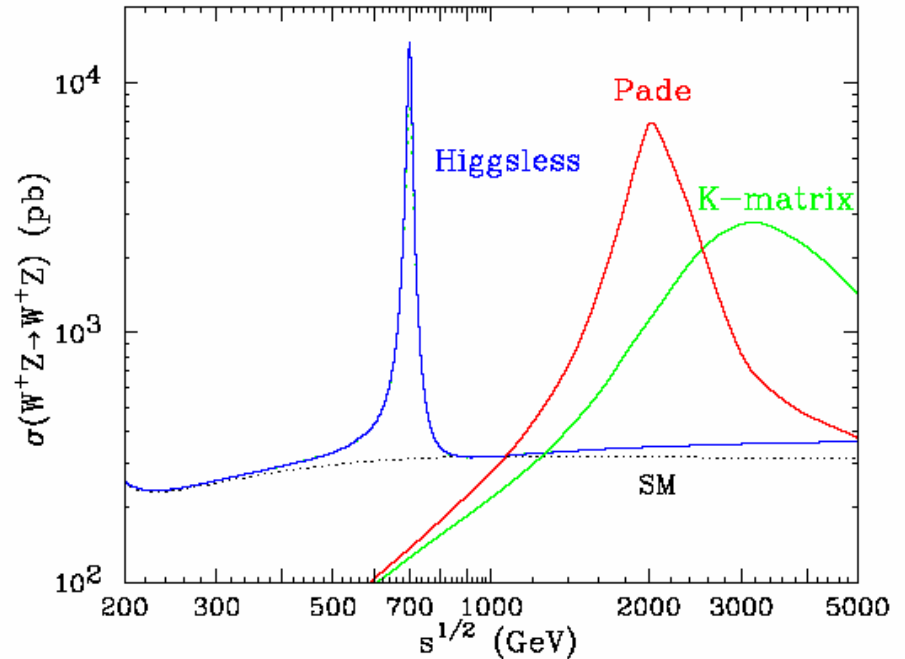
# Higgsless model

- EWSB from boundary conditions between branes in warped 5<sup>th</sup> dimension

C. Csaki, et al., Phys. Rev. D 69, 055006 (2004) [hep-ph/0305237]



Higher KK levels of vector bosons contribute and regularize the cross section, in the absence of a Higgs boson



A. Birkedal, K. T. Matchev, and M. Perelstein, hep-ph/0412278, 0508185

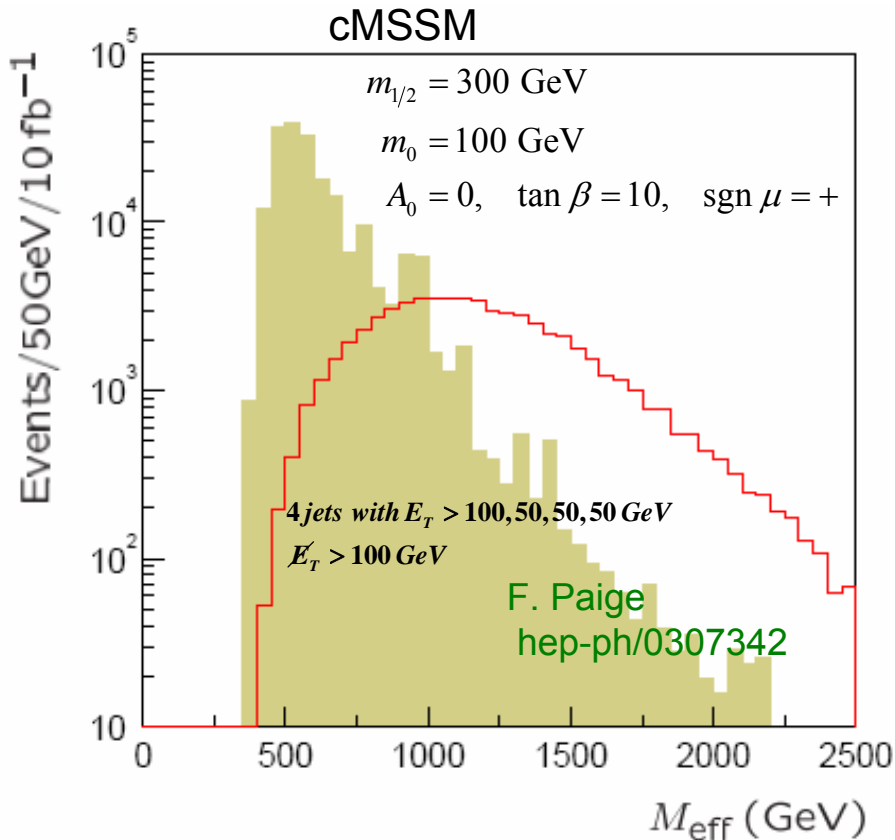
**LHC: How to distinguish from technicolor (QCD-like) resonance? need to see more KK states, but insufficient energy...**



## ➤ Inclusive signal with jets (R-conservation)

- trigger: Jets +  $\cancel{E}_T$
- main background: SUSY

$g g \rightarrow K$  cascade  $\rightarrow \cancel{E}_T \cancel{E}_T + jets + leptons$



$$M_{\text{eff}} \equiv \sum_{\text{jets}} E_T^{\text{jets}} + \cancel{E}_T : \text{good measure of } M_{\text{SUSY}}$$

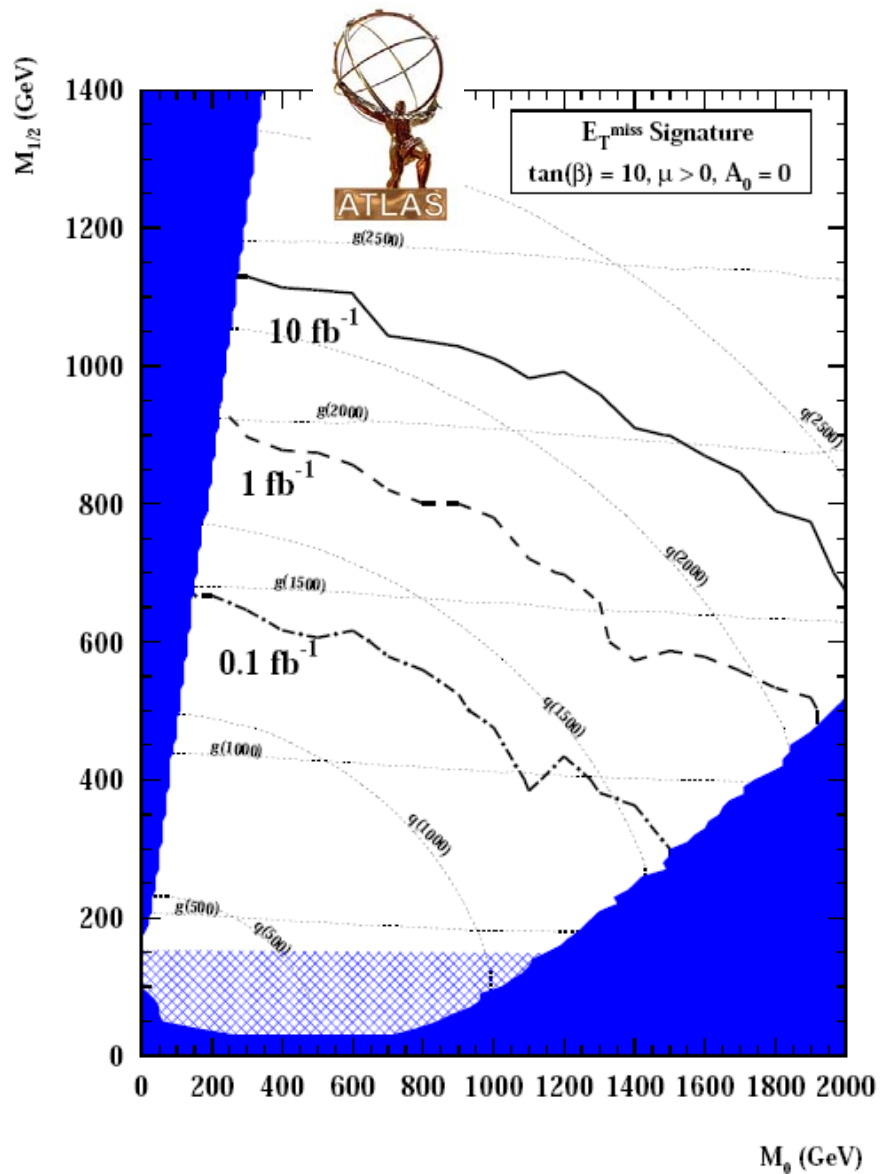
(linear relation)

⇒ preselect SUSY-rich sample

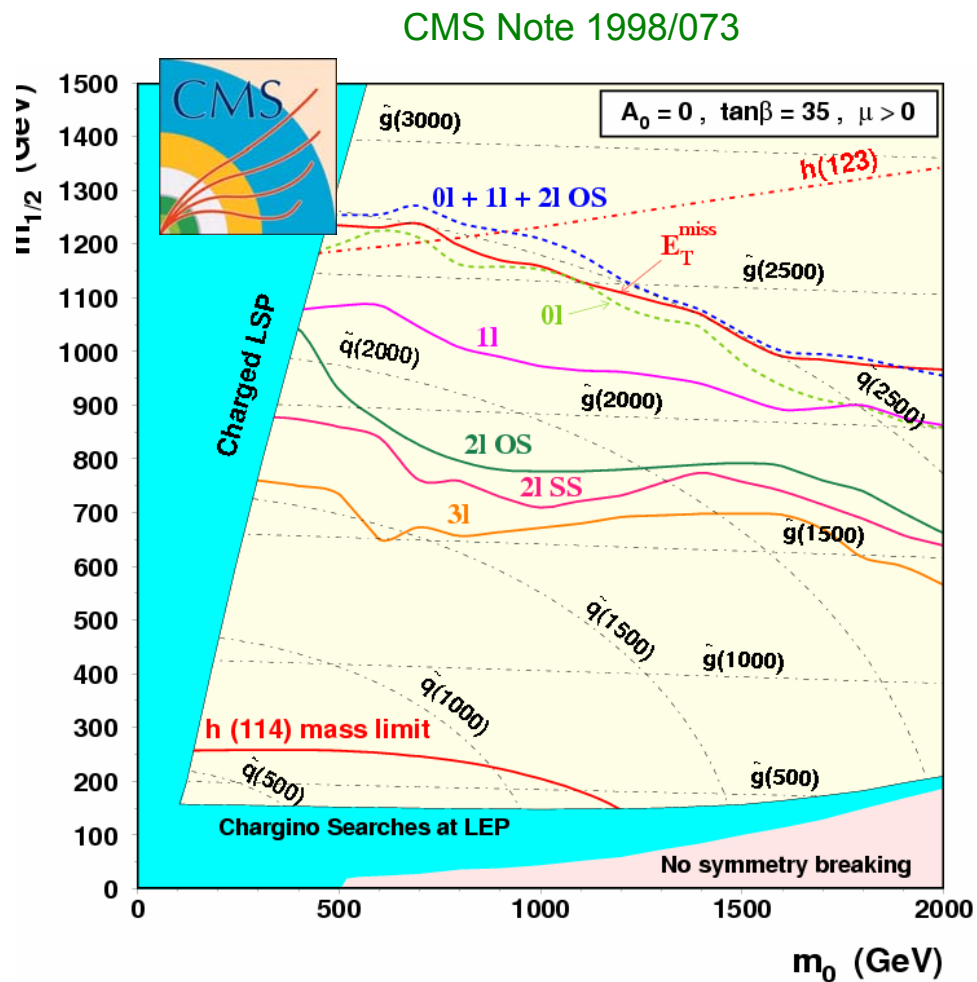
**note:**

background ( $t\bar{t}, W + j, Z + j, \text{QCD}$ )  
 increases by factor 2-4 with multi-parton  
 (ME-PS) generators

# SUSY - inclusive search



LHC reach,  $100 \text{ fb}^{-1}$ , various signals:

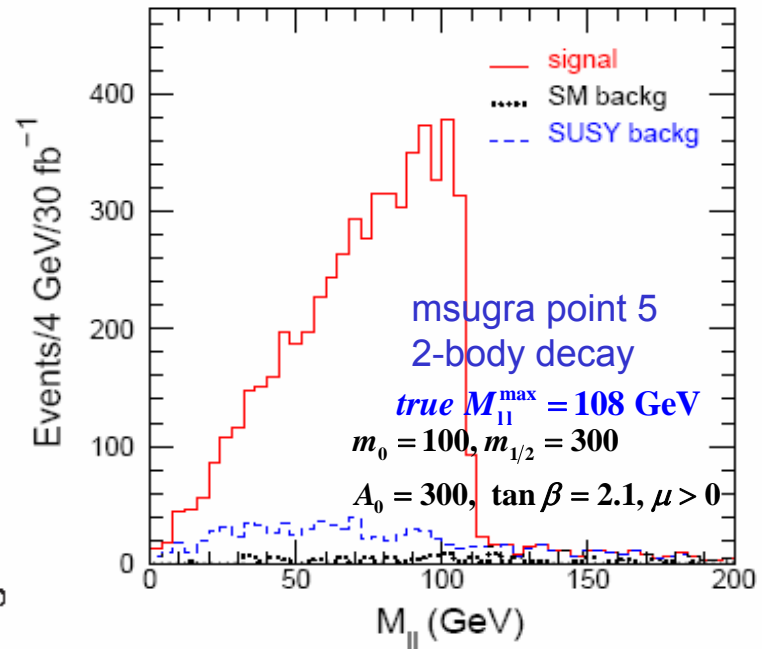
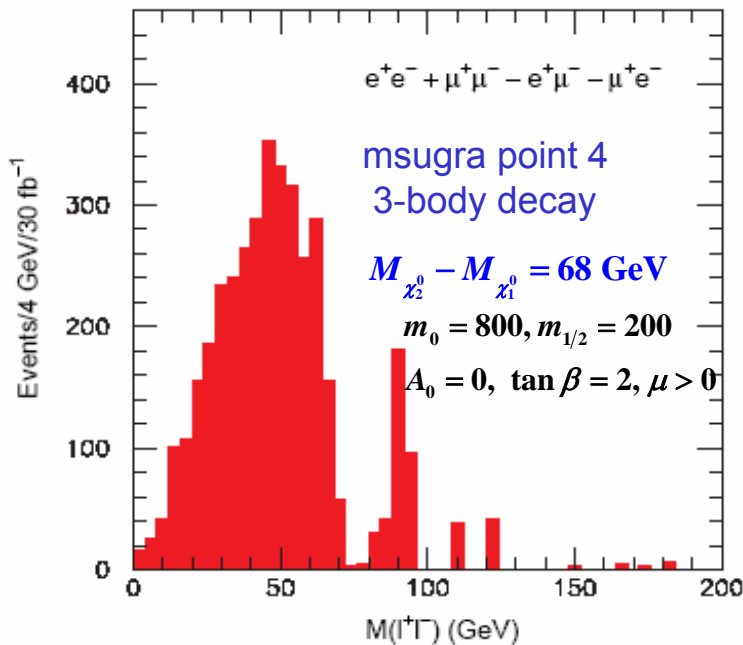


# Supersymmetry

## Mass reconstruction from di-lepton endpoints

3-body decay:  $\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \ell^+ \ell^- \Rightarrow M_{\parallel} < M_{\tilde{\chi}_2^0} - M_{\tilde{\chi}_1^0}$

2-body decay:  $\tilde{\chi}_2^0 \rightarrow \tilde{\nu}_1^m \rightarrow \tilde{\chi}_1^0 \ell^+ \ell^- \Rightarrow M_{\parallel} < \frac{\sqrt{(M_{\tilde{\chi}_2^0}^2 - M_{\tilde{\nu}_1^m}^2)(M_{\tilde{\nu}_1^m}^2 - M_{\tilde{\chi}_1^0}^2)}}{M_{\tilde{\nu}_1^m}}$

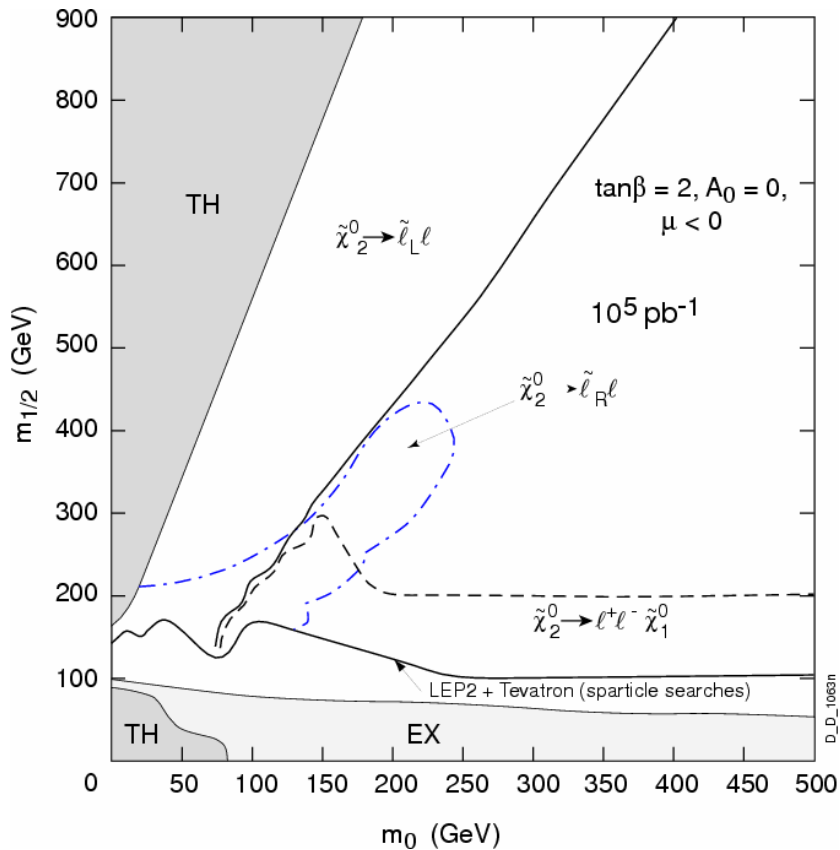


SM background removed by subtracting distributions with opposite flavor leptons  
 precision in end-point measurement:  $\sim 1\text{-}2\%$

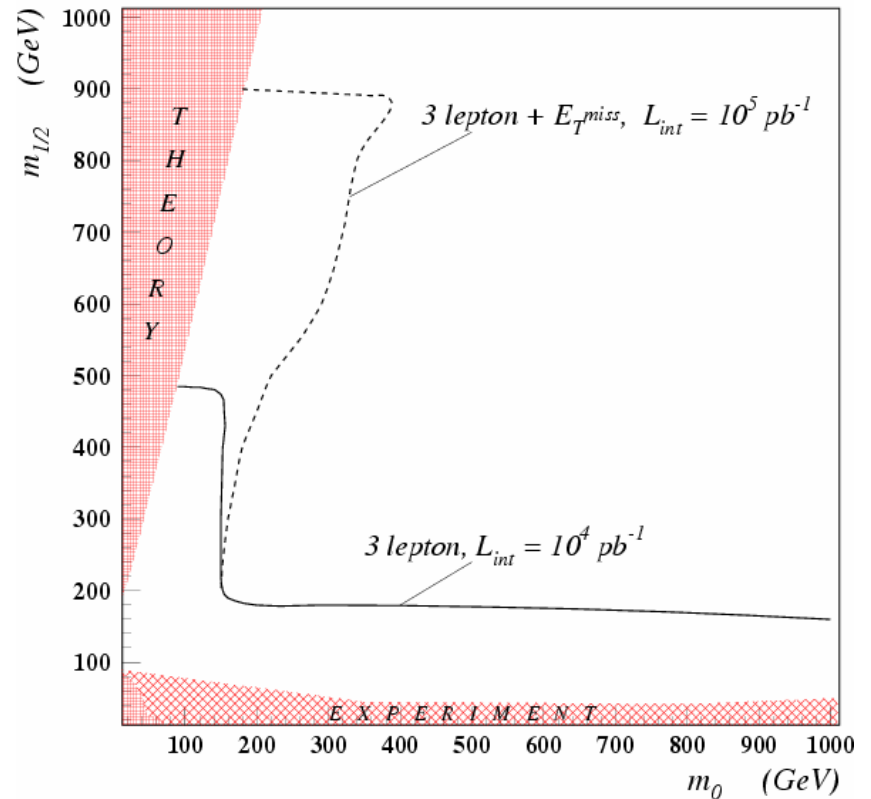
*example of SPS1a point fully studied in LHC/ILC report:*

## ➔ Inclusive signal from di-lepton endpoint

2 leptons +  $E_T^{\text{miss}}$  + jets



Direct  $\tilde{\chi}_1^0 \tilde{\chi}_2^0$  production  $\rightarrow 1\% \text{ @ } 1\% \tilde{\chi}_1^0$   
*mSUGRA parameters:  $\tan\beta=2, A_0=0, \mu<0$*



CMS

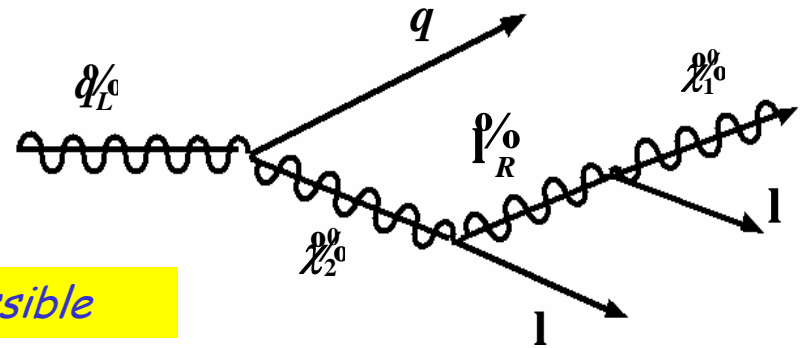
# Supersymmetry

reconstruction higher up in the chain possible

mass relations with  $\sim 1\%$  precision

$\tilde{\chi}_1^0$  mass:  $\sim 10\%$

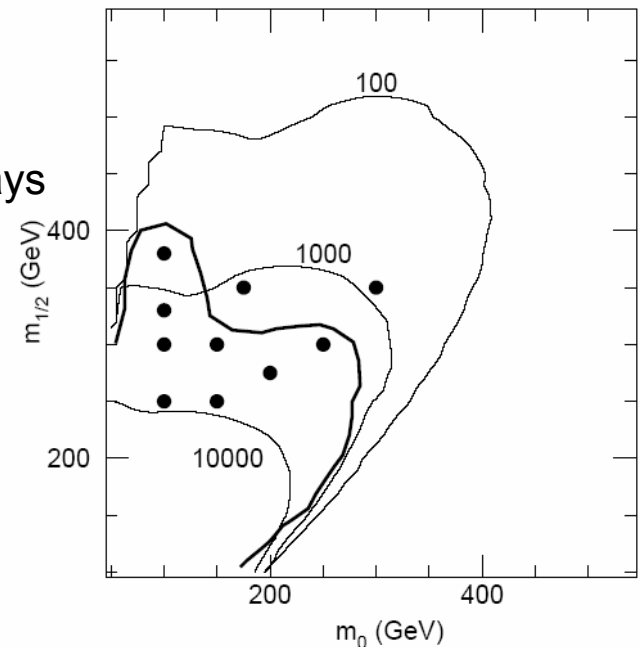
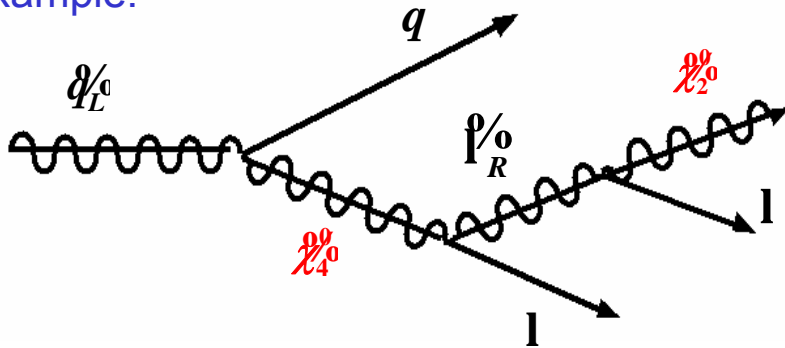
at ILC: precision of  $\leq 1\%$  possible



Heavy gauginos

$\tilde{\chi}_4^0, \tilde{\chi}_2^0$  have non-negligible gaugino content  
 $\rightarrow$  participate in squark and gluino decays

example:



LHC reach,  $100 \text{ fb}^{-1}$

F. Paige,  
 Physics at LHC, Vienna

# SUSY masses and parameters

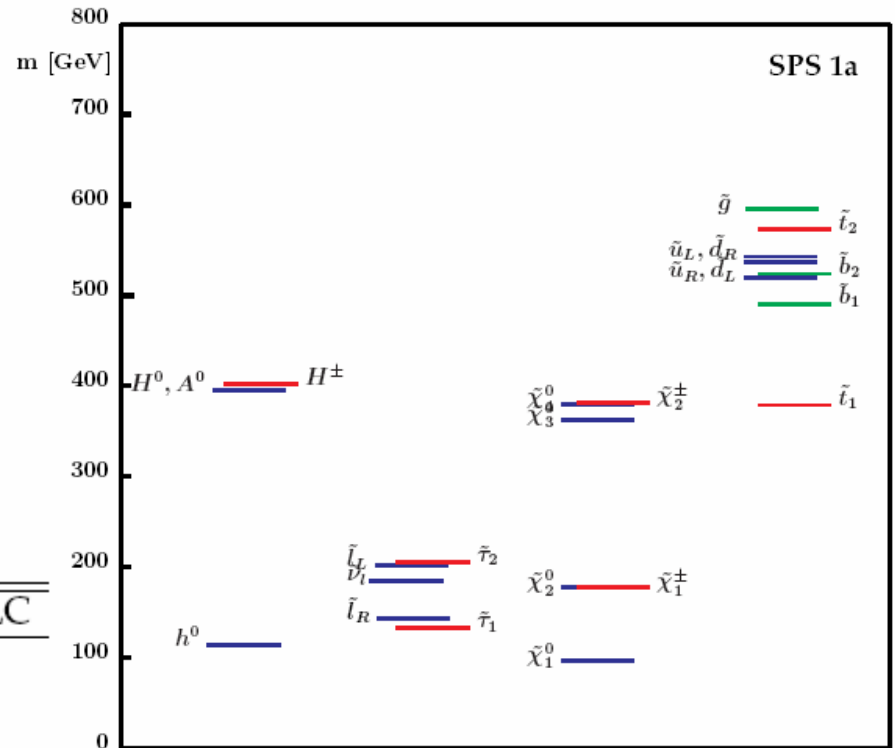
	LHC	LHC+LC
$\Delta m_{\tilde{\chi}_1^0}$	4.8	0.05 (LC input)
$\Delta m_{\tilde{\chi}_2^0}$	4.7	0.08
$\Delta m_{\tilde{\chi}_4^0}$	5.1	2.23
$\Delta m_{\tilde{l}_R}$	4.8	0.05 (LC input)
$\Delta m_{\tilde{l}_L}$	5.0	0.2 (LC input)
$\Delta m_{\tau_1}$	5-8	0.3 (LC input)
$\Delta m_{\tilde{q}_L}$	8.7	4.9
$\Delta m_{\tilde{q}_R}$	7-12	5-11
$\Delta m_{\tilde{b}_1}$	7.5	5.7
$\Delta m_{\tilde{b}_2}$	7.9	6.2
$\Delta m_{\tilde{g}}$	8.0	6.5

	LHC (0.5% jet scale)	LHC (0.5% jet scale) + LC
$\Delta m_{\tilde{q}_L}$	7.8	2.6
$\Delta m_{\tilde{b}_1}$	6.0	3.7
$\Delta m_{\tilde{b}_2}$	6.4	4.3
$\Delta m_{\tilde{g}}$	6.0	3.7

SUSY Parameters				Mass Predictions	
$M_1$	$M_2$	$\mu$	$\tan \beta$	$m_{\tilde{\chi}_2^\pm}$	$m_{\tilde{\chi}_3^0}$
$99.1 \pm 0.1$	$192.7 \pm 0.3$	$352.4 \pm 2.1$	$10.2 \pm 0.6$	$378.5 \pm 2.0$	$358.8 \pm 2.1$

Table 5.23: SUSY parameters with  $1\sigma$  errors derived from the combined analysis of the LHC and LC data with  $\delta m_{\tilde{\chi}_2^0} = 0.08$  GeV and  $\delta m_{\tilde{\chi}_4^0} = 2.23$  GeV derived from the LHC when using the LC input of  $\delta m_{\tilde{\chi}_1^0} = 0.05$  GeV.

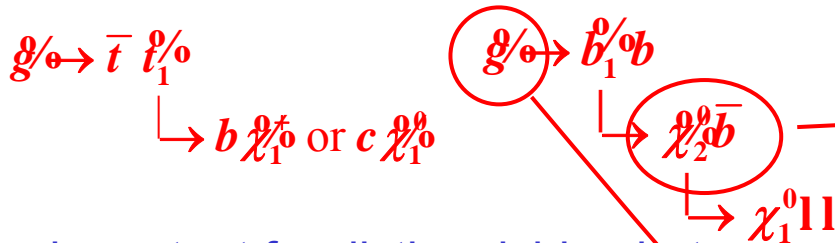
from LHC/ILC report



**SPS 1a :**  
 $m_0 = 100$  GeV,  
 $m_{1/2} = 250$  GeV,  
 $A_0 = -100$  GeV,  
 $\tan \beta = 10, \mu > 0$

# Supersymmetry

## 3<sup>rd</sup> generation squarks



- important for distinguishing between SUSY models
- difficult, but possible to measure combination of masses ( $m_{tb}$ )
- knowing  $\tilde{g}$  masses, can reconstruct  $g$  and  $\tilde{b}$
- mass relation method:
  - constraint: require same masses in several events

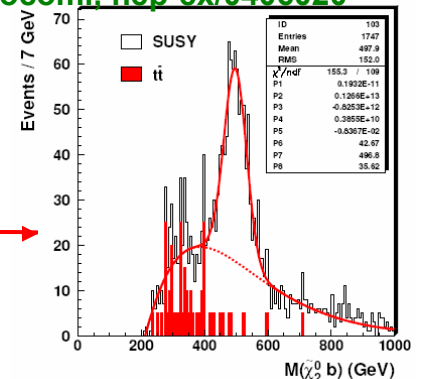
## $\tau$ 's

- identification of  $\tau$ 's in leptonic decays difficult because of  $E_{T, \text{miss}}$ 
  - use hadronic decays, but high background
- polarisation measurement appears possible

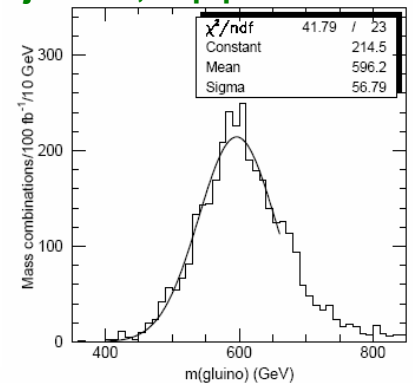
$$\chi_2^0 \rightarrow \tau \tilde{g}$$

$$\rightarrow \tau \chi_1^0$$

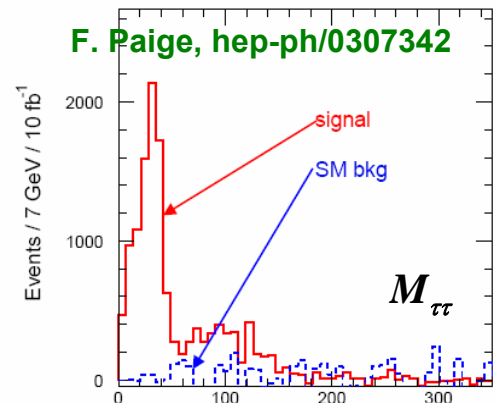
A. Trocomi, hep-ex/0406020



Nojiri et al, hep-ph/0312317



F. Paige, hep-ph/0307342



*mixing angles and masses possible by combining LHC and ILC data*

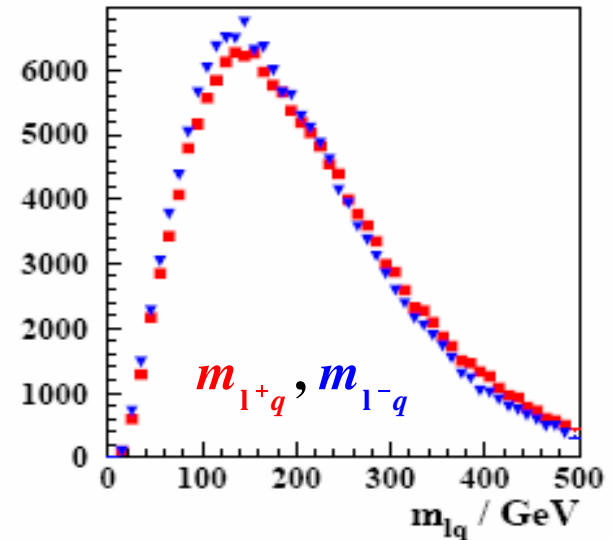
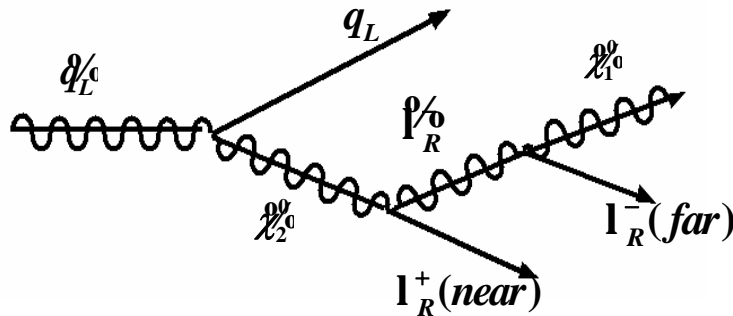
*heavy sfermions: → see talk by G. Moortgat-Pick*



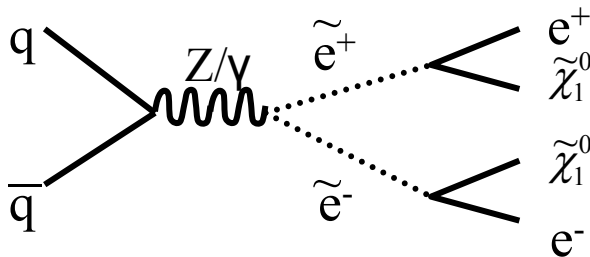
## Spin measurements

A.J. Barr, hep-ph/0405052

- evidence for supersymmetry (vs extra dimensions, for example)
  - polarization of  $\tilde{\chi}_2^0$  induces asymmetry
  - seems feasible, with  $150 \text{ fb}^{-1}$ , but very difficult



- slepton spin can be measured from angular distribution:  $1 - \cos^2 \theta$  for SUSY



*polarization measurements with ILC and threshold measurements*

A. Barr, ATLAS presentation, Oct.2005

## ➤ GMSB

- $LSP : \tilde{G}^0$ ;  $NLSP : \tilde{\chi}_1^0$  or  $\tilde{\nu}(\tilde{\nu}/\tilde{\rho})$
- jets +  $E_t^{\text{miss}}$ : ~ similar reach as MSugra
- signatures:
  - $\tilde{\chi}_1^0 \rightarrow \tilde{G}^0 \tau^+ \tau^-$  for short lifetimes, or
  - $\tilde{\chi}_1^0 \rightarrow \tilde{G}^0 \gamma$  : non-pointing photons for long lifetimes
  - $\tilde{\nu} \approx$  slow “muon”  $\Rightarrow$  TOF measurements
  - $\tilde{\nu} \rightarrow \tau \tilde{\chi}_1^0 \rightarrow \tau \gamma \tilde{G}^0$  : non-pointing photons

Kawagoe, Vienna 2004

*ILC: help discriminate between scenarios extrapolation to GUT scale*

## ➤ AMSB

- light  $\tilde{W}^0$  or  $\tilde{\nu}$   $p\bar{p} \rightarrow \tilde{\nu}_L \tilde{p}_L^+ \rightarrow \tau^+ \tilde{W}^0 \tau^+ \tilde{W}^0$ ; displaced vtx from  $\tilde{W}^0 \rightarrow \tilde{W}^0 \pi^\pm$
- $\tilde{\nu}$  pairs:

## ➤ Generic SUSY search by statistical method

Duchovni, SN-ATLAS-2004-043

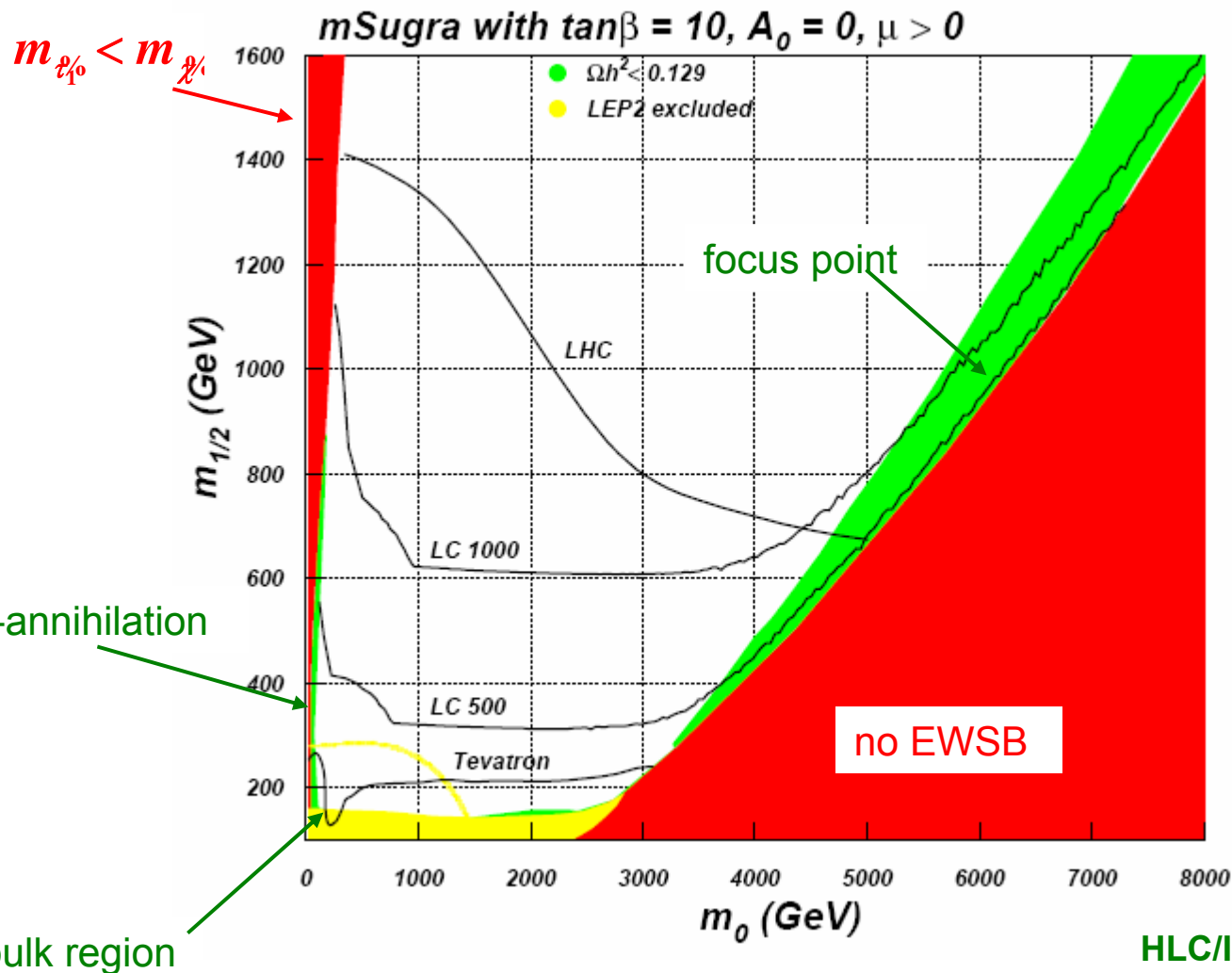
## ➤ R-parity violating processes

## ➤ Split Supersymmetry

- heavy, stable gluino  $\rightarrow$  R-hadron

## ➤ LFV ( $\tau - \mu$ sector)

# SUSY vs WMAP (and PLANCK)



see T. Lari, at Vienna 2004

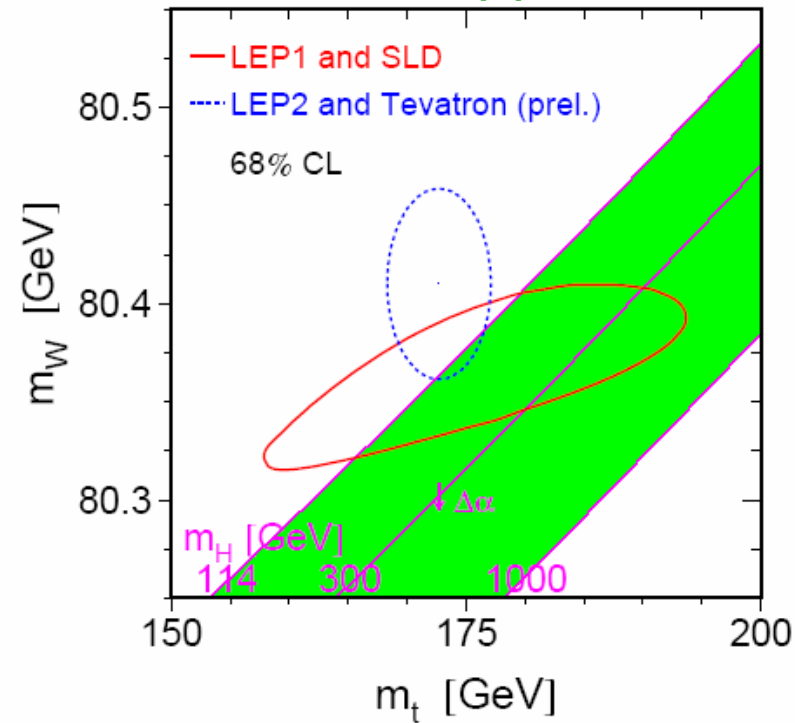
HLC/ILC workshop

*ILC: High precision measurement necessary to match experimental predictions from WMAP and PLANCK, assuming cosmological model*

## ➤ precision mass measurements:

- $m_W$ : use  $Z \rightarrow ee$  as a reference for an accurate measure of the em energy scale (N. Besson, M. Boonekamp, ATL-PHYS-2006-007)  $\Rightarrow$  syst. uncertainty of 4 MeV
- $m_t$ : from  $t\bar{t}$ ; various decay channels investigated. di-lepton channel clean, but more than one solution and MC weights required
- $V_{tb}$  from single top measurement

LEP EWWG, hep-ph/0511027

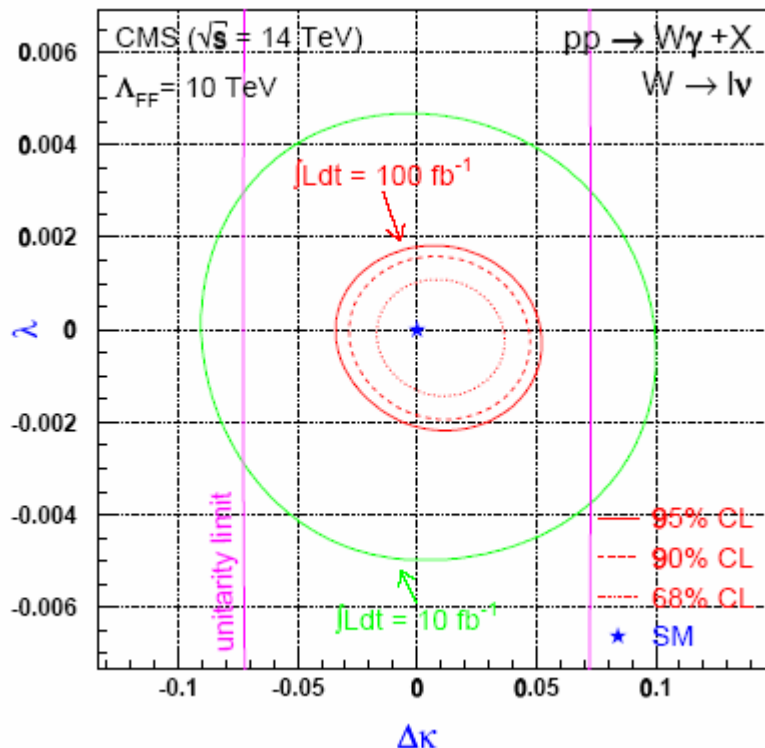


## LHC/ILC report (Table 6.5)

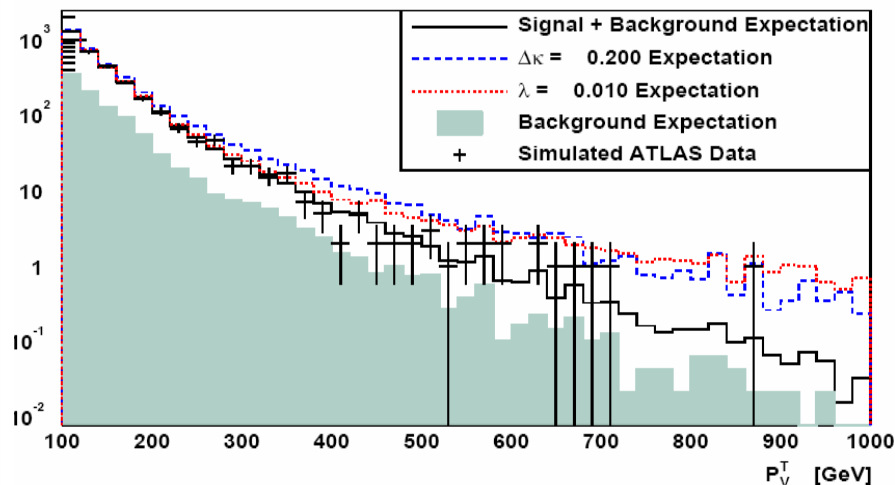
	now	Tev. Run IIA	Run IIB	LHC	LC	GigaZ
$\delta \sin^2 \theta_{\text{eff}} (\times 10^5)$	17	78	29	14–20	(6)	1.3
$\delta M_W$ [MeV]	34	27	16	15	10	7
$\delta m_t$ [GeV]	5.1	2.7	1.4	1.0	0.2–0.1	0.1
$\delta m_h$ [MeV]	—	—	$\mathcal{O}(2000)$	200	50	50

# Triple Gauge Boson Couplings

S. Haywood, CERN 2000-004



M. Dobbs, hep-ph/0506174



*ILC: considerable improvement for all parameters:  $\Delta\kappa_\gamma, \Delta\lambda_\gamma, \Delta\kappa_Z, \Delta\lambda_Z, \Delta g_Z$*   
*(LHC/ILC report, Sect. 6.2)*

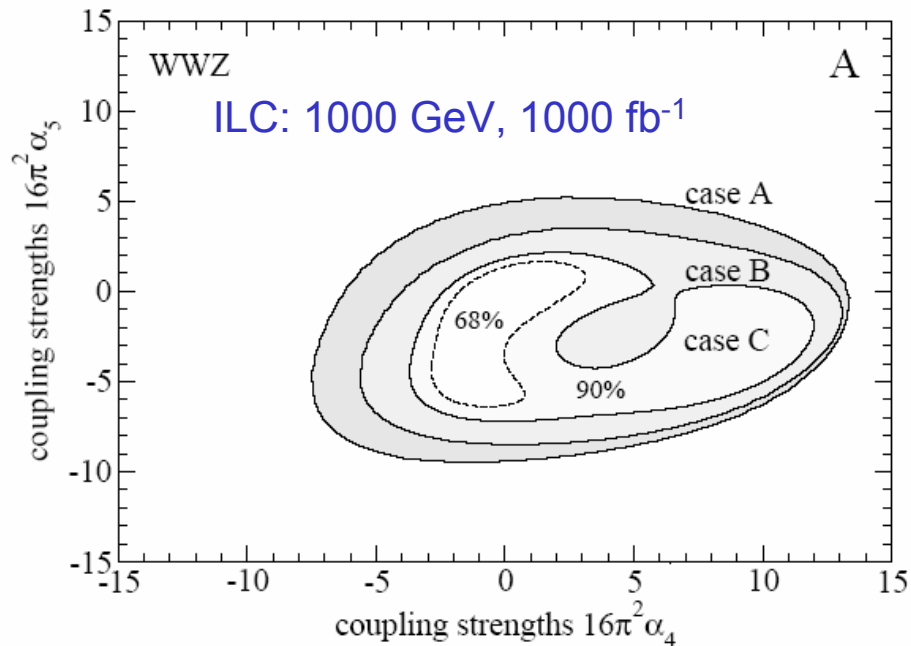
# sensitivity to quartic couplings

high mass vector boson scattering.

at LHC: c.m. energy sufficient to explore resonances (unitarization procedure)

→ see talk by T. Barklow

M. Beyer, hep-ph/0604048

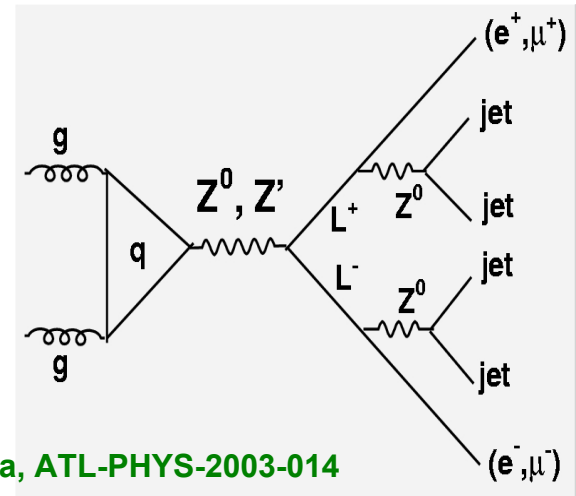


$$\begin{aligned}\mathcal{L}_4 &= \alpha_4 [\text{Tr}(V_\mu V_\nu)]^2, \\ \mathcal{L}_5 &= \alpha_5 [\text{Tr}(V_\mu V^\mu)]^2, \\ \mathcal{L}_6 &= \alpha_6 \text{Tr}(V_\mu V_\nu) \text{Tr}(TV^\mu) \text{Tr}(TV^\nu) \\ \mathcal{L}_7 &= \alpha_7 \text{Tr}(V_\mu V^\mu) [\text{Tr}(TV^\nu)]^2, \\ \mathcal{L}_{10} &= \frac{\alpha_{10}}{2} [\text{Tr}(TV_\mu) \text{Tr}(TV_\nu)]^2.\end{aligned}$$

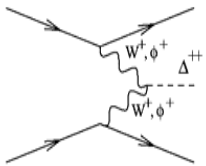
# Gauge theories with Extended symmetries

## ➔ E6 and other grand unified theories

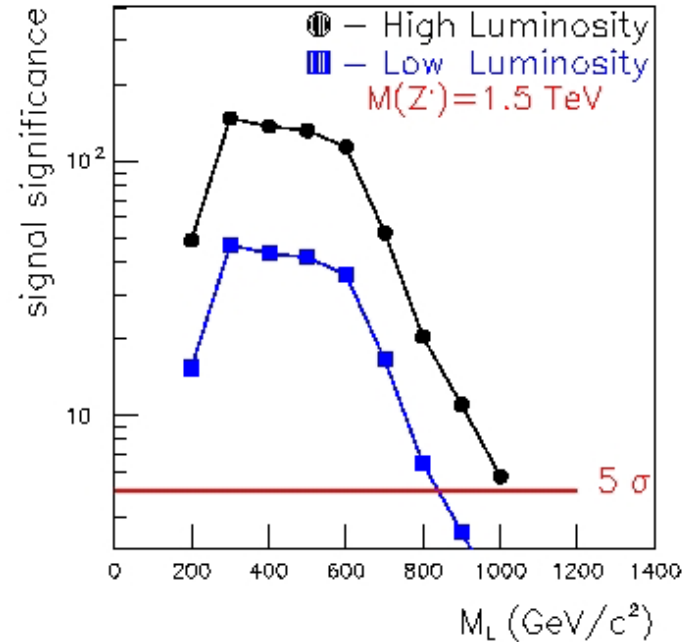
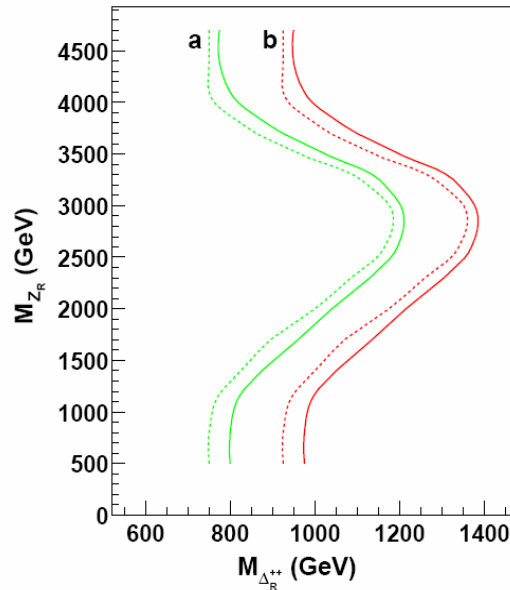
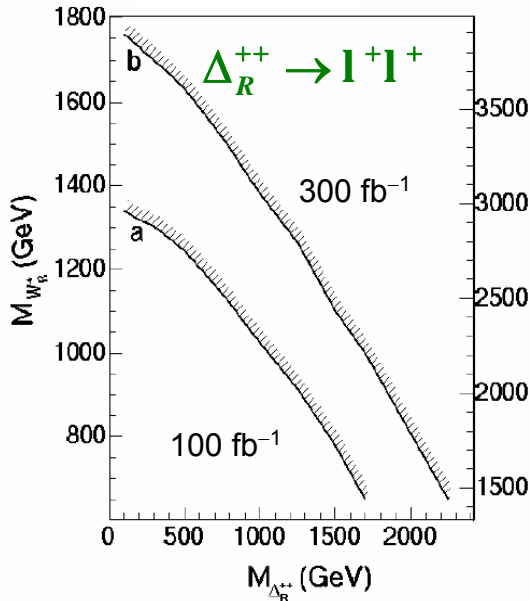
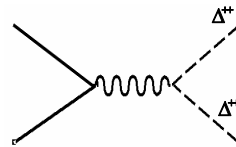
- include LR symmetric model
  - $Z', W'$
  - heavy leptons, Majorana neutrino
  - triplet Higgs ( $H^{++}$ )



C. Alexa, ATL-PHYS-2003-014

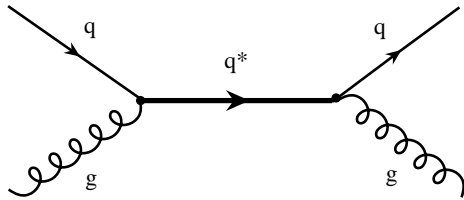


ATL-PHYS-2004-025





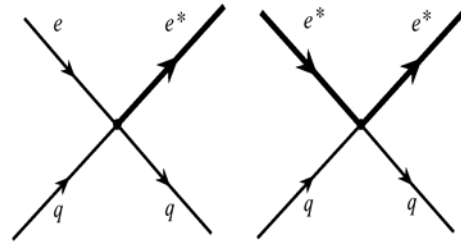
# Excited quarks and leptons



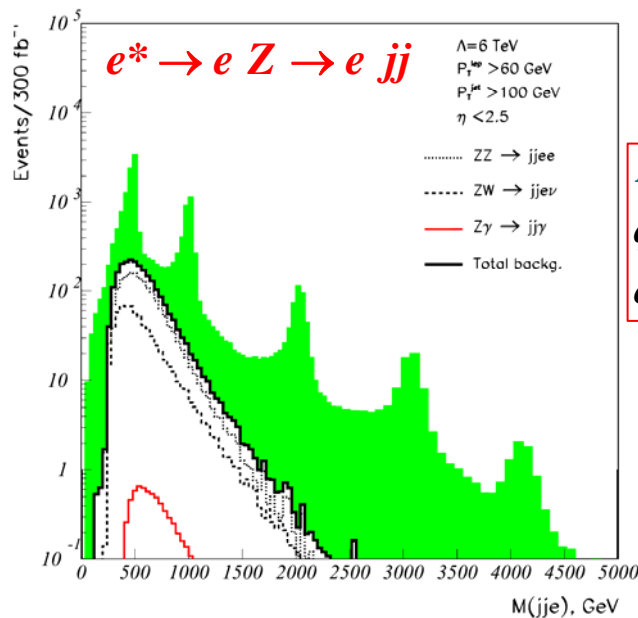
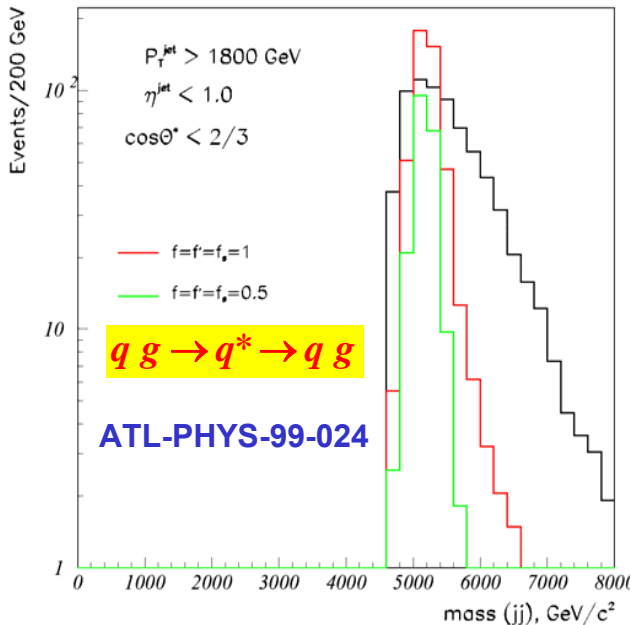
$$L = \frac{1}{2\Lambda} \bar{q}_R^* \sigma^{\mu\nu} \left( g_s f_s G_{\mu\nu}^a + g f \frac{\tau}{2} W_{\mu\nu} + g' f' \frac{Y}{2} B_{\mu\nu} \right) q_L + h.c.$$

take as reference :  $\Lambda = m^*$ ,  $f_s = f = f' = 1$

O. Çakir, C. Leroy, R. Mehdiev,  
ATL-PHYS-2002-014



contact interaction :  $L_C = \frac{g_*^2}{2\Lambda^2} J_\mu J^\mu$



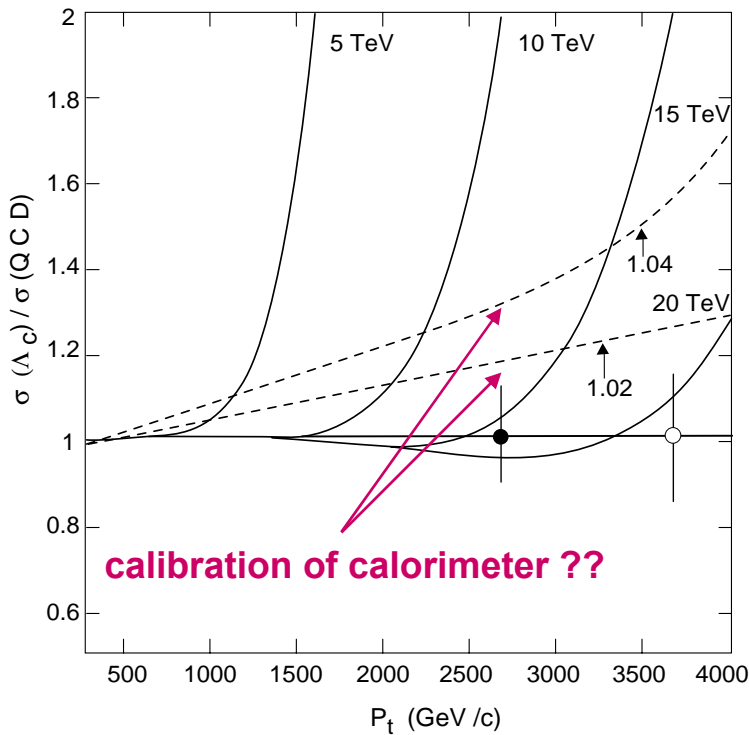
Also :

$q^* \rightarrow q\gamma$ ;  $q^* \rightarrow qZ$ ;  $q^* \rightarrow \bar{q}'W$   
 $e^* \rightarrow e\gamma$

clean signals up to few TeV, depending on  $\Lambda$

# quark substructure

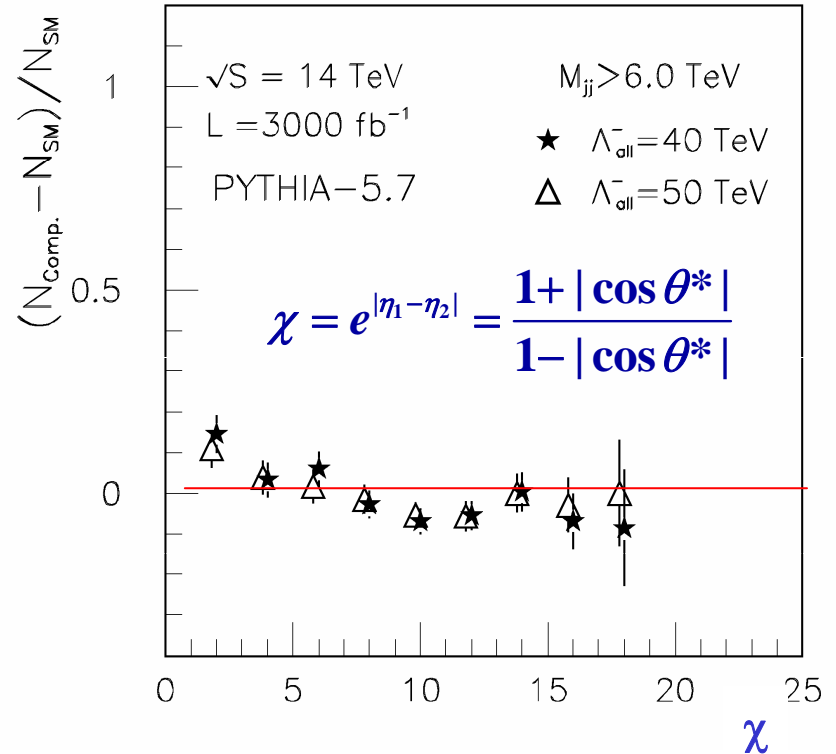
Contact interactions:  
scale of compositeness  $\Lambda \gg m^*$



di-jets, relative to SM QCD, from TP

extraD signal?

- virtual G exchange,
- virtual  $g^*$  exchange in UED?
- early gauge coupling unification?



angular distribution

less dependent on calibration

Z. Usubov

(from superLHC)

# Extra dimensions

## ➔ ADD: Large, flat compactified extra dimensions

- Gauss' law: →
- SM particles on 3-brane

$$G_N \frac{m_1 m_2}{r^2} = G' \frac{m_1 m_2}{r^2 (2\pi r)^n}, \text{ for } r = R \Rightarrow G' = G_N (2\pi R)^n$$

$$M_D^{2+n} = \frac{M_{Pl}^2}{8\pi (2\pi R)^n} \Rightarrow M_D : \text{TeV for values of } R \leq \text{mm}$$

for compactification in circles: graviton field periodic in extra dimensions ( $y$ )

$$\phi(x, y) = \sum_{k_1=-\infty}^{\infty} \sum_{k_2=-\infty}^{\infty} \dots \sum_{k_n=-\infty}^{\infty} \phi^{(k)}(x) e^{i \vec{k} \cdot \vec{y} / R}$$

towers of Kaluza-Klein states of graviton  
with  $p_T \sim \text{mass} = k/R$

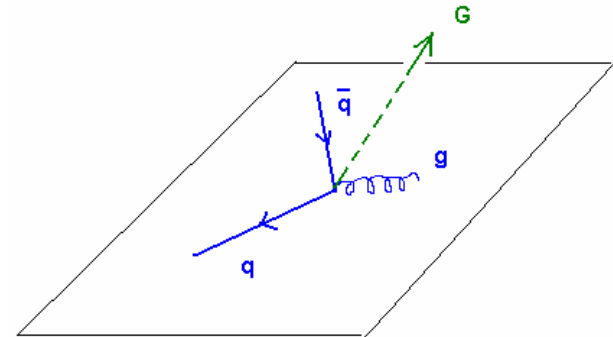
KK mode separation is very small: → continuous spectrum

→ high density of states compensates low coupling ( $\sim 1/M_{Pl}$ )

→ chance to observe effects at LHC

at LHC:

$$\left. \begin{aligned} \bar{q}q &\rightarrow gG^{(k)}, \gamma G^{(k)} \\ qg &\rightarrow qG^{(k)} \\ gg &\rightarrow gG^{(k)} \end{aligned} \right\} \text{jets} + \cancel{E}_T, \gamma + \cancel{E}_T$$



# Direct graviton production: Jets + missing $E_T$

## Detector effects: ATLFAST

- Jets and leptons reconstructed in  $|\eta| < 5, 2.5$

Minimum of validity



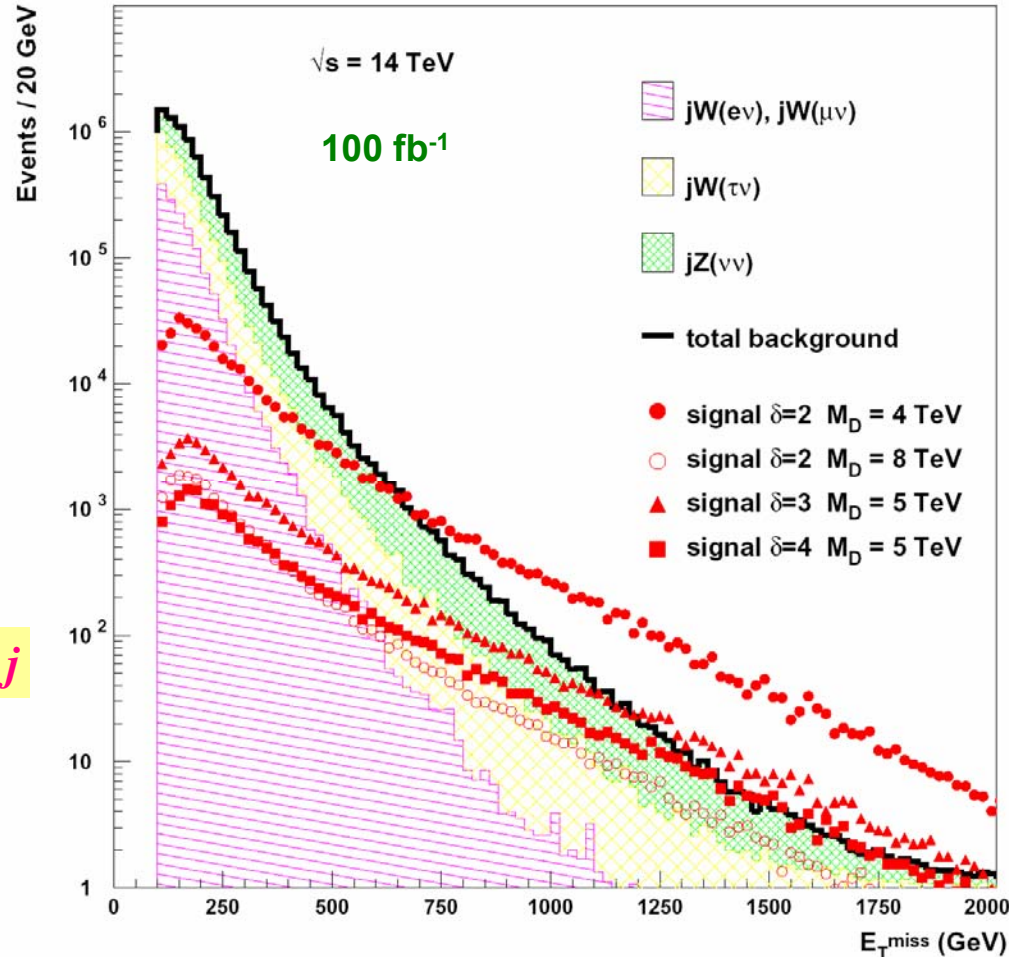
$\delta$	$M_D^{max}$ (TeV) LL, $30 \text{ fb}^{-1}$	$M_D^{max}$ (TeV) HL, $100 \text{ fb}^{-1}$	$M_D^{min}$ (TeV)
2	7.7	9.1	$\sim 4$
3	6.2	7.0	$\sim 4.5$
4	5.2	6.0	$\sim 5$

calibrate background level by  $Z + j \rightarrow l^+ l^- + j$

Uncertainty in  $\sigma(Z+jets)$  will lower the reach

Reach in  $M_D$  for  $\gamma G$

$\delta$	$M_D^{max}$ (TeV) HL, $100 \text{ fb}^{-1}$	$M_D^{min}$ (TeV)
2	4	$\sim 3.5$



ILC: disentangle  $\delta$  and  $M_D$  by running at 2 cm energies

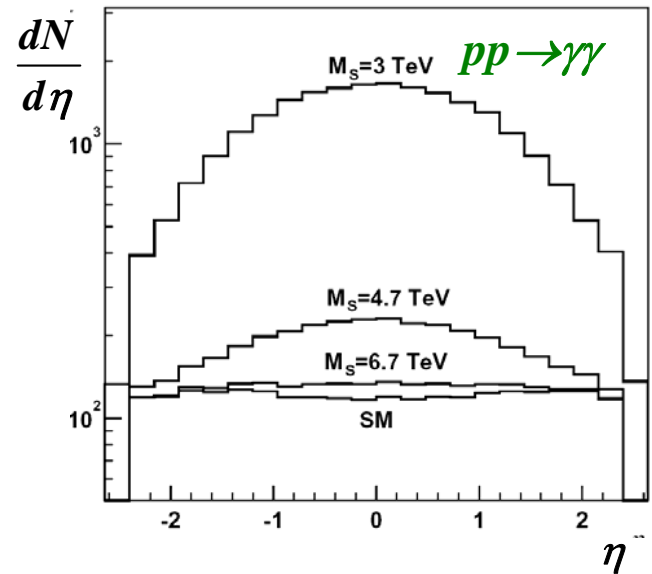
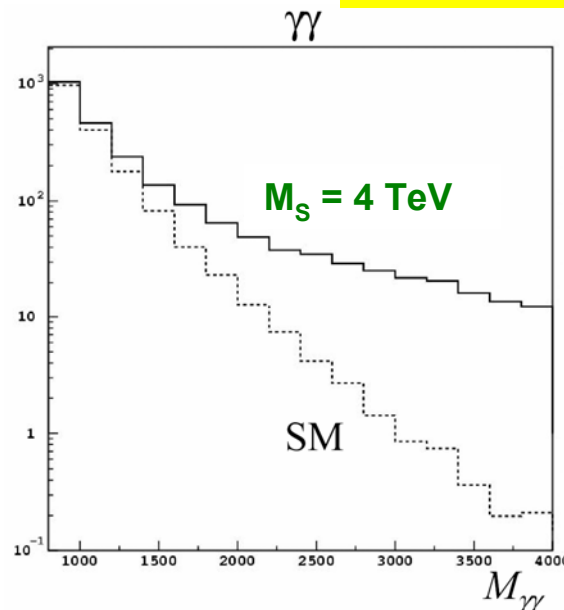
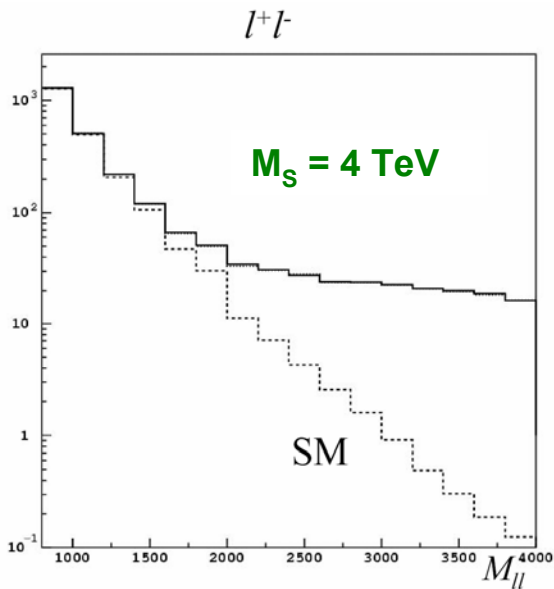
L. Vacavant and I. Hinchliffe, *J. Phys. G: Nucl. Part. Phys.* 27 (2001) 1839

# Virtual Graviton Exchange

➤ **Signals:**  $q\bar{q}, gg \rightarrow \gamma\gamma, l\bar{l}, (WW, t\bar{t} \dots)$

- We require that :  $M_{\gamma\gamma, l\bar{l}} < 0.9 M_S$  (effective theory)
- Excess in **di-leptons** mass distribution (same for di-photons)
- event distribution of  $\gamma\gamma$  (s-channel) more central than in SM (t and u channels)
- can measure FB asymmetry

**Sensitivity for  $100 \text{ fb}^{-1} \sim 5\text{-}6 \text{ TeV}$**



[detail](#)

# Extra dimensions of $\sim \text{TeV}^{-1}$ size

G. A. and G. Polesello *Proceedings of Les Houches 2001*

- Compactification radius is small enough to allow SM particles in the bulk

note:  $R = hc/1 \text{ TeV}^{-1} = 2 \times 10^{-4} \text{ fm}$  [ $M_D : 10^{15} \text{ GeV}$  for  $n=1$ ]

- **indirect** constraints from LEP EW measurements:

- limits from gauge boson couplings, mixings and virtual exchanges

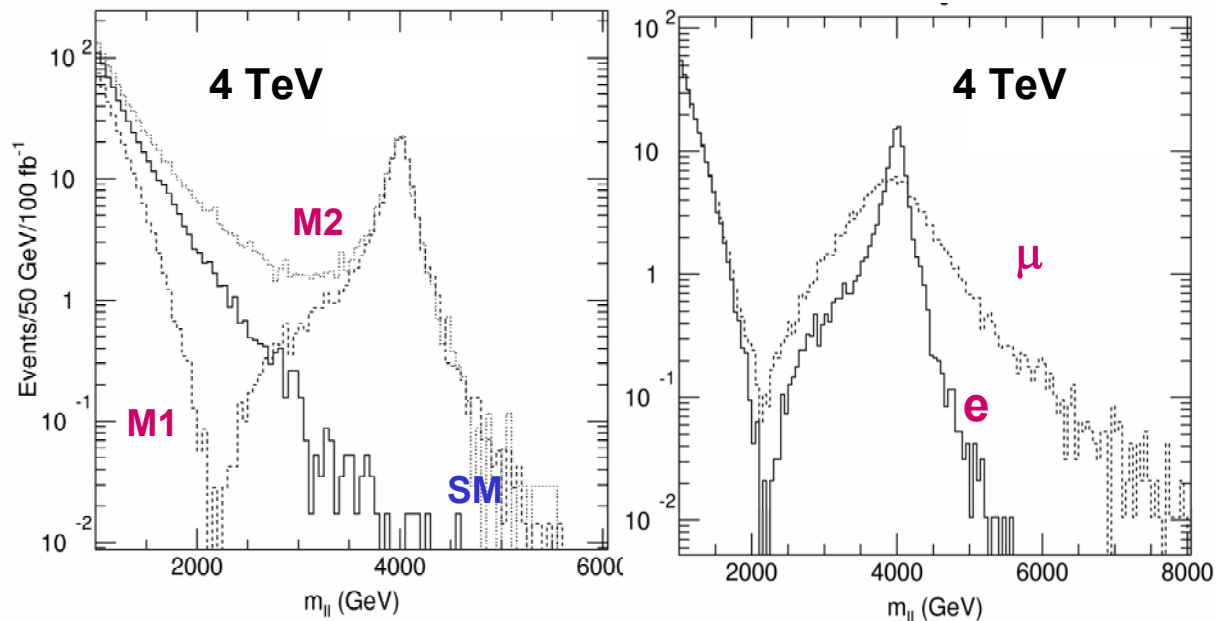
◦  $R^{-1} \geq 3.3 - 6.8 \text{ TeV}$  K. Cheung and G. Landsberg, Phys.Rev. D65 (2002) 076003 (*hep-ph/0110346*)

- **Z/ $\gamma$  excitation:**

( following T. Rizzo, Phys. Rev. D61 (2000) 055005)

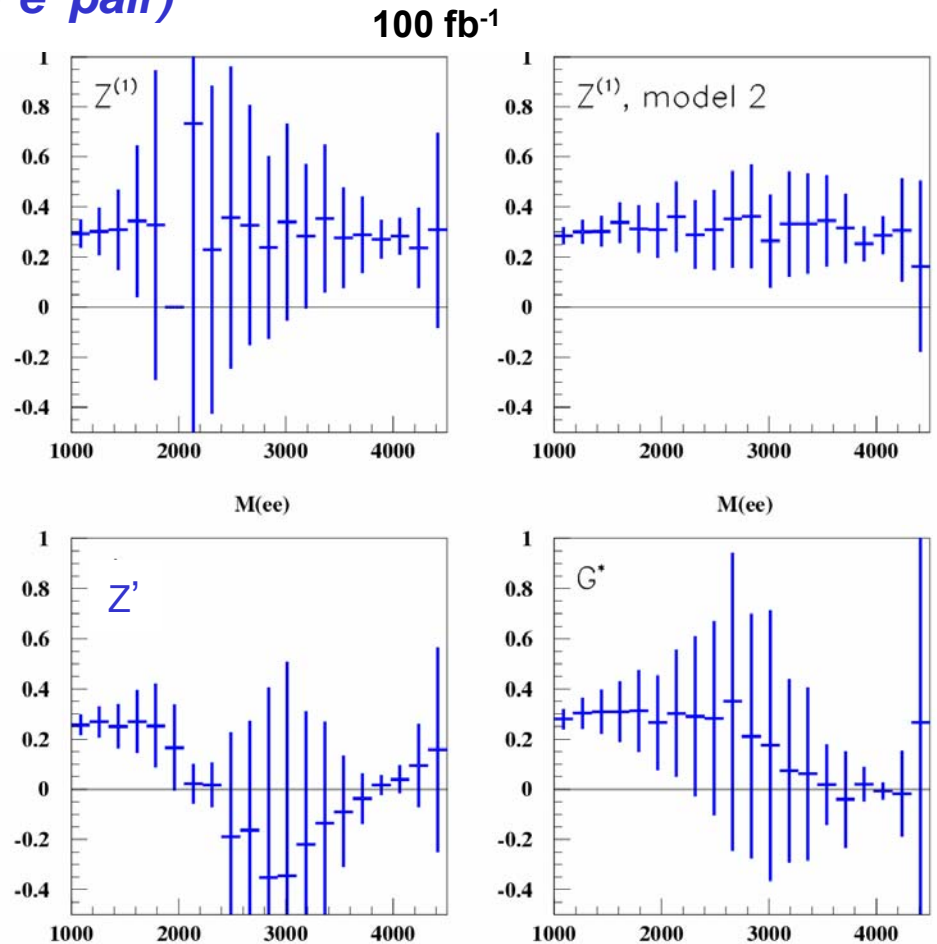
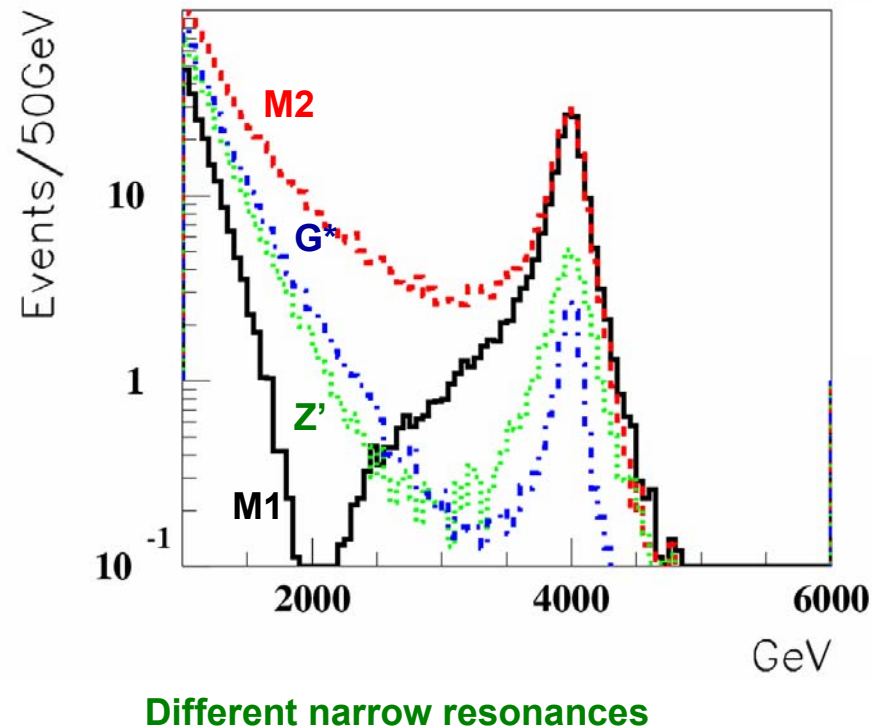
resonance observable up to  $\sim 5.8 \text{ TeV}$  with  $100 \text{ fb}^{-1}$

sensitivity to DY tail up to  $13.5 \text{ TeV}$  (e- $\mu$ ) combined



# Extra dimensions of $\sim\text{TeV}^{-1}$ size: $Z/\gamma$ excitation

- angular distribution and **FB asymmetry**  
(with respect to direction of  $e^+e^-$  pair)



*ILC: precision measurements of FB asymmetry and DY help discriminate between models*

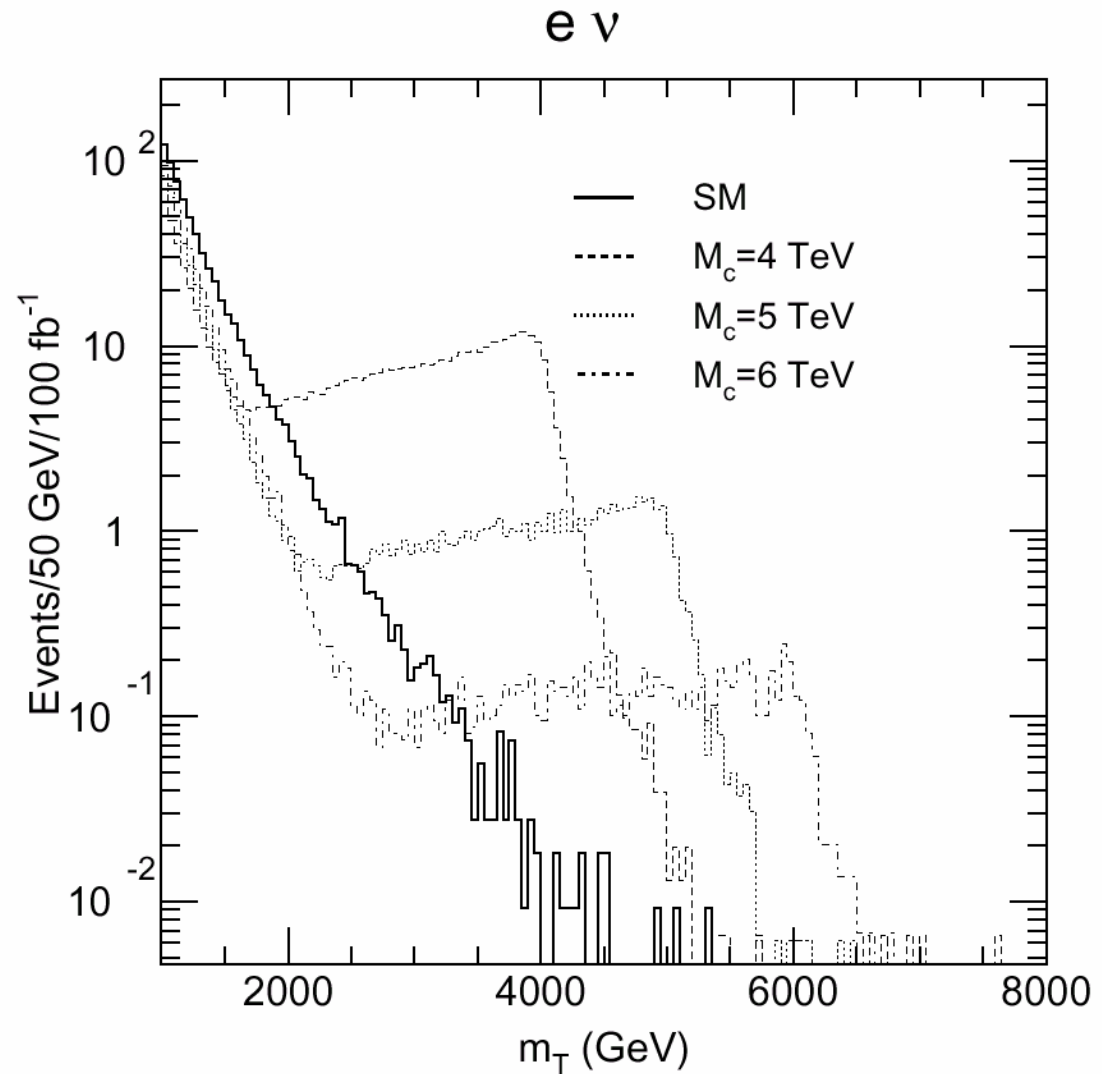


# Extra dimensions of $\sim \text{TeV}^{-1}$ size: $W$ excitation

## ➔ Same for $W^{(1)}$

- direct observation  
up to 6 TeV
- indirect  
up to  $\sim 9$  TeV

## ➔ more difficult to distinguish from a $W'$



G. Polesello et M. Prata  
*EPJDirect (SN-ATLAS-2003-036)*

# Randall-Sundrum Model

L. Randall and R. Sundrum, Phys.Rev.Lett. 83 (1999) 3370 (hep-ph/9905221)

## ➔ 1 extra dimension $y$ with non-factorizable metric

- 5D space, bounded by 2 branes
  - SM brane (TeV) at  $y = \pi r_c$
  - Planck brane at  $y = 0$

$$ds^2 = e^{-2ky} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2, \quad y = r_c \phi, \quad (k \sim M_{Pl}) \Rightarrow \text{3D distances shrink with } y$$

- 5-D Planck scale  $M_5 \sim M_{Pl}$ :  $M_{Pl}^2 = \frac{M_5^3}{k} (1 - e^{-2kr_c\pi})$  curvature scalar  $|R_5| = 20k^2$
- new physics scale at SM brane =  $\Lambda_\pi = M_{pl} e^{-kr_c\pi}$ ;  $kr_c\pi : 35 \Rightarrow \Lambda_\pi : \text{TeV}$ 
  - coupling of KK states  $\sim 1/\Lambda_\pi$
- KK masses:
  - $m_n = kx_n e^{-k\pi r_c}$ , with  $J_1(x_n) = 0 \Rightarrow m_1 = 3.83 \frac{k}{M_{Pl}} \Lambda_\pi$

3 parameters:  $k, r_c, \Lambda_\pi$   
(2 independent)

constraints:  $0.01 \leq \frac{k}{M_{Pl}} \leq 0.1$

*Note: astrophysical bounds not applicable in this model*

# Narrow graviton resonance

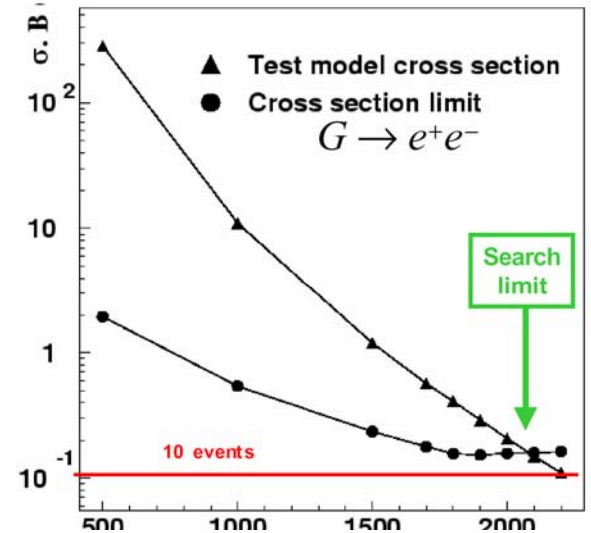
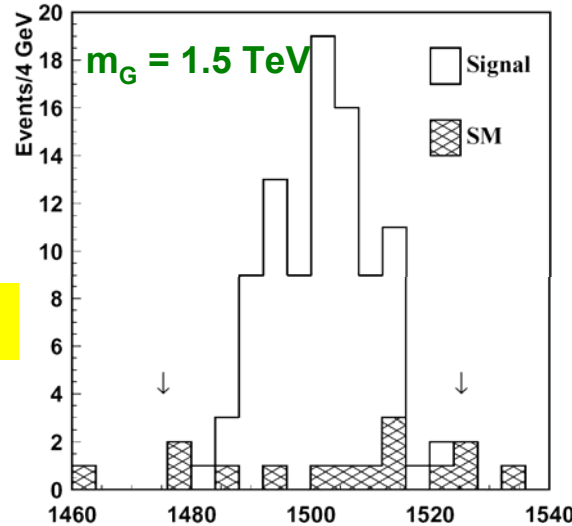
**$e^+e^-$  channel,**

**$100 \text{ fb}^{-1}$**

$$\frac{k}{M_{Pl}} = 0.01$$

**Sensitive at  $5\sigma$  up to 2080 GeV**

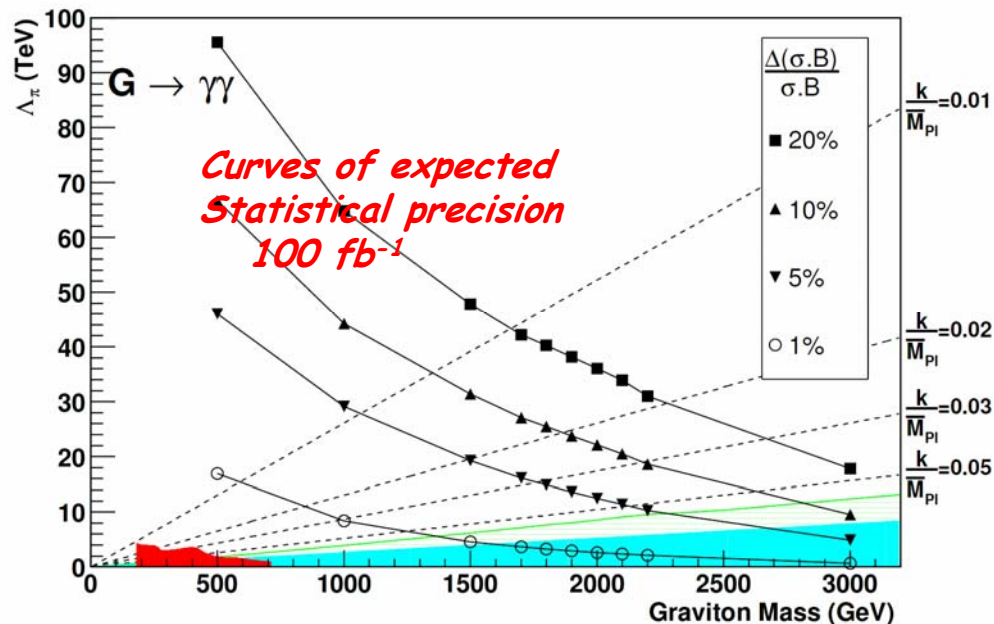
**B. C. Allanach, K. Odagiri, A. Parker and B. Webber,  
JHEP 9 (2000) 19**



*- also  $G \rightarrow WW, ZZ, jj, \mu\mu, tt, hh$*

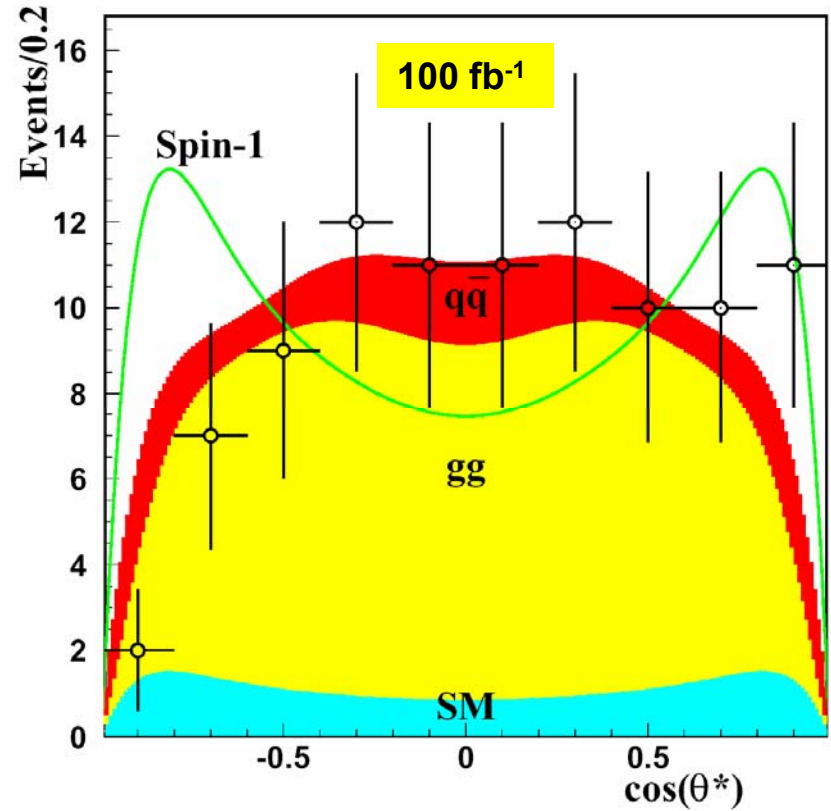
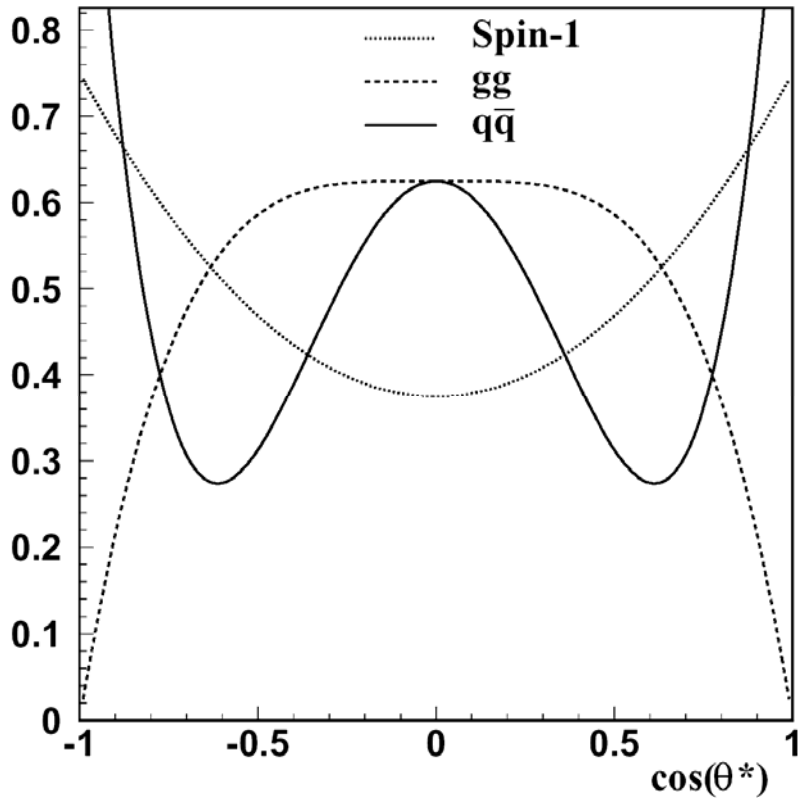
**e.g.:** for a resonance observed at  
 $m_G = 1.5 \text{ TeV}$  in  $ee$  channel  
 $\Delta m_G < 10.5 \text{ GeV}$  (energy scale error)  
 $\Delta \sigma.B \sim 18\%$   
 if  $k/r_c = 0.01$  (pessimistic)  
 $\Rightarrow r_c = (82 \pm 7) \times 10^{-33} \text{ m} !!$

**Allanach et al., ATL-PHYS-2002-031**



# Narrow Graviton Resonance

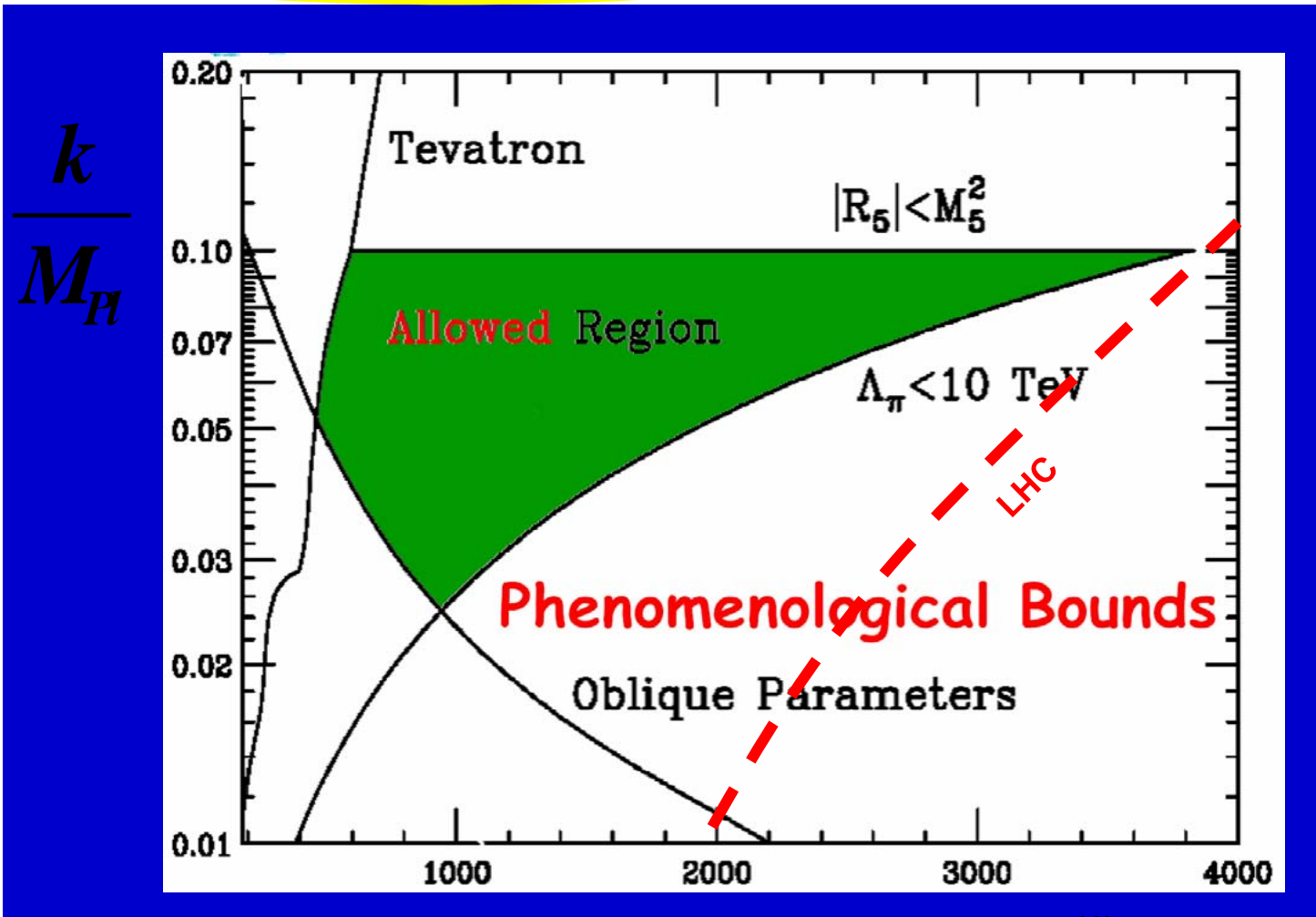
## Spin determination



spin-2 could be determined (spin-1 ruled out) with 90% CL up to graviton mass of **1720 GeV**

# Narrow Graviton Resonance

Interesting region covered by LHC

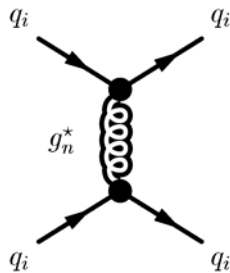


H. Davoudiasl, J. Hewett and T. Rizzo, Phys.Rev. D63 (2001) 075004 ([hep-ph/0006041](https://arxiv.org/abs/hep-ph/0006041))

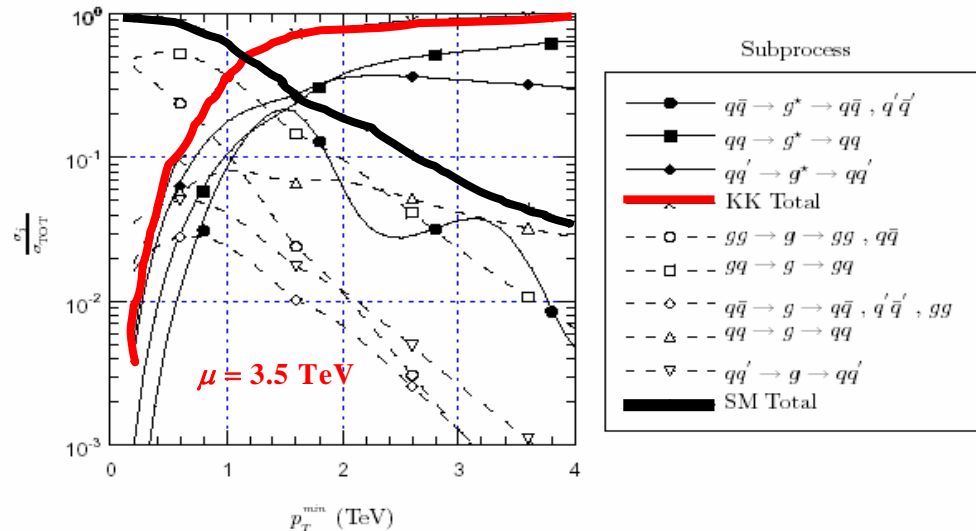
## TeV<sup>-1</sup>-size:

DA Dicus, CD McMullen and S. Nandi hep-ph/0012259

- virtual  $g^*$  excitation  $\Rightarrow$  enhanced di-jet cross section



ATLAS study in progress...



## Universal Extra dimensions

T. Appelquist, HC Cheng and BA Dobrescu, PR D64 (2001) 035002

- All SM particles in bulk

- $\Rightarrow$  conservation of momentum in extra dimensions
- $\Rightarrow$  conservation of KK number
- $\Rightarrow$  pair production of KK states
- $\Rightarrow$  lower collider bounds:  $\sim 350$ - $400$  GeV

# Universal Extra Dimensions

## ➤ dijet signals

- stable
- unstable:  
fat brane absorbs unbalanced momentum from KK number violation

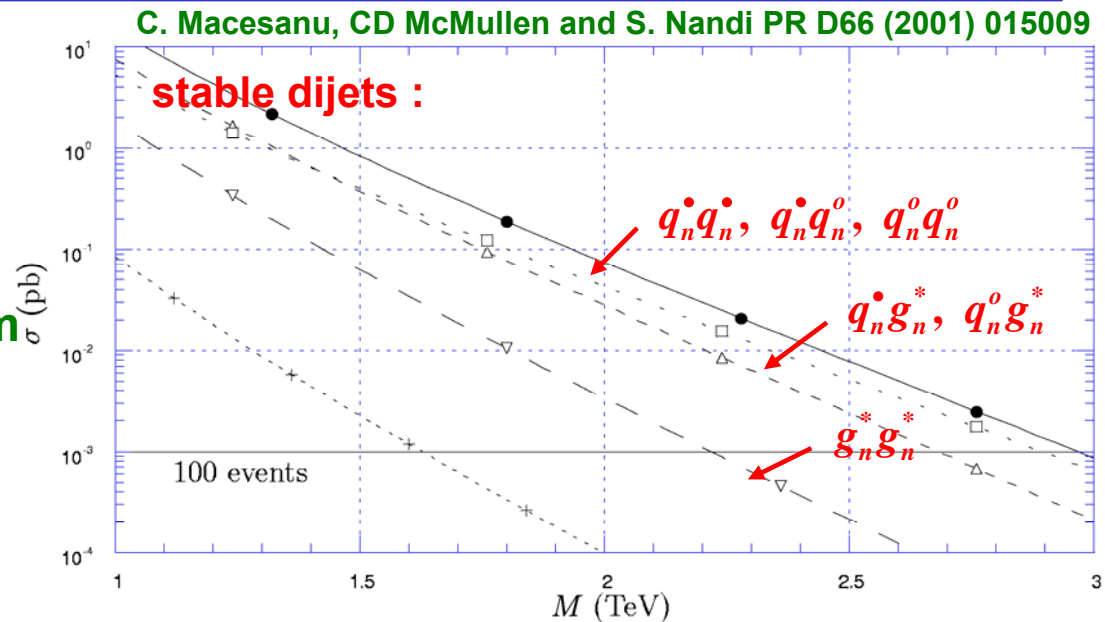
## ➤ LKP: $\gamma^*$

$$q_i q_i \rightarrow (\dots + \gamma^*) (\dots + \gamma^*) (\rightarrow \gamma G + \gamma G + X)$$

C. Macesanu, CD McMullen and S. Nandi  
Phys.Lett. B546 (2002) 253

- can be fooled by SUSY

HC Cheng, KT Matchev, and M Schmaltz  
Phys.Rev. D66 (2002) 056006 (hep-ph/0205314)



*ILC: spin can distinguish between SUSY and UED better than at LHC*

## ➤ QCD:

- measurement of  $\alpha_s$  through high  $E_T$  jets (*ILC: use event shapes*)
- parton densities
- tests of NNLO corrections
- forward physics

## ➤ leptoquarks

- reach:  $\sim 1.5$  TeV

## ➤ scalars in the bulk

- $\rightarrow$  coupling to Higgs
- mixing with Higgs

## ➤ Lepton Flavor Violation:

- $\tau \rightarrow 3 \mu$
- in SUSY decays

## ➤ monopoles

## ➤ isosinglet quarks...

## ➤ black holes



- **LHC: a discovery machine with a very rich programme of physics**
  - At the start of ILC, we can expect**
    - some manifestation from source of EWSB, and resolution of hierarchy problem?  
(SM, SUSY, little Higgs? extra Dimensions?)
      - « guarantee » of some new phenomena to be observed
        - **ILC should refine understanding and help discriminate between models through clean precision measurements and new processes**
          - **telescope and microscope for much better measurement of couplings, consistency**
    - exotic physics
      - extended gauge groups?
      - extra dimensions?
      - compositeness?
      - 3<sup>rd</sup> family couplings?
- **Observation of a resonance or other signal is not enough. What is the nature of the new physics ??**
  - **ILC will allow to observe the new discoveries from a different angle and correlate precision measurements of different observables**

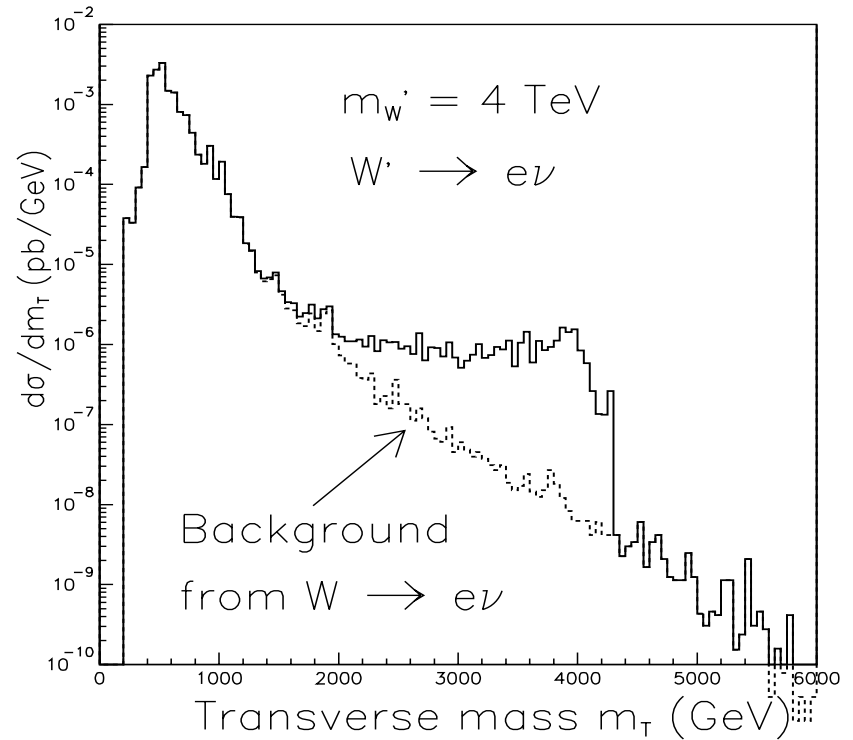
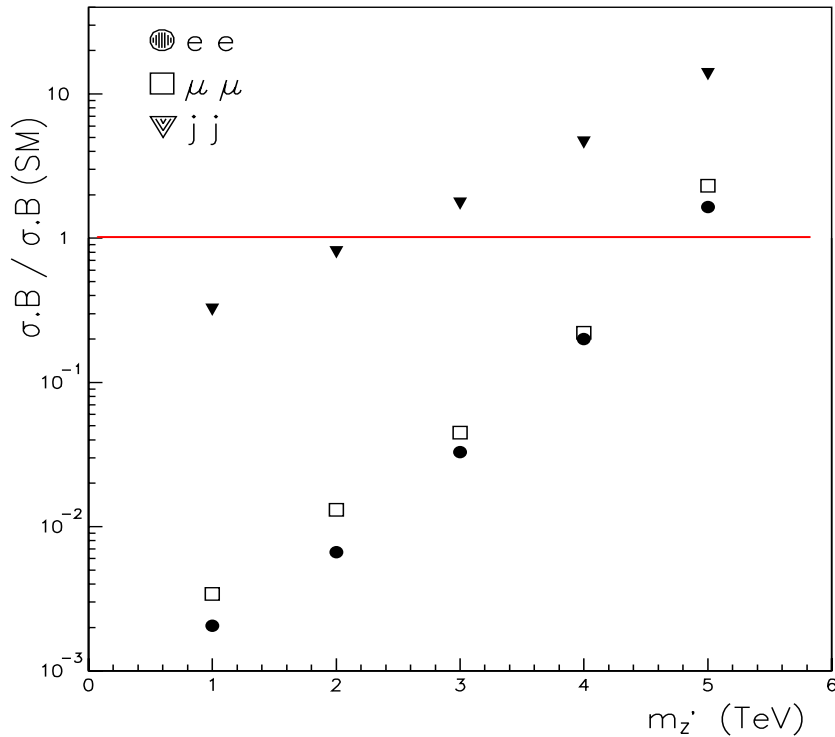
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**backup**

# Sequential Z', W'

## Generic Z', W'

*clean signals, little background*



from TDR  $100 \text{ fb}^{-1}$

### experimental considerations:

- energy calibration, resolution, ID, efficiencies for high energy e,  $\mu$ , jets
- charge assignment for e,  $\mu$

## Z' : first sign of extended gauge group ?

### ➤ Left-Right Symmetric Model:

$$\begin{array}{ccccc}
 SU(2)_L \times SU(2)_R \times U(1)_{B-L} & \xrightarrow{\Delta} & SU(2)_L \times U(1)_Y & & \\
 \downarrow & & \downarrow & & \downarrow \\
 W_L^i & W_R^i & C & W_L^i & B \\
 \mathfrak{g}_L & \mathfrak{g}_R & \mathfrak{g}' & \mathfrak{g}_L & \mathfrak{g}_Y
 \end{array}$$

triplet Higgs

- right-handed fermions in doublets → heavy Majorana  $\nu_R = N$ 
  - explains low mass of  $\nu_L$  (see-saw mechanism)
- $W_R, Z_R$  associated with right-handed sector  $W_R \rightarrow eN, Z_R \rightarrow ee$

### ➤ larger GUT groups (includes LRSM)

$$\begin{aligned}
 E_6 &\rightarrow SO(10) + U(1)_\psi \\
 &\quad \downarrow \\
 &\quad \rightarrow SU(5) + U(1)_x
 \end{aligned}$$

$$Q_{E_6} = \cos \beta Q_x + \sin \beta Q_\psi : \Rightarrow Q_\eta = \sqrt{\frac{3}{8}} Q_x - \sqrt{\frac{5}{8}} Q_\psi$$

# Z', W' in LRSM

$p p \rightarrow Z_R \rightarrow N_1 N_1 \rightarrow l j j l j j$

J. Collot, A. Ferrari  
ATL-PHYS-98-124,  
ATL-PHYS-99-034

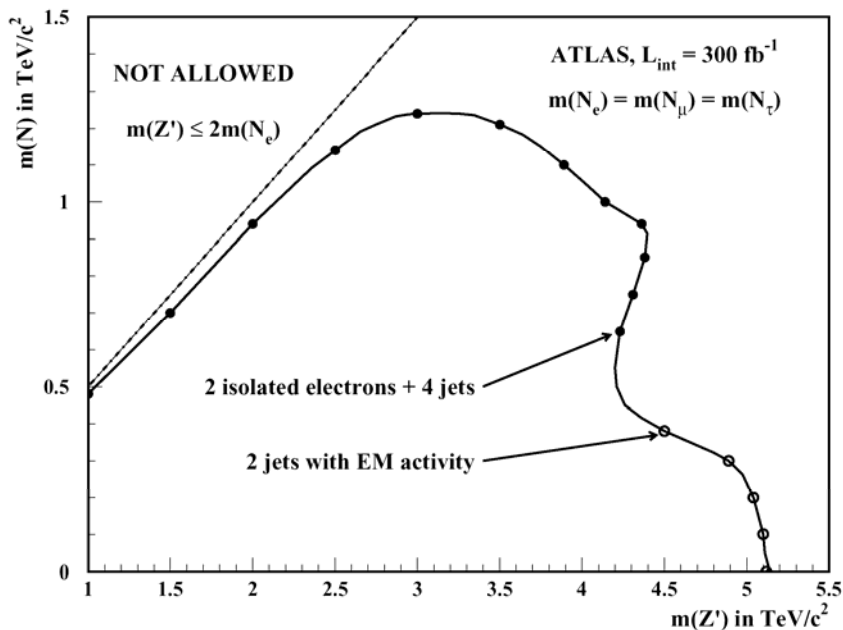
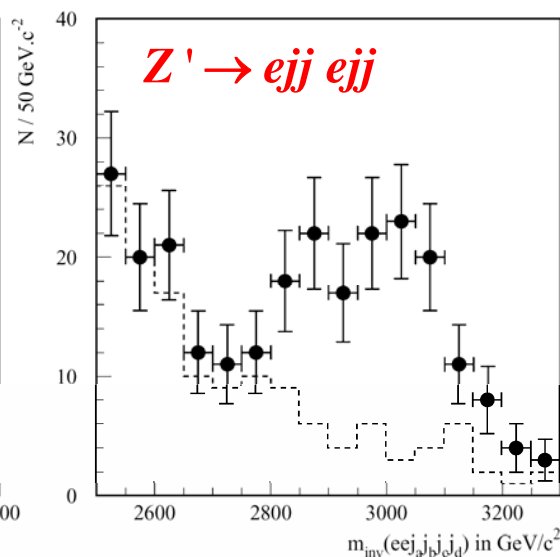
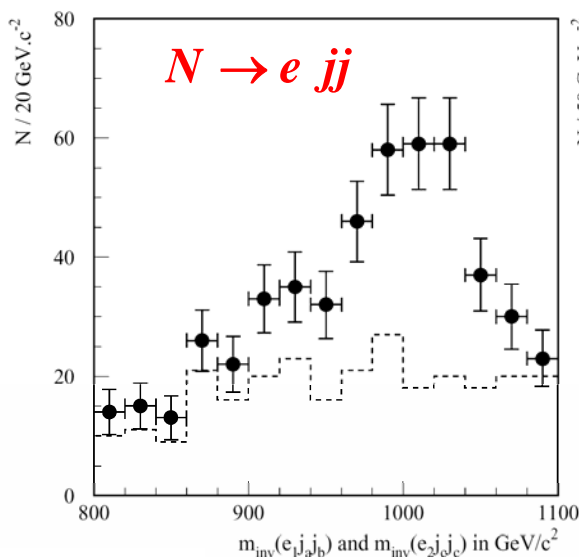
backgrounds:

$t \bar{t}$

$DY, WW, ZW, ZZ$

LRSM bckg:  $W_R, \dots$

$\Rightarrow$  cuts on  
 $m_{ee}, pT(\text{jets})$



FB asymmetry gives a  
measure of  $\kappa = g_R/g_L$

$$m_{Z'} = \sqrt{\frac{2\kappa^2 \cot^2 \theta_W}{\kappa^2 \cot^2 \theta_W - 1}} m_{W_R}$$

$= 1.7 m_{W_R}$  if  $\kappa=1$

# Z', W'

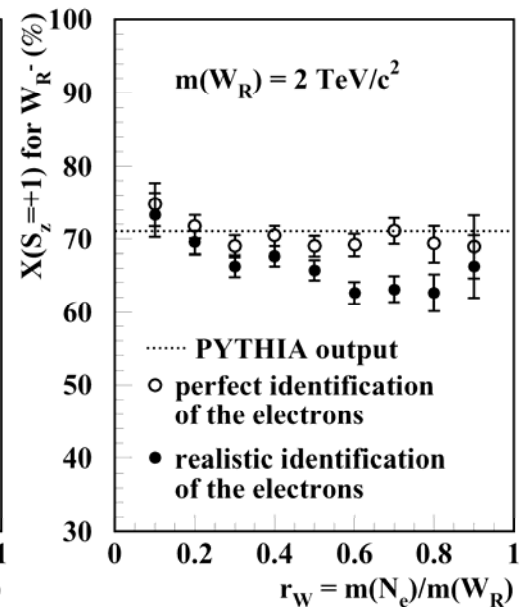
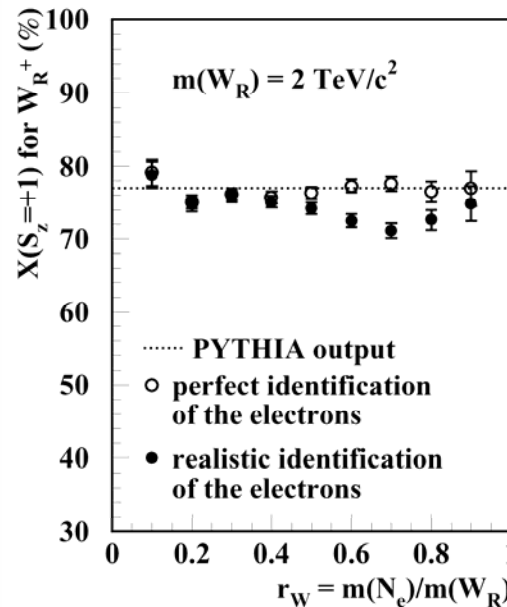
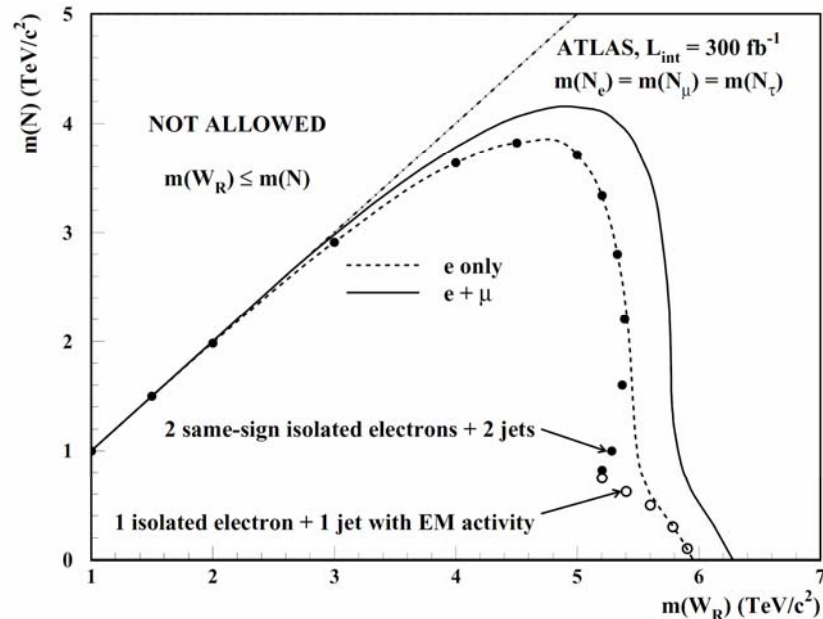
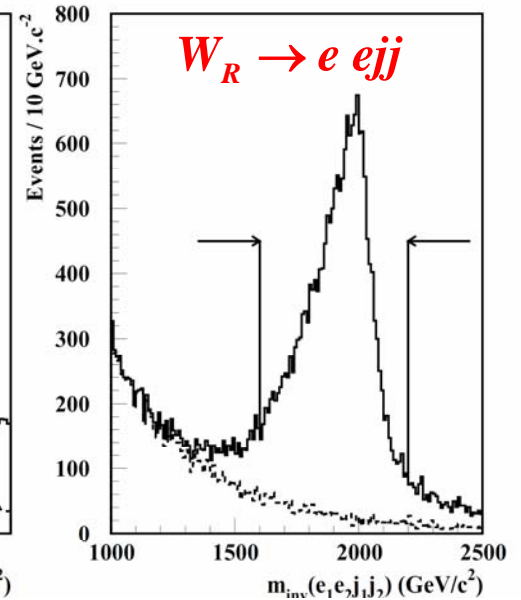
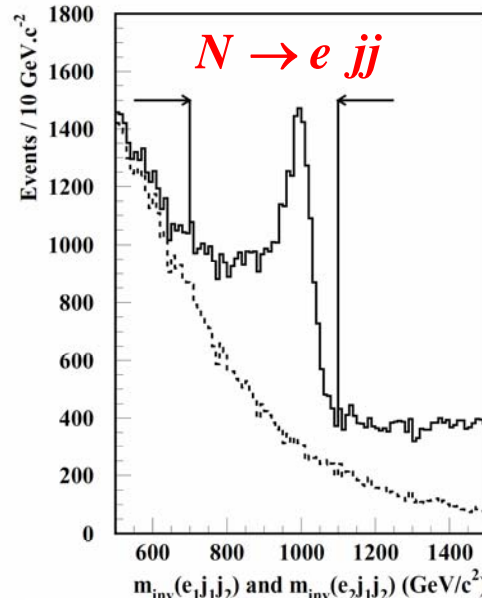
$$p p \rightarrow W_R \rightarrow l N_1 \rightarrow l l j j$$

J. Collot, A. Ferrari  
ATL-PHYS-98-124,  
ATL-PHYS-99-018

backgrounds:

*t tbar*

*DY, WW, ZW, ZZ*



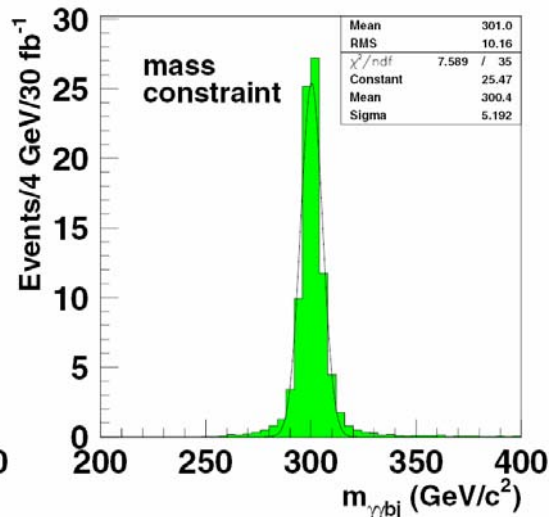
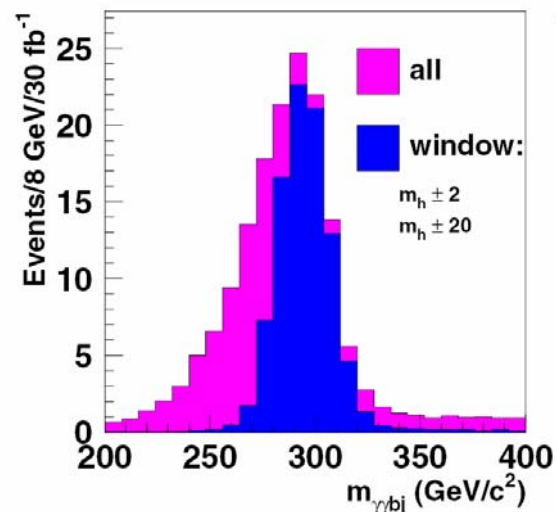
# Randall-Sundrum model: the radion

## signal: $\phi \rightarrow hh \rightarrow \gamma\gamma b\bar{b}$

- similar to MSSM, but with appropriate corrections for width and branching ratios
- consider cases:  $m_\phi = 300, 600$  GeV,  $m_h = 125$  GeV

## backgrounds negligible

- $gg \rightarrow \gamma\gamma$  with QCD radiation
- $\gamma j$ , with jet misidentified as photon

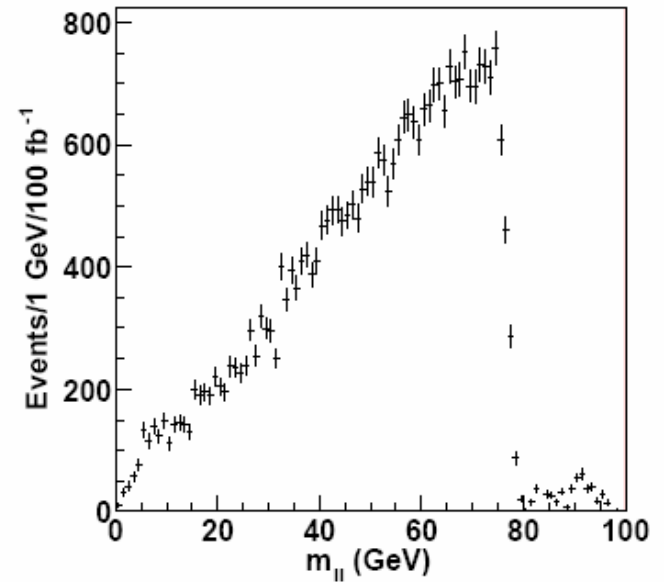
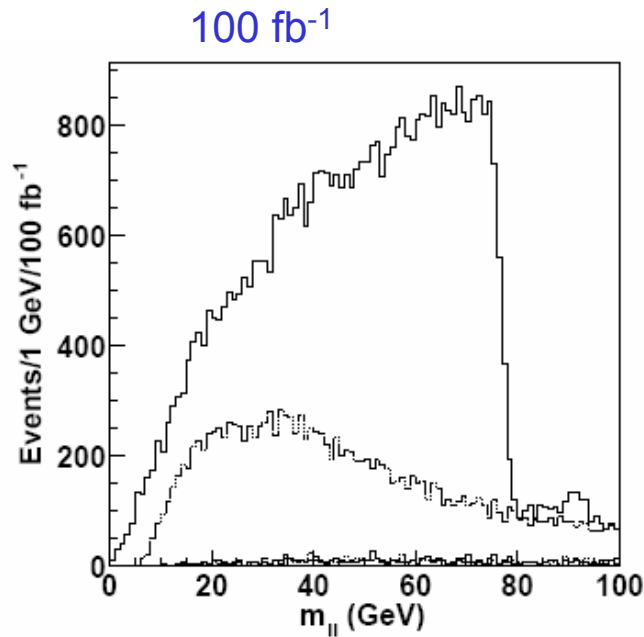


$\xi$	$\Lambda_\phi$ (TeV)	$m_\phi=300$ GeV	$m_\phi=600$ GeV
0	1	4	43
0	10	333	-
1/6	1	2	57
1/6	10	250	-

Luminosity ( $\text{fb}^{-1}$ ) required for  $5\sigma$  discovery

reach: 2.2 TeV or 0.6 TeV for  $m_\phi = 300$  or 600 GeV, respectively, with  $30 \text{ fb}^{-1}$

**SPS 1a :**  
 $m_0 = 100$  GeV,  
 $m_{\frac{1}{2}} = 250$  GeV,  
 $A_0 = -100$  GeV,  
 $\tan \beta = 10, \mu > 0$





A. Belyaev, Phys Rev D59 1998 015022

